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# High-Resolution Thin-Film Device to Sense Texture by Touch

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# High-Resolution Thin-Film Device to Sense Texture by Touch

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## Abstract

Touch (or tactile) sensors are gaining renewed interest as the level of sophistication in the application of minimum invasive surgery and humanoid robots increases. The spatial resolution of current large-area (greater than 1 cm<sup>2</sup>) tactile sensor lags by more than an order of magnitude compared with the human finger. By using metal and semi conducting nanoparticles, a ~100-nm-thick, large-area thin-film device is self-assembled such that the change in current density through the film and the electroluminescent light intensity are linearly proportional to the local stress. A stress image is obtained by pressing a copper grid and a United States 1-cent coin on the device and focusing the resulting electroluminescent light directly on the charge-coupled device. Both the lateral and height resolution of texture are comparable to the human finger at similar stress levels of ~10 kilopascals.

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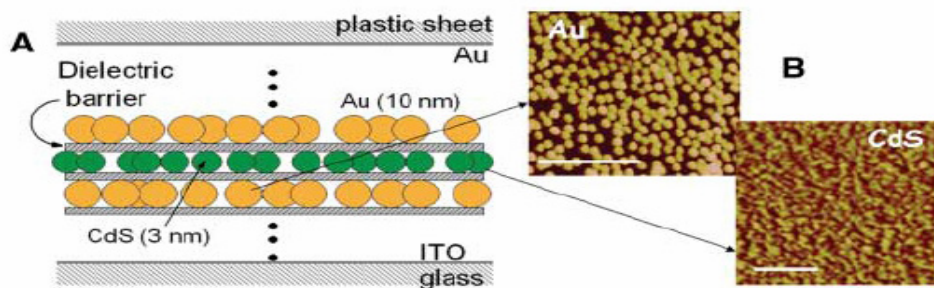
Sensation of touch, primarily the determination of stress distribution over the area of physical contact between the sensor and the object surfaces (1), is a critical component to advance minimum invasive surgical procedures by giving the surgeon a "touch sensation" (2–6) to decipher, for example, cancer tissue (3) and gallstone (2) using the signal from tactile sensor (7). Moreover, there is great interest in developing humanoid robots (8, 9) that can sense shapes (10, 11), textures (12, 13), and hardness (14) and manipulate complex objects (11), which are not readily possible by vision alone. Touch (or tactile) sensors are usually made as micro-electromechanical systems composed of micromachined deformable components (15) or by integrating chip with electronic circuit and strain sensitive materials, such as magento-resistive ceramics (16), piezoelectric polymers (17, 18), and strain sensitive conducting elastomers (19, 20). Tactile sensors from optical data have been demonstrated in which the contact stress distribution is calculated from the change in shape of the deformable sensor surface obtained by a camera (21). For small-area devices, such as an 8 by 8 array of capacitance sensors, spatial resolution of 100  $\mu\text{m}$  is demonstrated (22). However, for largearea devices,  $\sim 1\text{ cm}^2$  or larger, the spatial resolution for stress distribution is, at best, in the  $\sim 2\text{ mm}$  range (20) compared with the  $\sim 40\text{-}\mu\text{m}$  resolution achieved by the surface of a human finger (23).

We report a  $\sim 100\text{-nm}$ -thick,  $2.5\text{-cm}^2$  device, based on the principle of electron tunneling, capable of imaging stress distribution with spatial resolution of  $\sim 40\text{ }\mu\text{m}$ . The height resolution is  $15\text{ }\mu\text{m}$ , comparable to the  $\sim 2\text{ }\mu\text{m}$  for a human finger (24). Our device, composed of metal and semiconducting nanoparticles, directly converts local stress into electroluminescent light and modulation in local current density. Both the electroluminescence intensity and current density are linearly proportional to the local stress.

To demonstrate the high-resolution capability, we image stress by directly focusing the electroluminescent light signal on a chargecoupled device (CCD) camera. The sensitivity or minimum stress required for direct imaging of touch on CCD is  $\sim 9\text{ kPa}$ , which is well within the 10 to 40 kPa range that a human finger applies to sense texture and shape (25). To demonstrate the spatial resolution of the device, we use a copper grid to hold the specimen for transmission electron microscope (TEM) analysis. When the Cu

grid composed of free-standing 40- $\mu\text{m}$ -wide lines at a pitch of 220  $\mu\text{m}$  is pressed on the sample, the stress distribution from the features produces spatial variation in electroluminescent light that is directly imaged on the CCD. Furthermore, we press a U.S. 1-cent coin against the device and image the stress distribution from the embossing on the coin directly onto the CCD. The fabrication requires no lithography, and the transduction of stress distribution is virtually continuous at 100-nm scale. A simple model based on electron tunneling between the nanoparticle layers and ionic conductivity through the dielectric medium explains the characteristics of the device.

