Computer simulation of radio frequency heating of model fruit immersed in water

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1. Introduction

Radio frequency (RF) heating has been explored for rapid postharvest disinfestation of agricultural commodities (Frings, 1952; Nelson & Payne, 1982; Wang, Tang, Johnson, Mitcham, Hansen, Cavalieri et al., 2002). A number of problems remain to be addressed before RF heat treatments can be successfully used in commercial applications. The most significant problem associated with RF heating is that heating is not uniform in fresh fruits (Tang, Ikediala, Wang, Hansen & Cavalieri, 2000). Large temperature variations among and within fresh fruit reduce the effectiveness of a treatment and may cause severe thermal damage to its quality. In order to overcome non-uniform RF heating of fruits, Ikediala, Hansen, Tang, Drake & Wang (2002) suggested a saline water immersion technique, but for large fruits such as apples and oranges uneven heating inside the fruit was still unacceptable. Birla, Wang, Tang & Hallman (2004) showed improvement in heating uniformity in oranges and apples when the fruits were immersed in water and kept in motion by water jets during RF heating. However, even after rotating and moving fruits in the RF field, non-uniform heating was evident in some fruits. Computer simulation model could be an effective tool to examine the causes of non-uniform heating and in finding solutions to overcome this problem.

Neophytou and Metaxas (1996; 1997; 1998; 1999) started electromagnetic field modeling for industrial scale RF heating systems. Their work involved comparison of solution from both electrostatic and wave equations. They recommended that for small sized applicators, solution of Laplace equation is adequate, whereas for large size electrodes, wave equations should be used. Recently Chan, Tang & Younce (2004) developed...
Nomenclature

<table>
<thead>
<tr>
<th>Symbol</th>
<th>Definition</th>
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<tbody>
<tr>
<td>$C_p$</td>
<td>Specific heat (J kg$^{-1}$ K$^{-1}$)</td>
</tr>
<tr>
<td>$d$</td>
<td>Thickness, (m)</td>
</tr>
<tr>
<td>$E$</td>
<td>Electric field strength (V m$^{-1}$)</td>
</tr>
<tr>
<td>$f$</td>
<td>Frequency (Hz)</td>
</tr>
<tr>
<td>$F_b$</td>
<td>Buoyancy force (N m$^{-3}$)</td>
</tr>
<tr>
<td>$g$</td>
<td>Gravitational acceleration (m s$^{-2}$)</td>
</tr>
<tr>
<td>$h$</td>
<td>Heat transfer coefficient (W m$^{-2}$ K$^{-1}$)</td>
</tr>
<tr>
<td>$j$</td>
<td>Complex number operator</td>
</tr>
<tr>
<td>$k$</td>
<td>Thermal conductivity (W m$^{-1}$ K$^{-1}$)</td>
</tr>
<tr>
<td>$P$</td>
<td>Pressure (N m$^{-2}$)</td>
</tr>
<tr>
<td>$Q$</td>
<td>Power Density (W m$^{-3}$)</td>
</tr>
<tr>
<td>$T$</td>
<td>Temperature (K)</td>
</tr>
<tr>
<td>$t$</td>
<td>Time (s)</td>
</tr>
<tr>
<td>$u$</td>
<td>Velocity (m s$^{-1}$)</td>
</tr>
<tr>
<td>$V$</td>
<td>Electric potential (Volt)</td>
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Greek symbols

