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## Simulation of the growth of CVD films

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J. J. V. Jansen

### 1. INTRODUCTION

Chemical vapor deposition (CVD) is the preferred method of manufacture for solid films used in many industrially important thin and thick film applications. Requirements for the physical, mechanical and electrical properties of these films are becoming increasingly difficult to achieve, and deposition morphology plays an important role in this regard. Recently, we proposed a continuum model describing the evolution of a gas-solid interface during atmospheric pressure CVD (Viljoen *et al.*, 1994). A linear stability analysis (LSA) was used to determine the effect of reactor conditions on planar growth stability. The present paper discusses numerical solution of this model, and uses simulation examples to illustrate interface evolution under typical deposition conditions and from arbitrary initial interface shapes.

### 2. MODEL FORMULATION

We consider the growth of an amorphous solid film during a typical high pressure (0.1–1 atm), low Knudsen number ( $Kn \ll 1$ ) CVD process. Transport of gaseous species to the interface is assumed to be diffusive. Details of the model derivation are given in Viljoen *et al.* (1994) and only the final dimensionless governing equations are presented here. The gas phase concentration balance becomes:

$$\frac{\partial \mathcal{C}}{\partial \tau} = \frac{1}{Pe} \left( \frac{\partial^2 \mathcal{C}}{\partial \xi^2} + \frac{\partial^2 \mathcal{C}}{\partial \zeta^2} \right) + \frac{\partial \mathcal{C}}{\partial \zeta} \quad (1)$$

and the interface evolution equation (solid phase):

$$\frac{\partial H}{\partial \tau} = \frac{\rho}{Pe} (\nabla \mathcal{C} \cdot \mathbf{n}) \sqrt{1 + H_\zeta^2} - 1 + \chi \left[ \frac{\kappa_\xi}{\sqrt{1 + H_\zeta^2}} \right]_\xi \quad (2)$$

with boundary conditions

$$\mathcal{C} = 1 \quad \text{at } \zeta = \frac{\rho}{L} \quad (3)$$

$$\left. \frac{\partial \mathcal{C}}{\partial \xi} \right|_{\xi=0} = \left. \frac{\partial \mathcal{C}}{\partial \xi} \right|_{\xi=n\ell}$$

$$\mathcal{C}(0, \zeta) = \mathcal{C}(n\ell, \zeta) \quad \text{where } n = 0, \pm 1, \pm 2, \dots \quad (4)$$

$$\frac{1}{Pe} (\nabla \mathcal{C} \cdot \mathbf{n}) = \mathcal{G} \left( 1 - \frac{\Gamma}{L} \kappa \right) \mathcal{C} \quad \text{at } \zeta = H \quad (5)$$

and

$$\begin{aligned} H(0) &= H(n\ell), \\ H_\zeta(0) &= H_\zeta(n\ell), \\ H_{\zeta\xi}(0) &= H_{\zeta\xi}(n\ell), \\ H_{\zeta\xi\xi}(0) &= H_{\zeta\xi\xi}(n\ell), \quad \text{where } n = 0, \pm 1, \pm 2, \dots \end{aligned} \quad (6)$$

and where

$$\kappa = - \frac{H_{\zeta\xi}}{(1 + H_\zeta^2)^{3/2}}, \quad Pe = \frac{V_0 L}{D_f}, \quad \mathcal{G} = \frac{kK_{AD}}{V_0}.$$

Note that the period over which changes are considered is  $d$ . The value of  $Pe$  is extremely small for most cases since film growth occurs at such a low rate. Another important parameter is the Damkohler number,  $Da$ , which we can define as  $kK_{AD}s/D_f$ . It gives an indication of the relative magnitude of gas diffusional and surface kinetic resistance. In our system a Damkohler number  $Da \ll 1$  represents kinetic control and practically no gas diffusional resistance, while  $Da \gg 1$  corresponds to deposition in the severely diffusion-limited regime. Note that the Damkohler number we define here is equivalent to the so-called CVD number defined by Van den Brekel and Jansen (1977). They used it to characterize the state of the deposition process, after their experiments of Si deposition showed that deposition uniformity decreased with increasing values of the CVD number. Our simulation examples will show similar phenomena.