The device ([Fig. 1](#)) consists of alternating layers of Au (10 nm thick) and CdS (3 nm thick) nanoparticles separated by dielectric layers (DL), composed of stacked alternating layers of poly(styrene sulfonate) (PSS) and poly(allylamine hydrochloride) (PAH). All the monolayers of polymer and nanoparticles are deposited from solution using layer-by-layer self-assembly ([26](#), [27](#)). The particular device described here has three Au layers and two CdS layers, with four layers each of PAH and PSS as the interlying DL [see details of structure and process in ([28](#))]. The atomic force microscope (AFM) images in [Fig. 1B](#) indicate that the nanoparticles do not form long-range, (electrically) percolating clusters. Consistent with the AFM image, the film is insulating in the inplane direction, as independently confirmed by measuring two-point electrical continuity using spring-loaded flat Au (pogo) pins.

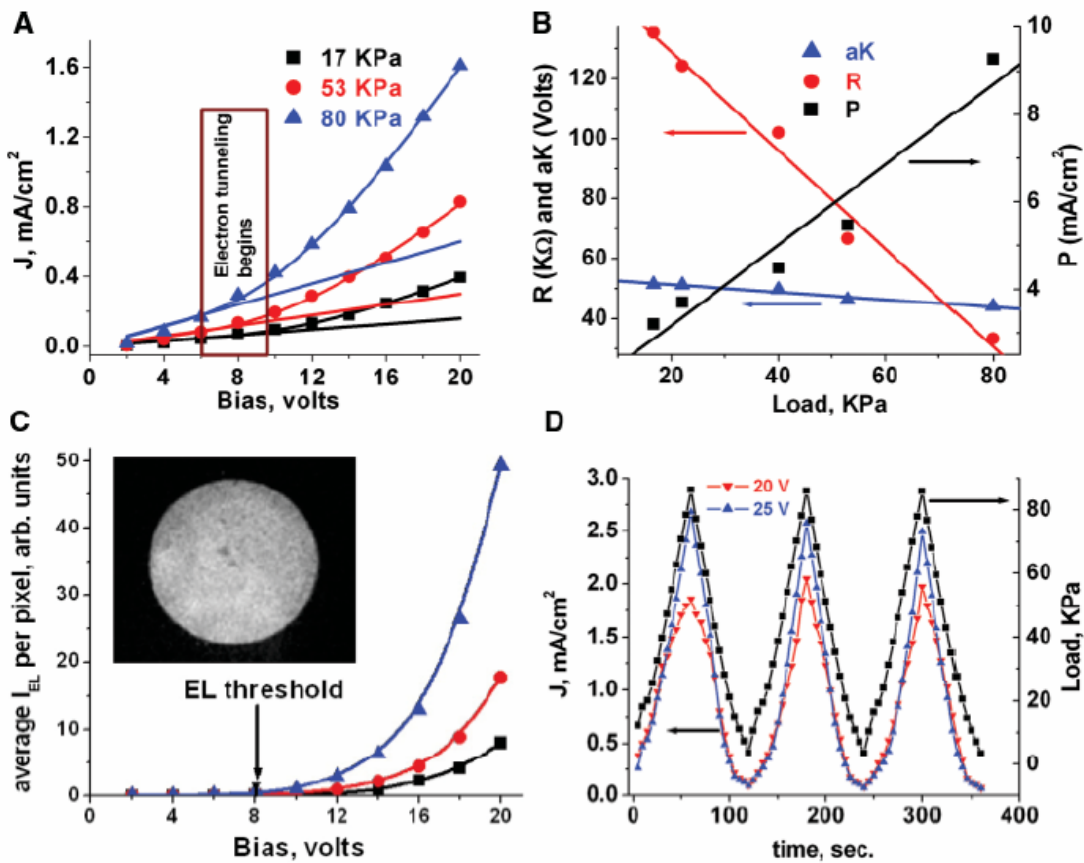


**Fig. 1.** Multilayer structure of electro-optical device. (A) Schematic of the unit structure of the device showing nanoparticle monolayers spaced by an organic DL composed of polyelectrolytes PSS and PAH. A total of three Au and two CdS nanoparticle monolayers with dielectric barriers are deposited. The top Au

electrode is coated with flexible plastic, and the bottom transparent Indium-Tin Oxide (ITO) electrode is on glass. The electroluminescent light from the CdS is measured through the bottom electrode. **(B)** AFM topography (i.e., height) image after deposition of first Au layer and the first CdS layer after that. The scale is 50 nm.

