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# Medium-sized double magic metal clusters: $\text{Al@Cu}_{54}^-$ and $\text{Al@Ag}_{54}^-$

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Medium-sized double magic metal clusters,  $\text{Al@Ag}_{54}^-$  and  $\text{Al@Cu}_{54}^-$ , are predicted based on unbiased global search and density functional calculation. Both bimetallic core-shell clusters have icosahedral symmetry, and they are much lower in energies than all other low-lying isomers. In contrast, the icosahedral cluster  $\text{Al@Au}_{54}^-$  is a high-energy isomer. Both  $\text{Al@Ag}_{54}^-$  and  $\text{Al@Cu}_{54}^-$  exhibit appreciable gaps between the highest occupied molecular orbital and the lowest unoccupied molecular orbital, and strong spherical aromaticity, which provide two additional evidences for the likelihood of their high stability. The simulated anion photoelectron spectra and optical absorption spectra are readily compared with future experiments. © 2008 American Institute of Physics.  
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## I. INTRODUCTION

Magic-number clusters are an important topic in cluster science. According to the spherical jellium model, near-spherical metal clusters with specific number of valence electrons, e.g., 8, 18, 20, 34, 40, 58,  $\dots$ , are very robust because of their closed electronic shells.<sup>1,2</sup> These clusters are called (electronic) magic-number clusters due to their high stability and much higher abundance in mass spectrometric studies than other sizes of metal clusters. Another hallmark of electronic magic-number clusters is that they typically have much larger energy gap between the highest occupied molecular orbital and the lowest unoccupied molecular orbital (HOMO-LUMO gap) than their neighboring-size clusters.

Besides the electronic factor, atomic packing and geometric symmetry is another important characteristic of highly stable clusters. Indeed, many known highly stable clusters exhibit very high group symmetry such as the icosahedral symmetry. Thus, highly symmetric clusters may be viewed as (geometric) magic-number clusters. The importance of both electronic and geometric factors to the stability of magic-number clusters has already been noted in previous studies.<sup>3,4</sup> Chen *et al.* recently proposed a viable way to design highly stable clusters based on both geometric factor (compact packing) and electronic factor (spherical aromaticity).<sup>5</sup> Clusters with both closed electronic shell and the icosahedral symmetry have been coined as the “double magic” clusters.<sup>6,7</sup> Known examples of double magic metal clusters are the icosahedral cluster  $I_h\text{-Al}_{13}^-$  (Ref. 8–13) and the core-shell clusters  $I_h\text{-W@Au}_{12}$  and  $I_h\text{-Mo@Au}_{12}$ .<sup>14,15</sup> These small-sized double magic clusters possess 40 and 18 valence electrons, respectively. Clusters with both closed electronic shell and the tetrahedral symmetry may be considered as the “nearly double magic” clusters.<sup>16</sup> For example, the tetrahedral cluster  $T_d\text{-Au}_{20}$  (Ref. 17) possesses 20 valence electrons which form four-center two-electron bond in each of ten tetrahedral cavities. Other examples of “near-double-magic” metal clusters are the tetrahedral cluster

$\text{Sc@Cu}_{16}^+$  (Ref. 18) and the core-shell cluster  $C_s\text{-Cu@Au}_{16}^-$ ,<sup>19</sup> both possessing 18 valence electrons. Note that there exists a second type of double magic metal clusters, namely, group-14 Zintl ion metal clusters such as the icosahedral core-shell clusters  $I_h\text{-Al@Pb}_{12}^+$  and  $I_h\text{-Pt@Pb}_{12}^{2-}$  (Refs. 20 and 21) and the stannaspherene  $I_h\text{-Sn}_{12}^{2-}$  and plumbaspherene  $I_h\text{-Pb}_{12}^{2-}$ .<sup>22,23</sup> These small-sized metal clusters possess 50 apparent valence electrons but their high stability can be rationalized by their 5*p* and 6*p* valence band contributions which contain a total of 26 electrons<sup>24</sup> and obey the Wade skeletal electron rule.<sup>25,26</sup> The third type of double magic metal clusters was recently proposed by Dognon *et al.*,<sup>24</sup> namely, the 32-electron bimetallic core-shell clusters such as  $\text{Pu@Pb}_{12}$  with the icosahedral symmetry.