### 3. SOLUTION OF GOVERNING EQUATIONS

A previous study (Viljoen *et al.*, 1994) showed that valuable information about a deposition process can be obtained from a linear stability analysis (LSA), without solving the full system of equations. The LSA is used to predict system conditions where planar film growth and thus film uniformity can be ensured. However, it is also important to be able to predict film morphology by solving the governing equations under different sets of operating conditions and from arbitrary initial conditions.

The governing equations represent a set of coupled, non-linear partial differential equations (PDEs), which need to be solved numerically. A time scale of change can be identified for each phase:  $\tau_g = L^2/D_f$  is the characteristic time for gas diffusion and  $\tau_s = L/V_0$  the characteristic time for interface evolution. The ratio of these characteristic times gives an indication of the time scale on which to solve the problem. This ratio is given by  $\tau_g/\tau_s \equiv Pe$ . Since  $Pe \ll 1$  for all cases we will consider, it means that the gas phase is in pseudo steady state with respect to growth in the solid phase. This greatly simplifies the solution procedure. Instead of solving the governing equations simultaneously, evolution of the interface can be followed on the time scale of solid growth and at each instant in time the steady state gas phase problem can be solved independently.

The gas phase problem was solved with e.g. the Galerkin Finite Element Method, using quadrilateral elements and bilinear basis functions. For the following interface evolution, we discretized the horizontal coordinate into an equidistant grid and approximated spatial derivatives with finite difference formula. This transforms the problem into a set of ordinary differential equations (ODEs), for which we used an

explicit solution method. Stability requirements dictate a small time step, and the maximum time step can be determined with a formula similar to that defined for parabolic equations [see e.g. Finlayson (1980)]. In our case, however, spatial derivatives up to fourth order are present in the surface diffusion term. This is not common in chemical engineering applications, where typically reaction-diffusion or convection-diffusion problems are encountered. Higher derivatives tend to amplify the effect of kinks in the interface. Computational problems can usually be avoided by flooding problem areas with many points. Unfortunately, this decreases the maximum allowable time step. For example, if the number of points are doubled, the time step has to be reduced about sixteen times. For time integration a fourth order Runge-Kutta method with automatic time step adjustment was used.

#### 4. NUMERICAL EXAMPLES

With periodic boundary conditions, one can in principle choose any length of substrate on which to study deposition. However, if this length is too small, the results may not be representative of a real case. The horizontal length is therefore chosen 5 to 20 times larger than the wavelength associated with the "most" unstable perturbation. The latter is the perturbation associated with the largest eigenvalue in the LSA under the specific conditions [see Viljoen et al. (1994)]. Consider a CVD process with parameters in the range typical for the CVD growth of e.g. SiC and CrB<sub>2</sub>, thus atmospheric pressure and temperatures higher than 900 K. Typical parameter values are given in Table 1. For these parameters and at a temperature of 1100 K the wavelength of most unstable perturbation (as determined from LSA) is 85.2 μm. Therefore, an acceptable horizontal length to choose as solution domain is 500 μm.

Another important characteristic is the following: We can use the LSA to determine the time needed to double the amplitude of a protrusion with most unstable wavelength, which is the dominant one at short deposition times. At 1100 K this time is 142 h, which means that in the typical duration of deposition we will not observe a significant change in planar morphology even though planar deposition is unstable under those conditions. This situation changes drastically for deposition at 1600 K, where the wavelength of most unstable perturbation is 7.6 μm. The time to double the amplitude is now only 140 s, which means very significant changes in the morphology will take place in normal operating times.

Now consider a case where the substrate has a random initial roughness onto which a film is deposited. The effect of temperature and Damkohler number on the morphology of film growth is shown in Fig. 1. Note that the duration of each run is different, since the deposition rate increases with increasing Damköhler number. The morphology of the films grown at  $Da = 1/20$  and  $Da = 1$  is very similar, but both cases are quite different than the run for  $Da = 20$ . At low to moderate Damkohler numbers the initial randomly rough surface develops into a characteristic nodular structure,

Table 1. System parameters for typical CVD process

Parameter	Value
$s$ (m)	$10^{-3}$
$\delta$ (m)	$10^{-10}$
$kK_{AD}$ (m/s)	$2 \times 10^7 \exp[-24056/T]$
$D_f$ (m <sup>2</sup> /s)	$10^{-10} T^{1.823}$
$D_e$ (m <sup>4</sup> /s)	$\frac{\exp[-2000/T(5.0 + 0.0033T)]}{T}$
$\Gamma$ (m)	$8 \times 10^{-9}/T$