In the vertical direction, the film is conducting due to tunneling between the Au and CdS layers. [Figure 2A](#) shows the current density  $J$  through the film as a function of bias  $V$  between the bottom ITO and top Au electrode under a uniform compressive stress  $\sigma$ . The uniform stress is applied by placing an optically flat quartz disk on the flexible Au electrode. The nonlinear  $J$ - $V$  curves fit based on a model combining the field-assisted electron tunneling current [i.e., the Fowler-Nordheim equation ([29–31](#))] through the nanoparticles and the ionic (leakage) current due to ions in the polyelectrolyte. According to the relation  $J_{\text{tunneling}} = P \exp(-aK/V)$ , where,  $a$  is the (vertical) interparticle distance,  $K$  is a critical field for activated tunneling that depends on the work functions of the particles, and  $P$  at constant temperature is proportional to  $V^2$  and the number density of carriers for conduction ([29, 32](#)). The ionic leakage-current,  $J_{\text{leakage}} = V/R$ , is due to mobile ions ( $\text{H}^+$ ,  $\text{Na}^+$  and  $\text{OH}^-$ ,  $\text{Cl}^-$ ) in the polyelectrolytes, where  $R$ , the ohmic resistance, is proportional to the distance between the electrodes. A three-parameter (i.e.,  $R$ ,  $aK$ , and  $P$ ) fit for three sets of data, with corresponding straight line due to leakage  $I$ - $V$ , is an excellent match to the experiment for the whole range of  $V$  and  $\sigma$  studied ([Fig. 2A](#)). The three fitting parameters are self-consistent as applied stress increases ([Fig. 2B](#));  $R$  from the fit decreases linearly with  $\sigma$ , implying that the film deforms linearly. The estimated resistivity of  $\sim 1.6 \cdot 10^8 \Omega\text{-m}$  is reasonable for ion conductivity in PSS and PAH under ambient humidity ([33, 34](#)). The linear decrease in  $Ka$  with  $\sigma$  indicates that the interparticle spacing (i.e.,  $a$ ) also decreases linearly with  $\sigma$ , which is consistent with linear deformation of overall thickness. Assuming  $a$  (at  $\sigma = 0$ )  $\sim 5$  nm, the critical field  $K$  of  $\sim 10^9$  V/m is reasonable ([29, 32](#)). Because  $P$  depends on the carrier density, at fixed  $V$ , the linear increase in  $P$  may be attributed to a linear increase in the number of percolating channels between the electrodes as the film is deformed. Furthermore, although not shown, consistent with the Fowler-Nordheim equation,  $P \sim V^2$  (at constant load) for the complete range of operating conditions ( $>8$  V). The curve fit in [Fig. 2A](#) indicates that the tunneling