Noble-metal clusters have been extensively studied for many years. Early experiments have established that  $\text{Au}_n$ ,  $\text{Ag}_n$ , and  $\text{Cu}_n$  clusters ( $n=2, 8, 20, 34, 58, \dots$ ) are all electronic magic-number clusters.<sup>27,28</sup> In particular, gold clusters in the size range of  $\text{Au}_3\text{--Au}_{34}$  have received considerable attention.<sup>27–39</sup> The near-double-magic cluster  $T_d\text{-Au}_{20}$  has a very large HOMO-LUMO gap ( $\Delta_{\text{HL}}=1.77$  eV).<sup>17</sup> For medium-sized gold clusters (with number of atoms  $>20$ ),  $\text{Au}_{34}$  is an electronic magic-number cluster with a reasonably large HOMO-LUMO gap ( $\Delta_{\text{HL}}=0.94$  eV) but low symmetry.<sup>37,38</sup> Copper and silver clusters have also been studied extensively.<sup>40–48</sup> In particular, the medium-sized clusters  $\text{Cu}_{55}$  and  $\text{Ag}_{55}$  are geometric magic-number clusters with the icosahedral symmetry, whereas the icosahedral cluster  $\text{Au}_{55}$  is not the global minimum.<sup>46–49</sup>

The bimetallic double magic cluster,  $I_h\text{-W@Au}_{12}$ , was first predicted by Pyykkö and Runeberg<sup>14</sup> and later experimentally produced and detected by Li *et al.*<sup>15</sup> This core-shell metal cluster has a very large HOMO-LUMO gap  $\Delta_{\text{HL}}=1.68$  eV. The central “impurity” atom can be used to fine-tuning electronic properties of core-shell clusters.<sup>50</sup> Similar double magic clusters  $I_h\text{-M@Au}_{12}^-$  ( $M=\text{V, Nb, Ta}$ ) were also observed experimentally.<sup>51</sup> A number of magic-number metal clusters with 18 valence electrons have been proposed

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and studied theoretically.<sup>52–55</sup> These core-shell metal clusters are predicted to be stable due to the “spherical aromaticity” of the shell as well as the geometric compatibility between the shell and the core.<sup>5</sup>

In this paper, we report a theoretical prediction of two medium-sized double magic metal clusters of the first type, namely, the bimetallic core-shell clusters  $I_h\text{-Al}@M_{54}^-$  ( $M=\text{Cu, Ag}$ ). In addition, we show that icosahedral cluster  $I_h\text{-Al}@Au_{54}^-$  is a high-energy isomer. To our knowledge, double magic metal clusters with the size greater than 20 or with valence electrons greater than 50 have not been reported previously.

## II. MODELS AND METHODS

The initial structure of  $I_h\text{-Al}@M_{54}^-$  ( $M=\text{Cu, Ag, Au}$ ) was constructed based on the icosahedral clusters  $I_h\text{-}M_{55}$  ( $M=\text{Cu, Ag, Au}$ ), all having 42 atoms in the outer shell, 12 atoms in the inner shell, and 1 atom at the center. We simply replaced the central metal atom  $M$  by an Al atom and added a negative charge to the entire cluster so that the metal clusters possess 58 valence electrons. We then applied the basin-hopping global minimization method<sup>56</sup> coupled with density functional theory (DFT) (Ref. 57) to search for the lowest-energy isomer. Specifically, we used the generalized gradient approximation in the form of Perdew-Burke-Ernzerhof (PBE) functional<sup>58</sup> together with the double numerical plus  $d$  (DND) basis set with the semicore pseudopotential, which are implemented in the DMOL3 software package.<sup>59,60</sup> Harmonic frequency calculations were performed for the lowest-energy structure. The reliability of the combination of basin-hopping Monte Carlo and DFT method has been verified for the search of several highly stable metal clusters by us.<sup>33,36,38</sup>