<table>
<thead>
<tr>
<th>Symbol</th>
<th>Definition</th>
</tr>
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<tbody>
<tr>
<td>$\rho$</td>
<td>Density (kg m$^{-3}$)</td>
</tr>
<tr>
<td>$\sigma$</td>
<td>Ionic conductivity (S m$^{-1}$)</td>
</tr>
<tr>
<td>$\epsilon'$</td>
<td>Dielectric constant</td>
</tr>
<tr>
<td>$\epsilon''$</td>
<td>Dielectric loss factor</td>
</tr>
<tr>
<td>$\varepsilon_o$</td>
<td>Permittivity of free space (F m$^{-1}$)</td>
</tr>
<tr>
<td>$\omega$</td>
<td>Angular frequency, (rad s$^{-1}$)</td>
</tr>
<tr>
<td>$\nabla$</td>
<td>Delta operator</td>
</tr>
<tr>
<td>$\beta$</td>
<td>Expansion coefficient (K$^{-1}$)</td>
</tr>
<tr>
<td>$\mu$</td>
<td>Viscosity (Pa s)</td>
</tr>
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</table>

Subscripts

<table>
<thead>
<tr>
<th>Symbol</th>
<th>Definition</th>
</tr>
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<tbody>
<tr>
<td>air</td>
<td>Air</td>
</tr>
<tr>
<td>mat</td>
<td>Material</td>
</tr>
<tr>
<td>rms</td>
<td>Root mean square</td>
</tr>
<tr>
<td>w</td>
<td>Water</td>
</tr>
</tbody>
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2. Materials and Methods

2.1. Physical model

A 12 kW parallel plate RF heating system (operating at 27.12 MHz, Strayfield Fastran with E-200, Strayfield International Limited, Wokingham, UK) was used in this study. The RF heating system was consisted of a generator, an applicator and a metallic enclosure. The applicator had a pair of parallel plates in which the top electrode, with adjustable height, was inductively coupled to the tank oscillator circuit via feed strips. The dielectric material was sandwiched between two electrodes to form a working circuit (Wang, Ikediala, Tang, Hansen, Mitcham, Mao et al., 2001). RF heating involves EM interaction of the lossy material (dielectric material having considerable power dissipation capability) placed between the two plate electrodes. The absorbed RF power per unit volume ($Q$, W m$^{-3}$) in the material is proportional to the square of the electric field strength ($E$, V m$^{-1}$) and directly proportional to the dielectric loss factor ($\epsilon''$) and the frequency ($f$, Hz) (Choi & Konrad, 1991):

\[
Q = 2\pi f \varepsilon_0 \varepsilon'' E_{\text{rms}}^2 = \pi f \varepsilon_0 \varepsilon'' |E|^2
\]

where $\varepsilon_0$ is the permittivity in free space (8.85×10$^{-12}$ F m$^{-3}$), and $E_{\text{rms}}$ is the root mean square value of the electric field which is equal to $\sqrt{2}$ times of the E-field amplitude.

Evaluation of the absorbed RF power density at any point inside the material requires the value of ‘$E$’ which is the function of the geometry and dielectric properties of the object, and electrode configuration (Marshall & Metaxas, 1998). Moreover, dielectric properties are temperature dependent, hence estimation of $Q$ at any point requires simultaneous solving EM field and heat transfer equations.
EM waves between 1 and 200 MHz frequency fall within the radio band, and at 27 MHz the wavelength is about 11 m in free space. In the present study, the electrode size (1.05 m × 0.80 m) was very small compared with the wavelength of 27.12 MHz used in our study. Therefore the primary mode of RF energy interaction was modeled as a quasi-static electrical field between two electrodes. Our approach was to systematically study the effect of the parameters that would significantly influence the heating patterns. Those parameters included the geometry of the object, relative location of the object in between the electrodes, and dielectric properties of the object and its surroundings.

2.1.1. Governing equations

The electric field at any point inside the electrodes is governed by Eq. (2) derived from a quasi-static approximation of Maxwell’s equations (Choi & Konrad, 1991):

\[-\nabla \cdot \left( \sigma \nabla V \right) = 0\]

where \(j = \sqrt{-1}\), \(\varepsilon'\) is the dielectric constant, and \(\sigma\) is the temperature-dependent conductivity (S m\(^{-1}\)) and related to the dielectric loss factor as \(\sigma = 2 \pi f \varepsilon_{\infty} \varepsilon'\) in the RF range (Guan, Cheng, Wang, & Tang, 2004). The scalar voltage potential \((V, V)\) is related to the electric field by \(E = -\nabla V\), hence temperature dependent dielectric properties were sufficient to calculate the electric field strength at each point in the considered domain.