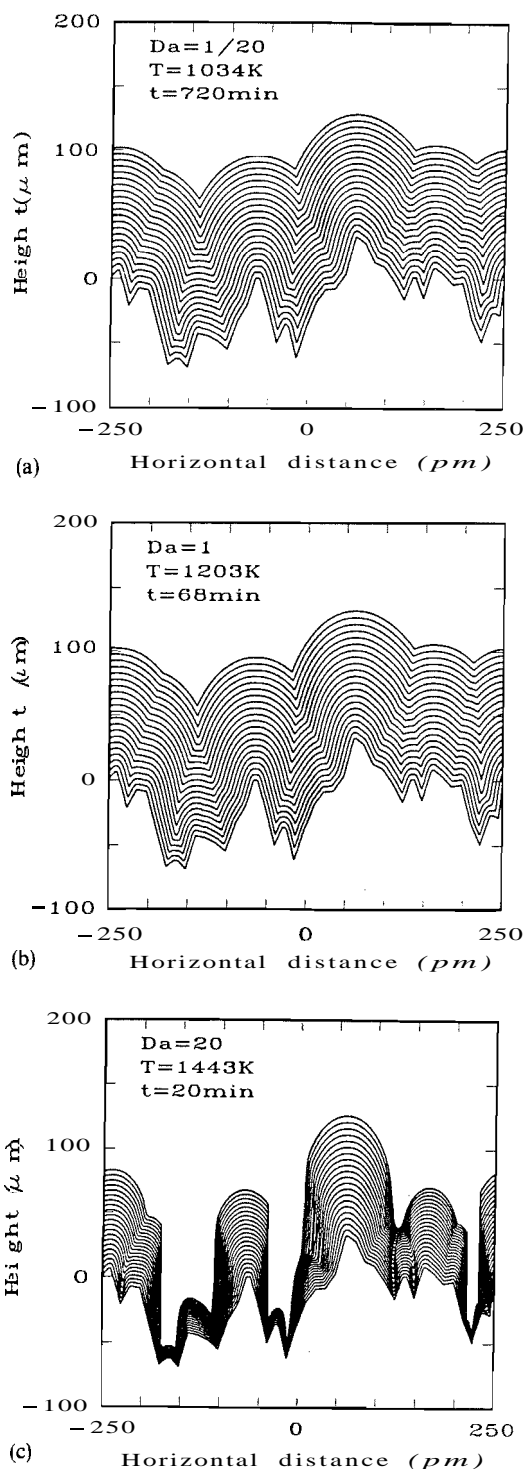
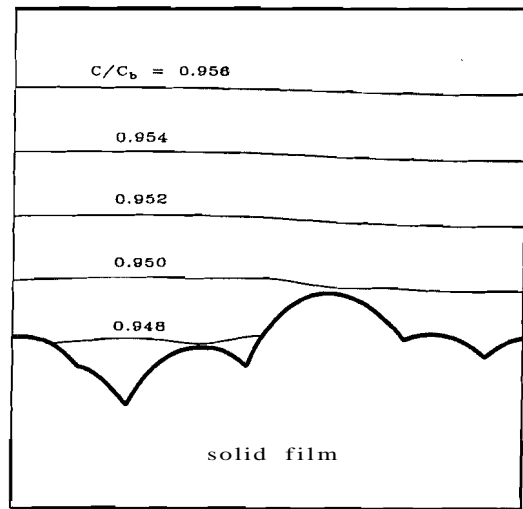
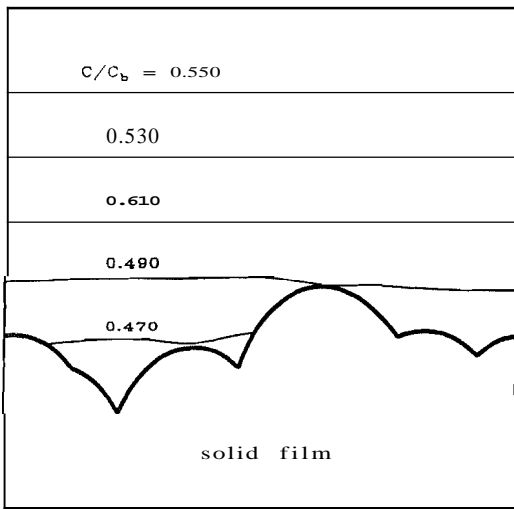


Fig. 1. Effect of temperature and thus Damkohler number on uniformity of film growth at a constant pressure of 1 atm: (a)  $Da = 1/20$ , (b)  $Da = 1$ , (c)  $Da = 20$ .