Nucleus-independent chemical shift (NICS) (Refs. 61 and 62) values were computed using BP86/LAN2DZ (Refs. 63 and 64) method implemented in GAUSSIAN03 package.<sup>65</sup> The same level of theory was used to compute anion photoelectron spectra. The optical adsorption spectra were computed using the PBE/TZP (triple zeta polarization function) method implemented in the ADF2007 package.<sup>66</sup>

## III. RESULTS AND DISCUSSION

The unbiased basin-hopping run generated more than 50 isomers for each of  $\text{Al}@M_{54}^-$  ( $M=\text{Cu, Ag}$ ) cluster and the top-5 lowest-energy isomers of  $\text{Al}@M_{54}^-$  ( $M=\text{Cu, Ag}$ ) are shown in Table I. The search suggests that the icosahedral isomer  $I_h\text{-Al}@M_{54}^-$  ( $M=\text{Cu, Ag}$ ) is likely the lowest-energy structure (see 1a and 1b in Fig. 1). The second lowest-energy isomers are also displayed in Fig. 1 (2a and 2b). As shown in Table I, the second lowest-energy isomer has appreciably higher energy ( $>0.7$  eV) than the lowest-energy isomer. This large energy difference further supports that the icosahedral structure  $I_h\text{-Al}@M_{54}^-$  ( $M=\text{Cu, Ag}$ ) is most likely the *global minimum*. It is worth noting that all the top-4 lowest-lying isomers exhibit appreciable HOMO-LUMO gaps  $\Delta_{\text{HL}} \sim 0.40\text{--}0.57$  eV. However, for  $\text{Al}@Au_{54}^-$ , the icosahedral structure 1c is a high-energy isomer. In fact, starting from 1c, the basin-hopping run showed that within only four Monte Carlo steps the structure of 1c was changed to 2c which has

TABLE I. (Color online) Outer shells of top-5 lowest-energy isomers of  $\text{Al}@Cu_{54}^-$  and  $\text{Al}@Ag_{54}^-$ . The core of all these isomers is  $I_h\text{-Al}@M_{12}$  ( $M=\text{Cu, Ag}$ ). The energy rankings are based on PBE/DND calculation. The HOMO-LUMO gap  $\Delta_{\text{HL}}$  (in unit of eV) are given in the parenthesis.  $\Delta E$  is the relative energy with respect to the lowest-energy isomer.

No.	$\text{Al}@Cu_{54}^-$	$\Delta E(\text{eV})$ ( $\Delta_{\text{HL}}$ )	$\text{Al}@Ag_{54}^-$	$\Delta E(\text{eV})$ ( $\Delta_{\text{HL}}$ )
1		0.00 (0.43)		0.00 (0.40)
2		0.69 (0.57)		0.65 (0.53)
3		0.92 (0.54)		0.66 (0.54)
4		1.81 (0.52)		0.72 (0.53)
5		1.93 (0.56)		0.98 (0.17)

a much lower energy (by 1.35 eV) than 1c. Interestingly, the core of 2c has only 11 atoms, two less than the 13-atom core of 1c. We note that the distinct structural difference between  $\text{Al}@Au_{54}^-$  and  $I_h\text{-Al}@M_{54}^-$  ( $M=\text{Cu, Ag}$ ) is most likely due to the strong relativistic effect entailed in the gold cluster, as thoroughly discussed in Ref. 47.