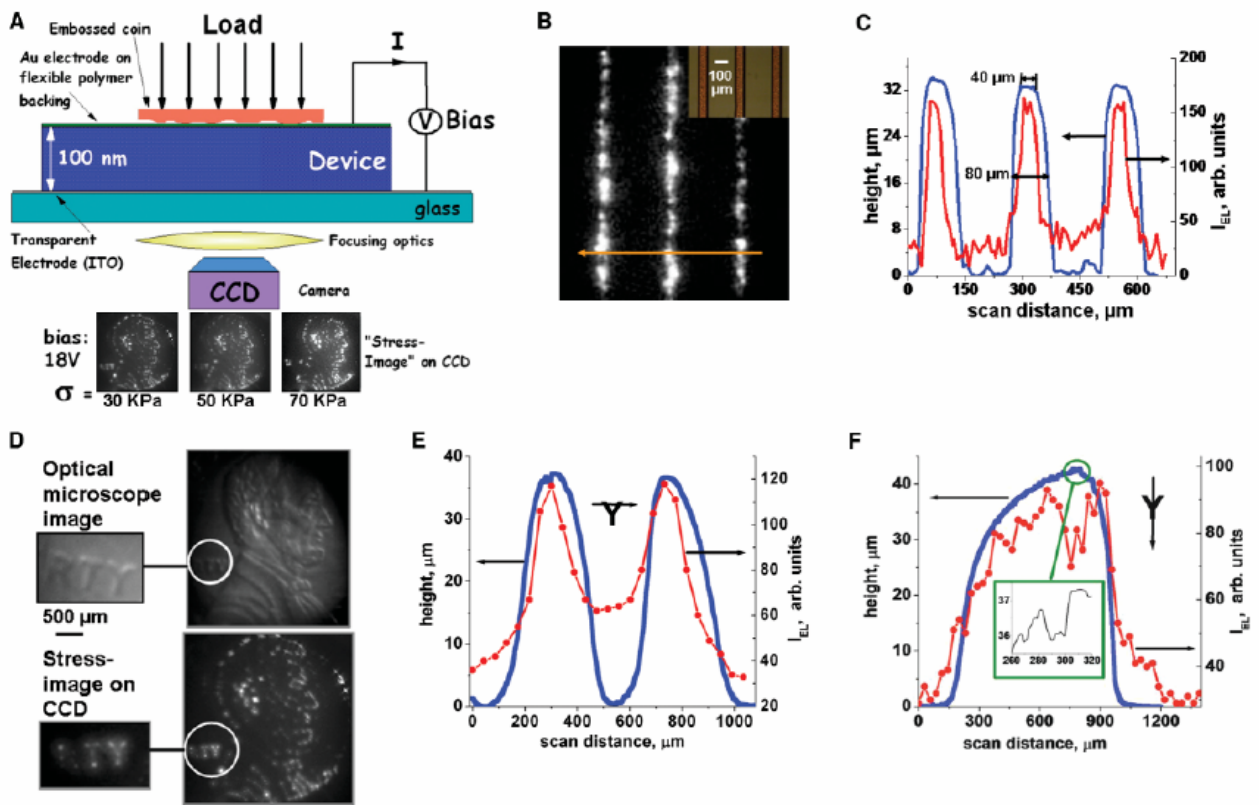
process becomes important beyond  $\sim 8$  V (indicated as a box). Electroluminescence will occur only when tunneling through CdS occurs. Hence, the electroluminescence intensity  $I_{EL}$  should be measurable above the bias of  $\sim 8 \pm 1$  V.  $I_{EL}$  as reported is intensity (in arbitrary units) per 16 by 16  $\mu\text{m}$  pixel of the CCD camera. The noise level based on imaging a blank sample (i.e., dark current) is 1 arbitrary unit. The  $I_{EL}$  threshold at  $\sim 8$  V further supports the charge-transport model (Fig. 2C). The average  $I_{EL}$  per pixel is limited to 80 kPa because the electroluminescence from the edge of the disk is enhanced as a result of "square-edge-punch" stress concentration (35). Furthermore, as  $J_{\text{tunneling}}$  increases exponentially with stress compared with a linear increase for  $J_{\text{leakage}}$ , at larger bias (i.e.,  $>18$  V) tunneling dominates over leakage current. The deformation is reversible with nominal hysteresis (Fig. 2D).



**Fig. 2.** Electro-optical characteristics of the device and electron-transport mechanism. The device is compressed using an optically flat quartz disk on the flexible Au electrode. (A)  $J$ - $V$  data points at various  $\sigma$ . The fitted line is based on

*a model combining charge transport by electron tunneling and ions. The straight line corresponds to  $J_{leakage}$ . The box indicates the expected onset of electroluminescence. (B) The three fitting parameters estimated from fitting the  $J$ - $V$  curve show linear dependence, indicating that the deformation is linear with respect to load. (C) The electroluminescence from the CdS nanoparticles for  $\sigma$  in 17 to 80 kPa range commences at  $\sim 8$  V. The inset is the electroluminescence image on the CCD camera below the ITO electrode corresponding to the 1.1-cm-diameter quartz disk pressing the top electrode (Fig. 3A). Uniform light indicates uniform load distribution. The data symbols correspond to Fig. 2A. (D) The current  $J$  monitored under dynamic loading of the device at a fixed bias of 20 and 25 V indicates reasonably low hysteresis. The line is to guide the eye.*