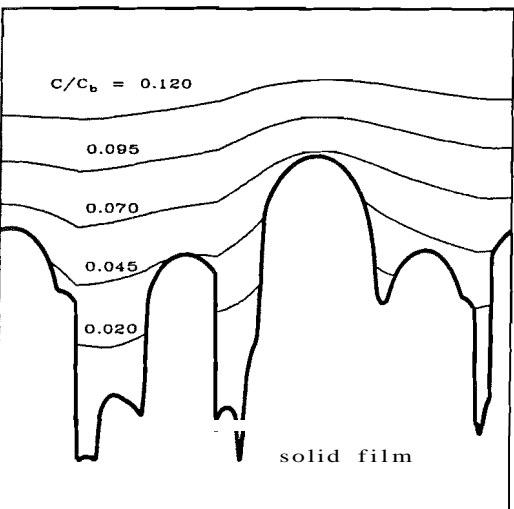
which is commonly observed in CVD experiments (Vandenbulcke and Vuillard, 1976; Chin et al., 1977; Schmidt, 1991). The results indicate that larger nodules grow at the expense of smaller ones. This growth-death phenomenon is commonly observed in deposition processes (Messier, 1986). It



(a)



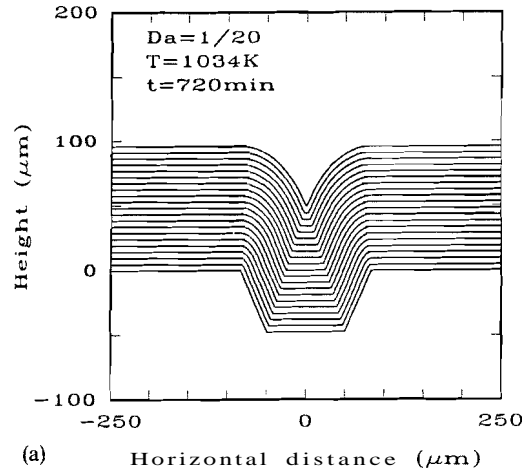
(b)



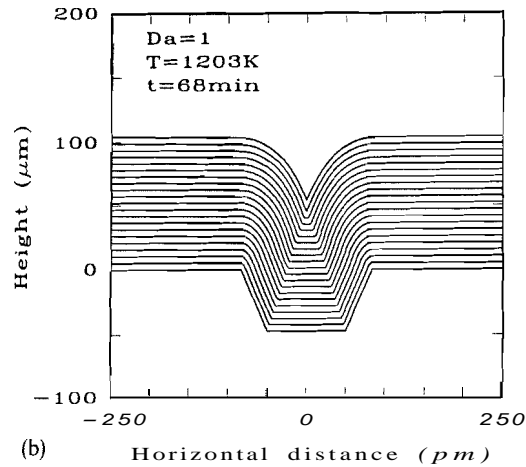
(c)

Fig. 2. Iso-concentration lines for deposition conditions in Fig. 1.

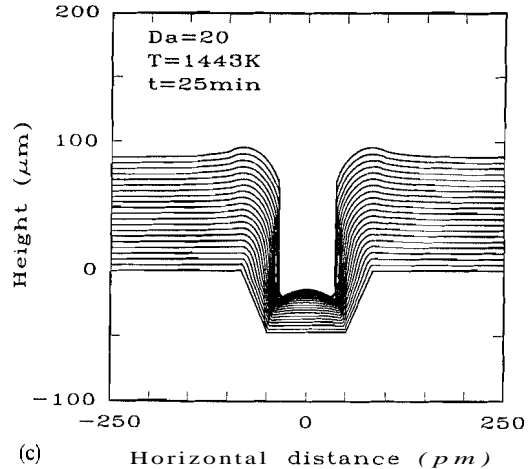
has been shown that the dominant nodule size follows a power law relation with respect to the film thickness, thus  $D_n = t^e$ , where  $D_n$  is the dominant nodule size,  $t$  the time of deposition or film thickness, and  $e$  an exponent which depends on deposition conditions. For the run at  $Da = 1$  the exponent is found to be  $e = 0.64$ , which compares well with