### A. Structural, electronic, and vibrational properties

Structural, electronic, and vibrational properties of  $I_h\text{-Al}@Ag_{54}^-$  and  $I_h\text{-Al}@Cu_{54}^-$  are collected in Table II. Also included in Table II are properties of the known double magic cluster  $I_h\text{-W}@Au_{12}$  for comparison. As expected, the medium-sized clusters  $I_h\text{-Al}@Ag_{54}^-$  and  $I_h\text{-Al}@Cu_{54}^-$  have much smaller HOMO-LUMO gaps ( $\Delta_{\text{HL}} \sim 0.4$  eV) than that (1.80 eV) for small-sized cluster  $I_h\text{-W}@Au_{12}$ . We note that  $\Delta_{\text{HL}} \sim 0.4$  eV is close to the measured  $\Delta_{\text{HL}} \sim 0.65$  eV of  $Au_{58}$ . The latter is an electronic magic-number cluster with 58 valence electrons.<sup>47</sup>  $I_h\text{-Al}@Au_{54}^-$  is a high-energy isomer but still a local minimum. The harmonic frequency calculation shows that  $I_h\text{-Al}@Au_{54}^-$  has no imaginary frequency. The lowest vibrational frequency of  $I_h\text{-Al}@Au_{54}^-$  is  $19.3\text{ cm}^{-1}$ , much less than that of  $I_h\text{-Al}@Ag_{54}^-$  ( $33.9\text{ cm}^{-1}$ ) and

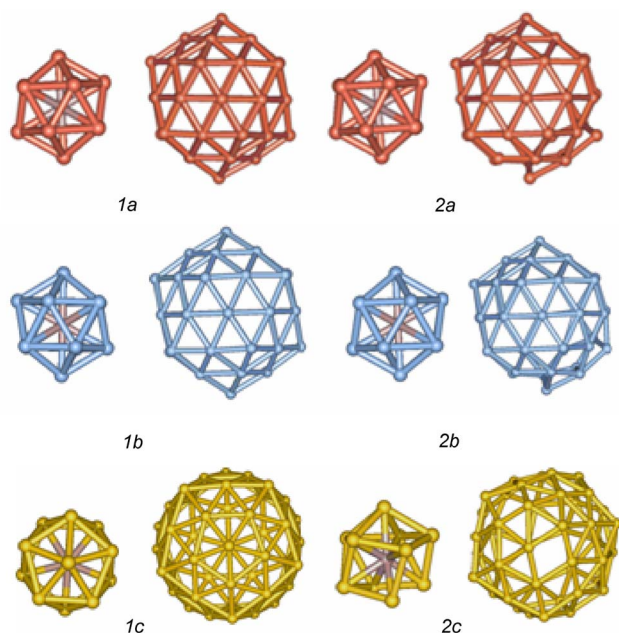


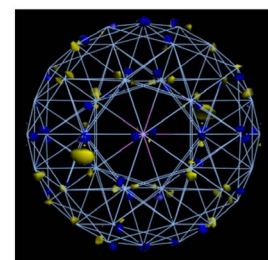
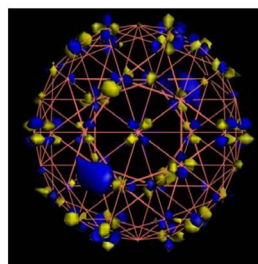
FIG. 1. (Color online) The predicted lowest-energy structure 1 and the second lowest-energy structure 2 of a Al@Cu<sub>54</sub><sup>-</sup> and b Al@Ag<sub>54</sub><sup>-</sup>. The inner 13-atom core and the outer 42-atom shell are plotted separately for ease of view. The Al atom (purple) is at the center of the core. 2c is an isomer of Al@Au<sub>54</sub><sup>-</sup> with lower energy by 1.35 eV than the icosahedral isomer 1c. The core of 2c has 11 atoms.

$I_h$ -Al@Cu<sub>54</sub><sup>-</sup> (55.9 cm<sup>-1</sup>) (Table II), indicating that  $I_h$ -Al@Ag<sub>54</sub><sup>-</sup> and  $I_h$ -Al@Cu<sub>54</sub><sup>-</sup> are less fluxional than  $I_h$ -Al@Au<sub>54</sub><sup>-</sup>. This result offers an explanation that  $I_h$ -Al@Au<sub>54</sub><sup>-</sup> may be too fluxional to maintain its icosahedral structure stable.