To assess the capabilities of the device for measuring surface topography and texture, a TEM specimen holder from Electron Microscopy Sciences (i.e., a Cu grid) and a U.S. 1-cent coin was pressed against the flexible Au electrode of the device, and the resulting electroluminescent light from the device was focused on the CCD camera below the ITO/glass substrate (Fig. 3A). The embossing of President Lincoln's head on the coin is apparent in the stress image obtained on the CCD. The overall intensity of the stress image increases with average stress  $\sigma$  (Fig. 3A). The TEM grid lines have a trapezoidal cross-section with a top base of 40  $\mu\text{m}$  that is in contact with the device, a width of 80  $\mu\text{m}$  at half height, and interline spacing of 230  $\mu\text{m}$  (Fig. 3B, inset, and Fig. 3C). Lines of the Cu grid (seen in the optical microscope image) are clearly resolved in the stress image (Fig. 3B). Comparison between the topography and  $I_{EL}$  scan of the TEM Cu grid shows that the flat surface of 40- $\mu\text{m}$  line width in contact with the device is well resolved, leading to a lateral spatial resolution of at least 40  $\mu\text{m}$ .



**Fig. 3.** Characteristics of the pressure imaging device. (A) Schematic of the setup to image texture of a metal coin; also shown are the pressure images taken on CCD camera at three different compressive stresses. The diameter of the coin is  $\sim 2$  cm, and the device is a 2.5-cm square. The bias between the coin and the ITO electrode is fixed at 18 V. (B) Comparison between optical microscope (inset) and pressure image (at  $V = 18$  V and  $\sigma = 40$  kPa) of a TEM Cu grid. The horizontal line in the pressure image is the  $I_{EL}$  scan shown in Fig. 3C. (C) Comparison between the topographic and  $I_{EL}$  scans across the TEM Cu grid. (D) Comparison between optical microscope and pressure image (at  $V = 18$  V and  $\sigma = 30$  kPa) of the coin, showing the finer structure. The magnified image shows the letters "RTY" of the word "LIBERTY" on the coin. (E) and (F) Comparison between the topographic and  $I_{EL}$  scan across letter "Y" of "LIBERTY." The inset of (F) shows the local roughness of the coin.

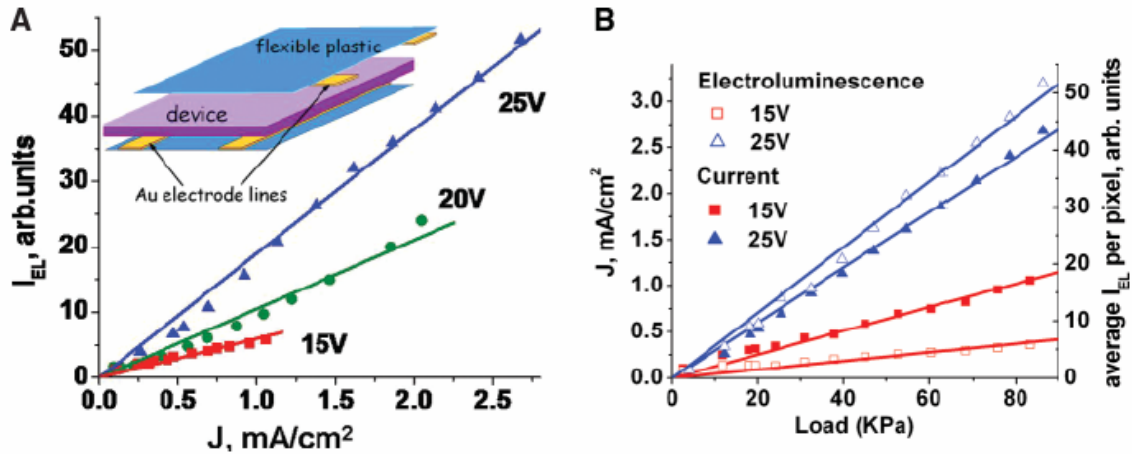
The remarkable correspondence between the optical microscope image and the stress image seen in Fig. 3, A, D, and E, indicates that the variation in stress distribution caused by the embossing on a U.S. 1-cent coin leads to significant change in local current density, which in turn is manifested as modulation in  $I_{EL}$ . The device can decipher fine