The geometry of most fruits such as cherries, apples, and oranges can be approximated as a sphere. Thermal diffusion in solid objects, for example fresh fruit, is governed by the Fourier transient transfer equation, but when immersed in fluid (water or air) it is governed by momentum transport and continuity equations (Zhang, Jackson, & Ungan, 2000):

Continuity equation \(\nabla \cdot \mathbf{u} = 0\)

Energy conservation equation \(\frac{\partial T}{\partial t} + \mathbf{u} \cdot \nabla T = \frac{k}{\rho C_p} \nabla^2 T + \frac{Q}{\rho C_p}\)

Momentum equation \(\rho \frac{\partial \mathbf{u}}{\partial t} + \rho (\mathbf{u} \cdot \nabla) \mathbf{u} = -\nabla P + \mu \nabla^2 \mathbf{u} + \mathbf{F}_b\)

where \(C_p\) is the specific heat (J kg\(^{-1}\)K\(^{-1}\)), \(k\) is the thermal conductivity (W m\(^{-1}\)K\(^{-1}\)), \(t\) is the time (s), \(T\) is the temperature (K), \(\rho\) is the density (kg m\(^{-3}\)), and \(\mu\) is the viscosity of the medium (Pa s). Similar to Eqs. (4) and (5) with the velocity component \(u\) (m s\(^{-1}\)) in x-direction, the momentum and energy transport in y and z directions can be written with corresponding velocity components \(v\) and \(w\). The numerical solution of Eq. (5) requires pressure \((P, N\ m^{-2})\) to specify at least one point of the domain.

Buoyancy force \(\mathbf{F}_b\) (N m\(^{-3}\)) in z-direction can be calculated from the Boussinesq approximation, whereby variations in density due to temperature are ignored but change in buoyancy force due to differential temperature is considered for the set up of a flow field.

\[\mathbf{F}_b = \rho_\infty g \beta \Delta T\]

where \(g\) is the gravitational acceleration (m s\(^{-2}\)), \(\beta\) is the volumetric thermal expansion coefficient (K\(^{-1}\)), \(\Delta T\) is the excess temperature from the reference one (K), and \(\rho_\infty\) is the fluid density at reference temperature (kg m\(^{-3}\)). If the buoyancy force is the sole cause of motion, the convection is defined as free convection.

2.1.2. Geometric model and boundary conditions

It is impossible to take into account every details of the RF machine parts in the computation domain, as it demands excessive computer resources (Chan et al., 2004). One quarter of the RF electrode and the container was modeled to take advantage of the system symmetry. The upper electrode plate was drawn as an embedded element in 3-D to avoid a large number of mesh elements for a thin electrode plate. Symmetric surfaces were assigned thermal and electrical insulating boundary conditions as shown in Fig. 1.

Source electric potential \((V)\) was applied to the upper electrode of the applicator and the bottom of the enclosure served as an electrical ground return \((V = 0)\). An electrically shielding boundary condition was applied to the metallic enclosure walls such that \(n \cdot (\sigma \nabla V) = 0\) where \(n\) is the unit vector normal to the surface of the wall. A thermal boundary condition of convective heat transfer was applied to the outer surfaces of the container (405 mm dia. × 127 mm height made of 6.25 mm thick polyethylene sheet). The water-air interface in the container was assigned to a convective heat transfer boundary condition. The convective heat transfer coefficient was assumed to be 20 W m\(^{-2}\) K\(^{-1}\) for water-air and container-air interfaces (Wang, Tang, & Cavalieri, 2001). At fruit-water and container wall-water interfaces no-slip boundaries were imposed in the momentum equations. At the water-air interface, the velocity in the normal direction \((n)\) and shear stress in the horizontal direction were assumed to be zero. Since solving Navier-Stokes equations needs to specify pressure at least at one point in the domain, atmospheric pressure was assigned at the top corner point of the container.