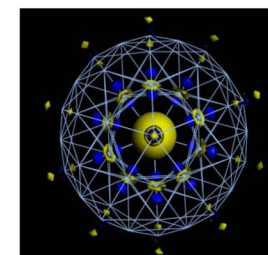
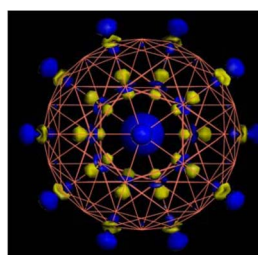
TABLE II. Computed properties of  $I_h$ -Al@Ag<sub>54</sub><sup>-</sup> and  $I_h$ -Al@Cu<sub>54</sub><sup>-</sup> vs  $I_h$ -W@Au<sub>12</sub>.

	$I_h$ -Al@Ag <sub>54</sub> <sup>-</sup>	$I_h$ -Al@Cu <sub>54</sub> <sup>-</sup>	$I_h$ -W@Au <sub>12</sub>
Diameter (nm)	1.1	1.0	0.6
$\Delta_{HL}$ (eV)	0.40	0.43	1.80
Frontier orbital	$(h_g)^{10}(g_g)^8(a_g)^0$	$(h_g)^{10}(g_g)^8(a_g)^0$	$(t_{2g})^6(h_g)^{10}(h_g)^0$
Outer-shell longest bond (Å)	3.014	2.620	...
Outer-shell shortest bond (Å)	2.922	2.546	2.918
Inner shell (Å)	2.910	2.577	...
Al-M or W-Au (Å)	2.766	2.450	2.775
Lowest vibrational frequency (cm <sup>-1</sup> )	33.9	55.9	30.1
Highest vibrational frequency (cm <sup>-1</sup> )	270.3	398.1	196.9
NICS at center excluding Al or W atom	-102.6	-113.7	-56.3

## HOMO



## LUMO



$I_h$ -Al@Cu<sub>54</sub><sup>-</sup>

$I_h$ -Al@Ag<sub>54</sub><sup>-</sup>

FIG. 2. (Color online) HOMO and LUMO diagrams of Al@Cu<sub>54</sub><sup>-</sup> and Al@Ag<sub>54</sub><sup>-</sup>. The symmetries of HOMO and LUMO are  $G_g$  and  $A_g$ , respectively.

dral structure stable. Note that  $I_h$ -W@Au<sub>12</sub>'s lowest vibrational frequency is 30.1 cm<sup>-1</sup> which is also higher than that of  $I_h$ -Al@Au<sub>54</sub><sup>-</sup>.

## B. Spherical aromaticity

Spherical aromaticity is likely an important factor to the stability of  $I_h$ -Al@Ag<sub>54</sub><sup>-</sup> and  $I_h$ -Al@Cu<sub>54</sub><sup>-</sup>. As shown in Table II, the NICS values at the center of 58-electron shells,  $I_h$ -Ag<sub>54</sub><sup>4+</sup> and  $I_h$ -Cu<sub>54</sub><sup>4+</sup>, are -102.6 and -113.7, respectively. The two NICS values are much more negative than that at the center of the 18-electron shell of  $I_h$ -Au<sub>12</sub><sup>6-</sup> (-56.3), 32-electron shell  $I_h$ -Au<sub>32</sub> (BP86: -80.7), and  $T_d$ -Au<sub>20</sub> (BP86: -36), suggesting the stronger aromaticity for the medium-sized double magic clusters, even though the 58 valence electrons do not obey the  $2(N+1)^2$  rule for spherical aromaticity. Note that the NICS value of  $I_h$ -Au<sub>54</sub><sup>4+</sup> is -111.9, very close to that of  $I_h$ -Cu<sub>54</sub><sup>4+</sup>, although it is not the global minimum due to strong relativistic effects.