(a)



(b)



(c)

Fig. 3. Effect of temperature and thus Damkohler number on step coverage in a trench at a constant pressure of 1 atm: (a)  $Da = 1/20$ , (b)  $Da = 1$ , (c)  $Da = 20$ .

the predicted value of Kardar *et al.* (1986). They solved a simplified growth model analytically and found the exponent to be  $e = 2/3$ .

At  $Da = 20$  interface growth does not seem to follow the typical growth-death phenomenon. Nodules tend to develop

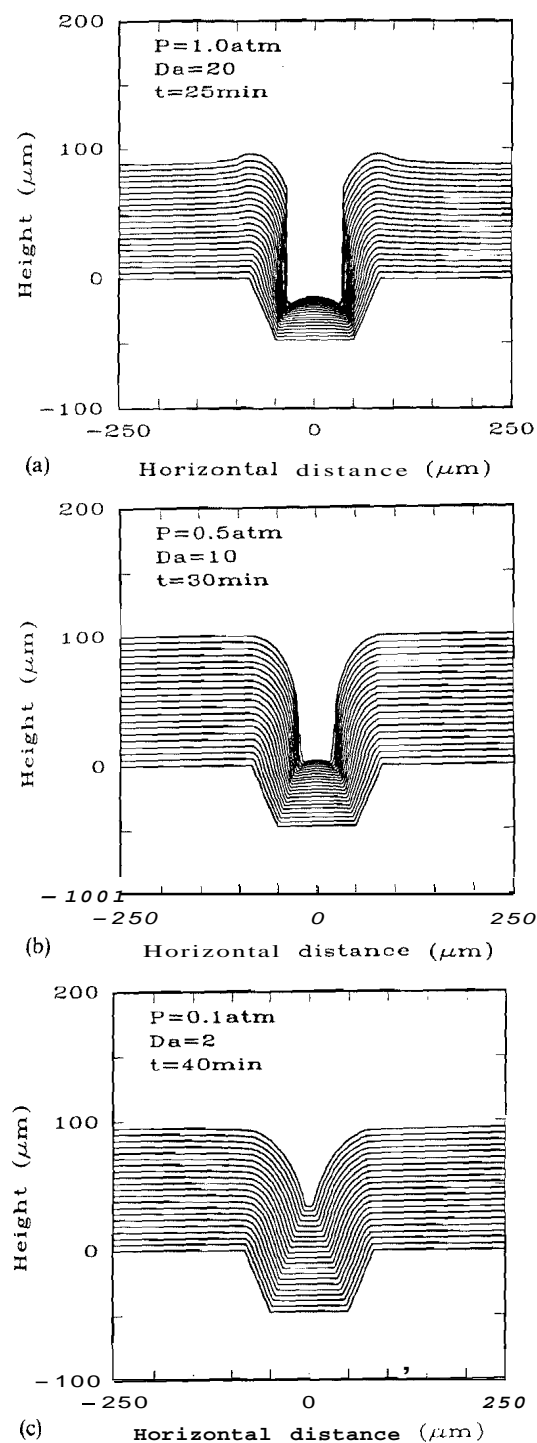


Fig. 4. Effect of reactor pressure on step coverage of a trench under diffusion-limited conditions: (a)  $P = 1.0$  atm, (b)  $P = 0.5$  atm, (c)  $P = 0.1$  atm.

into long fingers with deep grooves between them. This has been observed experimentally in the CVD growth of Si from  $\text{SiHCl}_3$  in the gas phase (van den Brekel, 1977) under diffusion-limited conditions. It is also similar to dendritic growth observed in solidification experiments where evolution of the liquid-solid interface is driven by diffusion of heat (Huang and Glicksman, 1981a, b). The formation of grooves is caused by depletion of reactant under diffusion-limited conditions. Note the tendency of some fingers to nose ahead of their neighbors. This is clear from Fig. 1(c), where reactant at the tip of the second nodule (at about the  $-140 \mu\text{m}$  position) is depleted and growth therefore inhibited. Iso-concentration lines for the runs mentioned above are shown in Fig. 2. Note how the reactant concentration close to the interface changes with increasing Damköhler number. At  $Da = 1/20$  it is close to the value in the bulk and at  $Da = 20$  it is close to zero.