2.1.3. Solution methodology

Commercial finite element method based software FEMLAB (V3.2, COMSOL Multiphysics, Burlington, MA, USA) was used to solve the coupled electromagnetic, momentum and transport equations. Various steps involved in the modeling with FEMLAB software are outlined in Fig. 2. The computational scheme was to simultaneously solve the highly non-linear equations. An unstructured mesh consisting of Lagrange quadratic elements was created. To solve temperature and velocity values at the interfaces of spherical objects and surrounding media accurately, a relatively fine mesh was generated near the interfaces. A FEM solution was considered converging until the difference in maximum temperature between successive calculations was less
than 0.1 % for a doubling of the number of elements. In the convergence studies, simulation results were found to be meshing independent when mesh sizes were reduced to 1/20th of the maximum dimension of the domain. The optimum mesh generation yielded 3,641 nodes and 26,880 elements. The ‘UMFPACK’ differential equation solver in FEMLAB was used to achieve convergence. All computer simulations were performed on a Dell 670 work station with two each Dual-Core, 2.80 GHz XEON processors and 12 GB RAM running a Windows XP 64-bit operating system. Since the time-dependent equations are solved implicit, the accuracy is the only limiting factor for the time-step determination. One can specify the limits(max/min) for the time-step control for desired accuracy (COMSOL FEMLAB 3.3 User Guide, 2006). Simultaneous solutions of coupled transient equations took 12 h for a 10 min RF heating process with 0.001s initial and 1 s maximum time steps.

2.1.4. Model input

Modeling of the RF heating process requires knowledge of dielectric, thermal and physical properties of the load and surrounding medium. Density, thermal conductivity, and specific heat were assumed to be temperature-independent, whereas the dielectric properties of the materials were assumed to be temperature-dependent over the treatment temperatures range from 20 to 60°C for postharvest pest controls. Regression expressions obtained from the analysis of reported dielectric properties data were used in the model. Table 1 summarizes the properties of the various materials used in the simulations.

In the RF range, it is impractical to measure impressed voltage between two electrodes without distorting the electric field (Marshall & Metaxas, 1998). Metaxas (1996) showed that for a typical industrial-scale system the voltage varies by only 7% between standby and full load. Therefore, a constant electric potential on the upper electrode is a realistic assumption. The analytical solution of Laplace equation for the electric potential between RF electrodes and heat dissipation in a slab
sandwiched between electrodes can be coupled together and an expression for electric potential on upper electrode can be obtained as (Metaxas, 1996):

\[
V = \left( \frac{d_{\text{air}}}{\sqrt{(\varepsilon')^2 + (\varepsilon'')^2}} + d_{\text{mat}} \right) \sqrt{\frac{\rho C_p}{\pi f \varepsilon_0 \varepsilon''}} \frac{dT}{dt} \quad (7)
\]

where \(d_{\text{air}}\) and \(d_{\text{mat}}\) are air gap and slab thickness (m), respectively.

The electric potential was estimated using the heating rate of a 1% gel slab (405 \times 405 \times 47 mm\(^3\)) in the 12 kW RF unit. Procedure details are elaborated in the next section.

2.2. Model validation

2.2.1. Model fruit preparation

As a first step to validate the simulation model, it was more convenient to use a consistent homogenous spherical object. Thus, we developed a homogeneous model fruit made of 1% gellan gel. Wang, Tang, Cavaliere & Davis (2003a) used gellan gel to prepare a model insect to study differential heating between insects and walnuts. One percent gellan gum dispersion (Kelco Division of Merck and Co., San Diego, CA) was added into the hot gel. Kelco Division of Merck and Co., San Diego, CA) was poured into the container that was completely immersed the gel balls.