## C. HOMO and LUMO diagrams

HOMO and LUMO diagrams of  $I_h$ -Al@Ag<sub>54</sub><sup>-</sup> and  $I_h$ -Al@Cu<sub>54</sub><sup>-</sup> are shown in Fig. 2. It can be seen that the HOMOs are quite similar, contributed mainly by the  $p$  orbitals of the noble-metal atoms but little from the central Al atom. The LUMOs are mainly contributed by the  $s$  and  $p$  orbitals of noble-metal atoms as well as the  $3s$  orbital of Al atom. Meanwhile, the population analysis shows that the Mülliken charges of central Al atom for  $I_h$ -Al@Ag<sub>54</sub><sup>-</sup> and  $I_h$ -Al@Cu<sub>54</sub><sup>-</sup> are 1.82e and 2.99e, respectively. More interestingly, the negative charges for middle-shell and outer-shell present different trend for the two clusters. For  $I_h$ -Al@Ag<sub>54</sub><sup>-</sup>, each Ag atom in the middle shell possesses a

charge of  $-0.048e$ , slightly less negative than that of the outer shell (vertex:  $-0.051e$ ; edge:  $-0.055e$ ). While for  $I_h\text{-Al@Cu}_{54}^-$ , each Cu atom in the middle shell possesses  $-0.144e$ , more negative than that of the outer shell (vertex:  $-0.051e$ ; edge:  $-0.055e$ ). Note also that for both clusters, the Mulliken charges of the outer surface are almost same ( $I_h\text{-Al@Ag}_{54}^-$ :  $-2.24e$ ;  $I_h\text{-Al@Cu}_{54}^-$ :  $-2.27e$ ), suggesting that both clusters have very similar electronic properties.

#### D. Simulated anion photoelectron spectra

Simulated anion photoelectron spectra of  $I_h\text{-Al@Ag}_{54}^-$  and  $I_h\text{-Al@Cu}_{54}^-$ , together with the simulated spectrum of  $I_h\text{-Ag}_{55}^-$  are shown in Fig. 3. The  $I_h\text{-Ag}_{55}^-$  is known to be the global minimum of  $\text{Ag}_{55}^-$  and the experimental photoelectron spectrum of  $I_h\text{-Ag}_{55}^-$  has been reported previously.<sup>46</sup> In the photoelectron spectrum, the location of the first peak characterizes the vertical detachment energies of the anion, which are  $\sim 3.1$  eV for  $I_h\text{-Al@Ag}_{54}^-$  and  $I_h\text{-Ag}_{55}^-$  and  $\sim 3.0$  eV for  $I_h\text{-Al@Cu}_{54}^-$ . The gap between the first two peaks gives rise to the HOMO1-HOMO2 gaps of the clusters, which are 0.68 eV for  $I_h\text{-Al@Cu}_{54}^-$ , 0.70 eV for  $I_h\text{-Al@Ag}_{54}^-$ , and 0.62 eV for  $I_h\text{-Ag}_{55}^-$ . Additionally, the photoelectron spectrum of  $I_h\text{-Al@Ag}_{54}^-$  and  $I_h\text{-Ag}_{55}^-$  exhibits the same sequence of the group symmetry for the first five peaks, suggesting that the replacement of the central Ag atom by Al in  $I_h\text{-Ag}_{55}^-$  does not change the order of the occupied orbitals, even though the orbital energies are changed [as can be seen in Fig. 3 that the third ( $T_{1u}$ ) and fourth ( $G_u$ ) peaks nearly overlap with each other for  $I_h\text{-Al@Ag}_{54}^-$ ].

#### E. Computed optical absorption spectra

Computed optical absorption spectra (less than 4 eV) of  $I_h\text{-Al@Ag}_{54}^-$  and  $I_h\text{-Al@Cu}_{54}^-$  are shown in Fig. 4. Unlike the photoelectron spectra, the two optical spectra are quite different. For  $\text{Al@Ag}_{54}^-$ , there are two weak peaks at 1.70 and 1.86 eV, one modest peak at 2.33 eV, and a strong peak at 2.65 eV. There are no apparent absorption peaks between 3 and 4 eV. However, for  $\text{Al@Cu}_{54}^-$ , there is one modest peak at  $\sim 1.52$  eV, one very strong peak at  $\sim 2.0$  eV, and one modest peak between 3 and 4 eV.