features, such as wrinkles on the clothing of President Lincoln and the letters "TY" in "LIBERTY," observed in the optical microscope and the stress image (Fig. 3D). Because the electroluminescence intensity from the device is virtually continuous, the resolution of the stress image is determined by the optics and the CCD camera. For the CCD with a 512 by 512 array of  $\sim 16\text{-}\mu\text{m}$  pixels, the  $I_{\text{EL}}$  is well over noise level (Fig. 3, C, E, and F), indicating that the electroluminescence is enough to achieve lateral spatial resolution of at least  $\sim 20\text{ }\mu\text{m}$  (as the area of chip and disk are comparable). However, the stress will be "smeared" as a result of the thickness of the flexible Au electrode, leading to higher resolution. Furthermore, from Fig. 3, E and F, height modulation of  $\sim 10\text{ }\mu\text{m}$  is measurable, indicating that the device is fairly deformable. The high compliance is perhaps due to some interpenetration among the nanoparticle layers. The inset in Fig. 3F shows that the embossing is rough at the micron scales. This local roughness will cause nonuniformity in the surface charge distribution, leading to "hot spots" of high electric field strength. The "pointillist" pressure images are attributed to these hot spots.

The relationship between the optical (i.e.,  $I_{\text{EL}}$ ) and electronic (i.e.,  $J$ ) signal from the device pressed by the quartz disk is linear (Fig. 4A), implying that the stress image from either signal would be similar. Above, we have demonstrated the tactile sensor principle by measuring  $I_{\text{EL}}$ . Extension to a fully electronic device by measuring  $J$ , required for certain applications (such as surgery), is possible using standard microelectronic circuitry. For example, an electronic device could be made by self-assembling the multilayer thin-film between mutually perpendicular parallel lines of Au (electrodes) on flexible plastic backing (i.e., flex-circuit). The stress distribution would be imaged by mapping  $J$  at the intersection points of the top and bottom electrode lines (similar to a display device). Both  $I_{\text{EL}}$  and  $J$  increase linearly with  $\sigma$ ; however, the sensitivity to bias for an optical device is larger (by a factor of about 3), which implies that the latter has a larger dynamic range (Fig. 4B). For  $I_{\text{EL}} \sim 5$  (corresponding to a signal-to-noise ratio of 5) the sensitivity of the optical device is  $\sim 9\text{ kPa}$  at a bias of 25 V. Assuming a minimum measurable current signal of 1 nA for the electronic device proposed in the inset of Fig. 4A, for the intersection point of  $\sim 20\text{ }\mu\text{m}^2$  (36) a sensitivity of  $\sim 10\text{ kPa}$  would be possible. Because the resolution of  $20\text{ }\mu\text{m}$  is comparable to the CCD

pixel size used in the optical device, a spatial resolution of  $\sim 50 \mu\text{m}$  at a sensitivity of 10 kPa should be possible for the electronic device.



**Fig. 4.** Comparison of electronic (i.e., current) and optical (i.e., electroluminescence) signals. (A)  $I_{EL}$  is linearly proportional to  $J$ . Each data point is at fixed bias with varying  $\sigma$ . The slope at the three biases is 6.0, 10.5, and 19.0 (arbitrary units)  $\text{cm}^2/\text{mA}$ . The efficiency of the device to convert current to photoluminescence increases with bias as indicated by the slope. (B) Both  $J$  and  $I_{EL}$  increase linearly with  $\sigma$  (i.e., load) for the 15 to 25 V bias range studied. The change in slope as bias increases from 15 to 25 V for  $J$  and  $I_{EL}$  is 2.34 and 7.63, respectively.

There are two practical advantages of the described tunneling-based device compared with the reported tactile imaging devices based on deformable membranes or stress-sensitive materials. First, we achieve high sensitivity because tunneling depends exponentially on displacement (i.e., strain). Second, the self-assembly process involves no complex lithography, making it easier to build the device directly on surfaces of large area and complex shapes.

## References and Notes

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by 1-inch ITO-coated (400 nm) glass surface in piranha solution ( $\text{H}_2\text{SO}_4/\text{H}_2\text{O}_2$ ; 3:1 by volume) to make the electrode surface negatively charged. The final structure of the device is glass/ITO/DL/[Au-NP/DL/CdS-NP/DL]<sub>2</sub> /Au-NP/DL/Au-electrode/plastic. The structure of the flexible electrode is Au (200 nm)/Cr (70 nm)/Al (12  $\mu\text{m}$ )/siloxane rubber ( $\sim 5 \mu\text{m}$ ). The Au is sputter-deposited on the Al foil with Cr as the adhesion layer. Poly(dimethyl siloxane) is spincoated on the other side for the electrode, followed by UV cross-linking. The resultant freestanding electrode is physically placed on the device with the Au side in physical contact with the device surface.

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