Experimental heating patterns were obtained by varying the vertical position of the model fruits between electrodes with respect to the bottom electrode. The core temperature of centrally placed fruit was recorded and logged during the RF heating experiment using a fiber optic sensor (UMI-8, FISO Technology, Quebec, Canada). In additional experiments, five model fruits were RF heated in a fruit mover developed by Birla et al. (2004) to study the effect of rotation and movement of fruits during RF heating. The fruit mover had twelve numbers of water spray nozzles mounted on the periphery of a square container to rotate fruits on their axes and in a circular path. Details of the fruit mover can be found in our earlier publication (Birla et al., 2004). The RF heated model fruit was immediately bisected vertically and a thermal image was recorded with the

### Table 1. Electrical and physical properties of gel ball and water used in simulation

<table>
<thead>
<tr>
<th>Material properties</th>
<th>Gel ball</th>
<th>Tap water</th>
<th>Polypropylene</th>
<th>Air</th>
</tr>
</thead>
<tbody>
<tr>
<td>Thermal conductivity (k), W m(^{-1})K(^{-1})</td>
<td>0.53</td>
<td>0.56</td>
<td>0.2</td>
<td>0.025</td>
</tr>
<tr>
<td>Density (\rho), kg m(^{-3})</td>
<td>1010</td>
<td>1000</td>
<td>900</td>
<td>1.2</td>
</tr>
<tr>
<td>Specific heat (C_p), J kg(^{-1})K(^{-1})</td>
<td>4160</td>
<td>4180</td>
<td>1800</td>
<td>1200</td>
</tr>
<tr>
<td>Viscosity (\mu), Pa s</td>
<td>N/A</td>
<td>0.001</td>
<td>N/A</td>
<td>0.000001</td>
</tr>
<tr>
<td>Exp. coefficient (\beta), K(^{-1})</td>
<td>N/A</td>
<td>0.0003</td>
<td>N/A</td>
<td>0.002</td>
</tr>
<tr>
<td>Dielectric constant (\varepsilon')</td>
<td>-0.21T+86.76</td>
<td>-0.48T+84.74</td>
<td>2.0</td>
<td>1</td>
</tr>
<tr>
<td>Loss factor (\varepsilon'')</td>
<td>4.36T+129.4</td>
<td>0.33T+11.1</td>
<td>0.0023</td>
<td>0.0023</td>
</tr>
</tbody>
</table>

* (Wang et al., 2003a; Wang et al., 2003b)  T - Temperature in °C
* (Rehman, 1995, pp. 179, 225)
* (Stogryn, 1971)
* (von Hippel, 1995, p. 327)
infrared camera for one of the cut surfaces within 10 s to avoid the surface cooling.

Simulation and experimental temperature data at core of the model fruit were subjected to statistical analysis. A parameter defined as root mean square (RMS) for temperature differences between simulation and experiment was calculated to validate the simulation results.

2.3. Modeling of fruit rotation

It has been demonstrated experimentally that rotation and movement of fruit during RF heating improves the heating uniformity (Birla et al., 2004). However, simulation of RF heating of rotating and moving object is very difficult because of the associated moving boundary conditions. To simplify simulation for this study, spherical objects were assumed to rotate on their own axes and each spot on any particular orbit was equally exposed to RF energy. The spherical object was further assumed to be composed of six concentric layers and RF energy absorbed on a particular layer to be uniform because of the 3-D rotation. This assumption was based on a concentric temperature contour over the fruit cross section observed with thermal images (Birla, Wang & Tang, 2006). In computer simulation, power density was integrated over the individual layer and the total power density was divided by volume of the layer to estimate an average power density over the individual layer.