Finally, we note that besides  $I_h\text{-Al@Ag}_{54}^-$  and  $I_h\text{-Al@Cu}_{54}^-$  we have performed structural optimization (using PBE/DND level of theory) for a series of bimetallic core-shell clusters  $I_h\text{-M@N}_{54}^\delta$  ( $M=\text{Pb, Sn, Zr, Ti, Sc, V, Ce, Th}$ ;  $N=\text{Ag, Cu}$ ;  $\delta=-1, 0, +1$ ), all having 58 valence electrons. It is found that only  $I_h\text{-Ti@Ag}_{54}$ ,  $I_h\text{-Zr@Ag}_{54}$ ,  $I_h\text{-Ti@Cu}_{54}$ ,  $I_h\text{-Zr@Cu}_{54}$ ,  $I_h\text{-Th@Ag}_{54}$ ,  $I_h\text{-Th@Cu}_{54}$ ,  $I_h\text{-Sc@Ag}_{54}^-$ , and  $I_h\text{-Sc@Cu}_{54}^-$  can be optimized successfully, and these clusters show smaller HOMO-LUMO gaps ( $\Delta_{\text{HL}} \sim 0.06\text{--}0.32$  eV) than the isoelectronic clusters  $I_h\text{-Al@Ag}_{54}^-$  and  $I_h\text{-Al@Cu}_{54}^-$ . The latter two have larger HOMO-LUMO gaps because their 58 valence electrons are composed of doubly occupied  $s$  and  $p$  shells and singly occupied  $d$ ,  $f$ , and  $g$  shells. Thus, the inner atom is more likely to be a simple  $sp$  atom (e.g., Al) which can strongly attract  $s$  and  $p$  electrons. Unbiased searches are needed to further confirm that these icosahedral core-shell clusters are the global

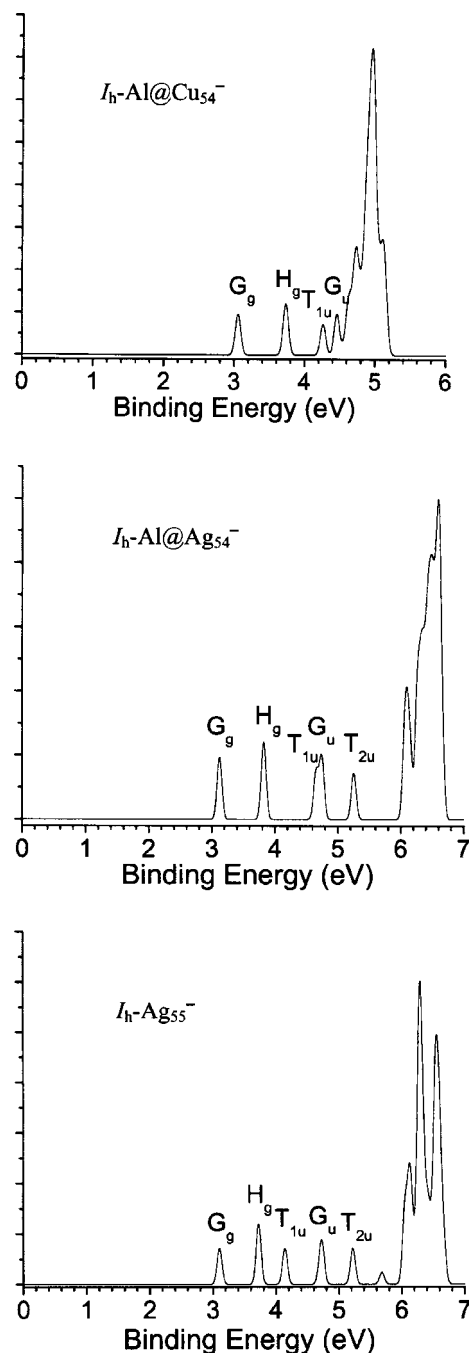


FIG. 3. Simulated anion photoelectron spectra of  $I_h\text{-Al@Cu}_{54}^-$ ,  $I_h\text{-Al@Ag}_{54}^-$ , and  $I_h\text{-Ag}_{55}^-$ , based on BP86/LANL2DZ level of theory implemented in GAUSSIAN03 program (Ref. 65).