When finger formation takes place, gradients of the interface ( $\partial H/\partial \xi$ ) along the sides of fingers become very large and can in fact approach infinity. The growth model represents height of the interface  $H$  as a single-valued function of horizontal coordinate  $\xi$ , and it is therefore assumed  $\partial H/\partial \xi$  is finite. As a result, computational problems are encountered when gradients become very large. To solve this problem, one can reformulate the model in a parametric form, where a different parameter than the horizontal coordinate is used. This will be addressed elsewhere.

Now consider a case where the initial interface is in the form of a duct or trench. This shape is typically found in the electronics industry, but the length scale may be quite different. It is often required that the thickness on all sides of the trench be the same. This requirement can be expressed in terms of the step coverage (Oh *et al.*, 1992), which represents the ratio of film thickness in the bottom of the trench to that at the top of the trench. The effect of temperature and thus Damköhler number on the film growth in a trench is shown in Fig. 3. Deposition parameters and physical properties were the same as in Table 1 and  $500 \mu\text{m}$  of the substrate was considered. The results clearly illustrate that step coverage decreases with increasing Damköhler number, which is similar to decrease in film uniformity observed in the previous example.

The runs in Fig. 3 were conducted at a constant pressure of 1 atm. Suppose one wants to maintain the high deposition temperature but improve step coverage at that temperature by adjusting other parameters, such as reactant dilution ratio or pressure. At such a high temperature the deposition is diffusion-limited, so step coverage should improve if the diffusional limitation can be decreased. Gas diffusion is inversely related to reactor pressure, which means step coverage should improve if the reactor pressure is decreased at constant dilution ratio. Three runs at different pressure and constant temperature of 1443 K were performed and the results are shown in Fig. 4. The graphs indicate that the step coverage indeed improves. Note that the Damköhler number decreases linearly with decreasing pressure.

## 5. CONCLUSIONS

Numerical solution of a continuum model derived for the evolution of a gas-solid interface during atmospheric pressure CVD, is described. Simulation results showed that deposition morphology is strongly dependent on the Damköhler number,  $Da$ . A significant decrease in film uniformity was observed for increasing values of the  $Da$ . This trend was also observed for growth in a trench, where step coverage was found to decrease with increasing values of  $Da$ . Under diffusion-limited conditions, step coverage was found to improve if the diffusional limitations could be decreased by decreasing system pressure at constant substrate temperature. The numerical results suggest that film uniformity and optimal step coverage can be achieved by changing operating condi-

tions in such a way that the lowest value of a Damköhler number of deposition is achieved.

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#### NOTATION

$\mathcal{C}$	dimensionless concentration, subscript: b—bulk
$D$	diffusion coefficient, subscripts: $f$ —gas, $m^2/s$ ; $e$ —surface, $m^4/s$
$Da$	Damköhler number, defined as $kK_{AD}s/D_f$ or $Pe\vartheta s/L$
$H$	dimensionless height of film
$kK_{AD}$	rate constant, $m/s$
$\ell$	dimensionless length of film considered, period
$L$	characteristic length scale, $\sqrt{D_e/\delta V_0}$ , $m$
$Pe$	Peclet number, $V_0 L/D_f$
$s$	boundary layer thickness, $m$
$V_0$	reference deposition rate, $m/s$

#### Greek letters

$\gamma$	molar volume, $m^3/\text{mole}$
$\Gamma$	capillary length, $m$
$\delta$	radius of molecule, $m$
$\zeta$	vertical coordinate, dimensionless
$\vartheta$	dimensionless parameter, defined as $kK_{AD}/V_0$
$\kappa$	dimensionless surface curvature, $\frac{H_{\xi\xi}}{(1 + H_\xi^2 + H_\eta^2)^{3/2}}$
$\xi$	horizontal coordinate, dimensionless
$\rho$	dimensionless parameter, equivalent to $\gamma C_b$
$\tau$	dimensionless time, $tV_0/L$
$\chi$	dimensionless parameter, equivalent to $D_e/V_0 L^3$

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