3. Results and Discussion

3.1. Estimation of electric potential

Fig. 3a shows the experimental temperature profile of the gel slab subjected to 10 min RF heating in a 160 mm electrode gap. Using the average slab temperature rise (12.7°C) and dielectric properties of the gel at 25°C in Eq. (7), the input voltage was estimated to be 8,162 V. Since Eq. (7) is valid for an infinitely large slab and electrode dimensions (based on the assumption of no fringing field), the estimated voltage was not the actual value. In simulation, the estimated value of \( V \) was used as a starting value and various input voltage values were tried for predicting the realistic transient temperature. A criterion for appropriate input voltage was based on 0.5°C RMS MS difference from measured temperature at the center of the gel slab. As a result of the simulations with a range of voltage, 9,500 V on the upper plate was selected for all further simulations. The simulated temperature distribution of the gel slab was in good agreement with the experimental one both in pattern and absolute temperatures (Fig. 3b). Marshall and Metaxas (1998) used a similar approach to estimate an appropriate voltage in modeling RF-assisted heating of particulates. Upon validation of the FEMLAB model for the slab, the model was further used to characterize the RF heating pattern in different geometries as influenced by the dielectric properties of medium and model fruit as well as locations of the model fruit relative to electrodes.

3.2. Electric field pattern in various object geometry

As the gel slab overheats at edges and corners of the slab because of fringe field effects, several simulation models were developed to better understand the effect of object shape on electric field distribution in the object. Fig. 4 shows the E-field distribution on the vertical cross section of sphere (80 mm dia.), cylinder (80 mm dia. × 80 mm height), and cube (80×80×80 mm³) immersed in water and situated in the center of the container with 25 mm above the bottom of the container. Under similar conditions, very high electric field concentrations occurred at the corners and edges of the cube, edges and the middle of the cylinder, whereas maximum E-field was at the bottom of the sphere. Interestingly, lowest electric field spot was always at the top central portion of all the objects. The spherical object is expected to heat more evenly as the electric field variation was only 37 % in sphere compared to 151 % and 173% variations in the cylindrical and the cubical objects respectively. The problem of uneven heating can be further aggravated by the dielectric properties’ dependency on temperature which may result in thermal runaway heating. Therefore, it is very important to devise means to eliminate these hot spots in the object.
3.3. RF heating of model fruit in air

We learned from the results presented in the previous section that spherical objects should have higher electric field concentration at the fruit section close to the electrode. Fig. 5a corroborates the result as the maximum heating occurred at lower segment of the model fruit, placed at 25 mm above the ground electrode or 25 mm below the upper electrode in 160 mm of air gap between two RF electrodes. Experimental temperature profile over the fruit cross section corroborated the simulated temperature profile of the 14 min RF heating of the model fruit in 160 mm electrode gap (Fig. 5b).

When the model fruit was placed in the middle of 160 mm electrode gap, the fruit sections near to electrodes were heated at faster rate than the core of the fruit as shown in Fig. 5. The uniform heating over the entire fruit section can be expected if the gap between two electrodes is very large. But, maintaining large air gap is not practical for industrial applications because a large gap reduces the electric field strength results in slow heating rates. It can be inferred from the above results that uniform RF heating of the fresh fruits would not be possible. Hallman & Sharp (1994) also summarized the experience of many researchers and concluded that localized and uneven heating of fresh fruits is the major obstacle in using RF energy for fresh fruits.