minima, i.e., double magic metal clusters. Nevertheless, the preliminary result suggests that the actinides seem capable of maintaining the clusters at the highest icosahedral symmetry, but the lanthanides seem not.<sup>24</sup>

#### IV. CONCLUSION

In conclusion, we present theoretical evidence of two medium-sized double magic metal clusters  $I_h\text{-Al@Ag}_{54}^-$  and  $I_h\text{-Al@Cu}_{54}^-$ . The unbiased search shows that the icosahedral isomer is markedly lower in energy than other low-lying isomers. Both  $I_h\text{-Al@Ag}_{54}^-$  and  $I_h\text{-Al@Cu}_{54}^-$  have appreciable HOMO-LUMO gaps. Moreover, the isoelectronic

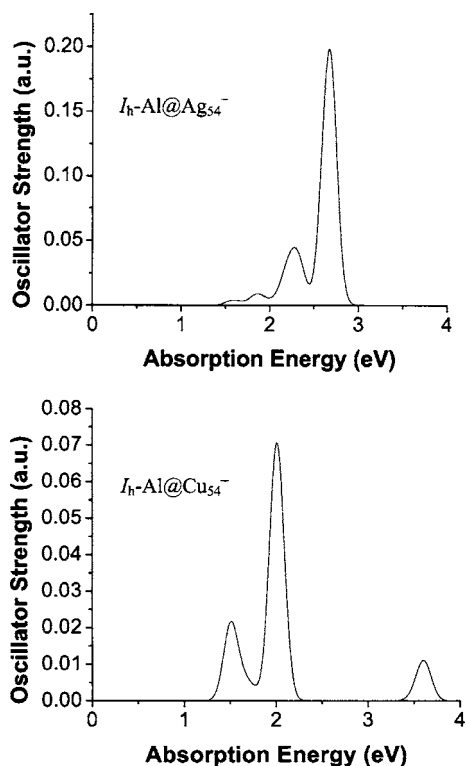


FIG. 4. Simulated optical absorption spectra of Al@Ag<sub>54</sub><sup>-</sup> and Al@Cu<sub>54</sub><sup>-</sup> based on the time-dependent DFT method at PBE/TZP level of theory implemented in the ADF program (Ref. 66).

outer shells  $I_h$ -Ag<sub>54</sub><sup>4-</sup> and  $I_h$ -Cu<sub>54</sub><sup>4-</sup> exhibit very large negative NICS values at their center, significantly more negative than other known double magic metal clusters. This result suggests that both  $I_h$ -Al@Ag<sub>54</sub><sup>-</sup> and  $I_h$ -Al@Cu<sub>54</sub><sup>-</sup> entail strong spherical aromaticity which may contribute to the stabilization of the icosahedral structure. The calculated anion photoelectron spectra as well as the optical absorption spectra of  $I_h$ -Al@Ag<sub>54</sub><sup>-</sup> and  $I_h$ -Al@Cu<sub>54</sub><sup>-</sup> are ready to compare with measured spectra in future. Finally, since  $I_h$ -Ag<sub>55</sub><sup>-</sup> and  $I_h$ -Cu<sub>55</sub><sup>-</sup> are stable at room temperature,<sup>46</sup> we expect that  $I_h$ -Al@Ag<sub>54</sub><sup>-</sup> and  $I_h$ -Al@Cu<sub>54</sub><sup>-</sup> are likely to be stable at the room temperature too. At elevated temperatures, however, isomers with lower symmetry would be more stable due to entropy effect. It will be interesting to find out the temperature at which the structural transition occurs. This study will be carried out in future.

## ACKNOWLEDGMENTS

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