3.4. RF heating pattern in fruit immersed in water

Immersion of the fresh fruits in water has been suggested by Ikediala et al. (2002) as a means to overcome the problems associated with RF heating of the fresh fruits, hence the simulation model was developed to study the effect of water on heating pattern. Fig. 6a

![Fig. 4. Simulated electric field (V/m) distribution in the vertical cross sections of sphere, cylinder and cube made of gellan gel (ε = 84 − j 220) placed in 195 mm RF electrodes and the objects were 25 mm above the bottom of the water (ε = 76 − j18 ) filled container (405 mm dia × 127 mm height).](image)

![Fig. 5. Simulated (a) and experimental (b) temperature distributions after 14 min of RF heating of the model fruit (Φ80 mm) placed at 25 mm above the ground electrode, middle, 25 mm below the top electrode in a 160 mm electrode gap.](image)
shows that immersion of the center-placed model fruit in water slightly shifted the hot spot toward the core of the model fruit. Moreover, presence of water radically enhanced power coupling as it took only half time (7 min) what required for RF heating of the fruits in air to reach ~50°C temperature. The increased heating rate was attributed by lossy nature of the water which offers less resistance to electric field compared to the air.

When the model fruit moved closer to the container wall from the center of the container, maximum heating occurred at the fruit section segment close to the container wall as shown in Fig 6b. Since water provides the path of least resistance, electric fields converge at the top of the water-wall interface. The very high electric field concentration at edges passes through the model fruit having large loss factor.

The temperature profile along the line A-A as shown in Fig. 6c corroborated that simulation and experimental results agreed. Root mean square (RMS) of temperature difference between the experimental and simulated distributions at line A-A was calculated and the RMS values were 0.98°C and 0.74°C for corner and centrally placed model fruits, respectively. It was also corroborated by the experimental and simulated time-temperature history (RMS value was 1.1°C) at the core of the centrally placed gel ball as shown in Fig. 7.

### 3.5. Effect of fruit rotation

It was clear from above discussion that immersion technique alone did not completely eliminate the uneven RF heating of large fruits such as apple and orange. To overcome this problem, Birla et al. (2004) suggested to move and rotate fruits in water during RF heating. Fig. 8 shows the effect of rotation and movement of a fruit-water mix on the heating pattern of model fruit. The concentric heating pattern over the fruit cross section was observed in both experimental and simulation

![Figure 6. Experimental (a) and simulated (b) temperature distributions inside model fruit (Φ80 mm) and at horizontal radial line A-A (c) after 7 min of RF heating in a 195 mm electrode gap.](image-url)
temperature profiles. However, there was a discrepancy in simulated and experimental temperatures over the cross section of the model fruit. In the simulation, spherical ball was kept in the center to avail the benefits of axial symmetry, whereas in the experiment model fruit moved along the peripheral of the rectangular fruit mover. This should explain the observed discrepancy in the final temperature. It is quite logical that the rotation of the spherical objects during RF heating evened out the effect of proximity of the sphere to electrodes and container walls. It should be noted that rotation and movement of fruit did not guarantee uniform heating all over the cross section of the homogenous fruit mass because of the difference in dielectric properties of the medium and fruit. A saline water immersion technique or differential startup temperature should be adopted along with fruit movement and rotation during RF heating to minimize differential heating between fruit and water.

3.6. Effect of dielectric properties of medium and the model fruit

Upon validating the simulation model, the developed model was used to study the effects of dielectric properties of the model fruit and medium. One model fruit was placed in the center of the water filled container at 25 mm above the bottom of the electrode. The value of loss factor was varied intentionally by changing the constant (C) in the regression Eqs. 8 and 9 such that the initial value at 20°C were 20, 40, 80 for water and 80, 160, 240, 300 for the model fruit. It was assumed here that in 20 to 50 °C temperature range, the loss factor of the water or the gel varies linearly and the effect of the temperature remains unchanged. It was also assumed that the regression Eqs. for the dielectric constant remains unchanged irrespective of the intentional increase in the loss factor value. The regression Eqs. for the dielectric constant of the model fruit and the water used in this study are listed in Table 1.

\[
\varepsilon''_w = 0.33T + C \tag{8}
\]

\[
\varepsilon''_m = 4.36T + C \tag{9}
\]

3.6.1. Loss factor ($\varepsilon''$) of water

Fig. 9a shows the influence of increasing loss factor of medium on the temperature profile at the vertical line drawn through the center of the model fruit placed in centre of 25 mm above the bottom of the container. An increase in the water loss factor (by increasing ‘C’ value in Eq. 8, and using all other properties of the fruit and the water as listed in Table 1) reduced the peak temperature inside the fruit, as saline water was likely to provide a conductive media for EM energy to pass through water as the path of least resistance.

Fig. 9b summarizes the results of a series of simulations with water loss factor initial values at 20°C ranging from 20 to 180. Increasing water loss factor increased the temperature but the temperature started to decline after a certain water $\varepsilon''$ value (~80). Altering the dielectric properties of water by addition of salt minimizes the differential heating between the fruit (oranges and cherries) and water (Birla et al., 2004; Ikediala et al., 2002). Differential heating between fruit and water can be minimized by appropriately matching dielectric properties of the water with that of fruit.

![Fig. 7. Simulated and experimental time-temperature histories at the core of the centrally placed model fruit (Φ80 mm) in a water filled container after 7 min of RF heating in a 195 mm electrode gap.](image)

![Fig. 8. Experimental and simulated temperature distributions of model fruits (Φ80 mm) rotating and moving in the fruit mover (Birla et al., 2006) after 6 min of RF heating in a 195 mm electrode gap.](image)
3.6.2. Loss factor of fruit

Fig. 10 shows the simulated effect of the increasing fruit loss factor on the temperature profile at the vertical line drawn through the center of the model fruit placed in centre of 25 mm above the bottom of the container (7 min of RF heating with a 195 mm electrodes gap). The simulation results showed that increasing the $\varepsilon''$ value of the model fruit (by increasing the constant ‘C’ in the Eq. 9, and using all other properties of the tap water and the model fruit as listed in Table 1) resulted in increasing heating rate only to a certain value ($\varepsilon'' = 180$), which was contrary to a general belief that increasing loss factor constantly increases heating rate inside the fruit. But beyond that value increasing loss factor diminished the heating rate inside the fruit. This can be explained by the fact that power density is proportional directly not only to the loss factor, but also to the square of electric field intensity. Increasing loss factor decreases the electric field intensity hence maximum heating rate occurs at particular loss factor value. An expression for limiting $\varepsilon''$ value as a function of the dielectric properties of a surrounding medium was derived by Birla (2006). Now with this simulation model and analytical expression a faster heating of apple ($\varepsilon'' = 120$) in comparison with oranges ($\varepsilon'' = 220$) can be explained.

3.7. Limitation of the model

The computer model developed in this study was able to predict the transient temperature profiles using physical and temperature-dependent dielectric properties, which are needed to design a RF heating process for disinfection of fruits. Having predicted the time-temperature profile in large fruit, one can decide the adequacy of the heating process based on information about the thermal death kinetics of insects/pests and quality degradation kinetics. In the present model, we assumed constant voltage over the entire upper electrode. However, in practice, this would not be true due to the effects of distributed system parameters and high frequency. To take these factors into account for the proper design of a whole RF system, an EM wave equation should be solved for domains including RF generator, tank circuit and applicator.

4. Conclusions

A computer model based on finite element method was developed to the quasi-static electric fields in an RF heating system. The coupled Maxwell’s EM equations and Navier-Stokes equations were solved for a 3-D model using FEMLAB software. The simulation results were validated with experimental temperature profiles of gellan gel model fruit. The experimental and simulation results were found to be in agreement, as the model
was able to predict transient temperature profiles using physical and temperature-dependent dielectric properties. The validated computer model was further used to study the effects of dielectric properties of the medium and fruits, and effect of fruits rotation. Non-uniform heating is attributed to shape, dielectric properties, and relative fruit position in the container. The predicted temperature profile may be useful for designing an RF heating process in disinfection of fruits. The developed model can be used as a tool to study the heating pattern of various fruits as influenced by dielectric properties of peel, pulp, and inner core.

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