2012

Measurement of the Ultrafast Temporal Response of a Plasmonic Antenna

Maria Becker
University of Nebraska - Lincoln, mgbecker3@gmail.com

Wayne Cheng-Wei Huang
University of Nebraska-Lincoln, wayneh@huskers.unl.edu

Herman Batelaan
University of Nebraska - Lincoln, hbatelaan@unl.edu

Elizabeth J. Smythe

Federico Capasso
Harvard University, capasso@seas.harvard.edu

Follow this and additional works at: https://digitalcommons.unl.edu/physicsfacpub

Part of the Physics Commons

Becker, Maria; Huang, Wayne Cheng-Wei; Batelaan, Herman; Smythe, Elizabeth J.; and Capasso, Federico, "Measurement of the Ultrafast Temporal Response of a Plasmonic Antenna" (2012). Faculty Publications, Department of Physics and Astronomy. 110. https://digitalcommons.unl.edu/physicsfacpub/110

This Article is brought to you for free and open access by the Research Papers in Physics and Astronomy at DigitalCommons@University of Nebraska - Lincoln. It has been accepted for inclusion in Faculty Publications, Department of Physics and Astronomy by an authorized administrator of DigitalCommons@University of Nebraska - Lincoln.
Measurement of the Ultrafast Temporal Response of a Plasmonic Antenna

Maria Becker,* Wayne Cheng-Wei Huang,† and Herman Batelaan‡
Department of Physics and Astronomy, University of Nebraska-Lincoln, Lincoln, Nebraska 68588, USA

Elizabeth J. Smythe§ and Federico Capasso
Harvard School of Engineering and Applied Science, Harvard University, Cambridge, Massachusetts 02138, USA

We report a measurement on the temporal response of a plasmonic antenna at the femtosecond time scale. The antenna consists of a square array of nanometer-size gold rods. We find that the far-field dispersion of light reflected from the plasmonic antenna is less than that of a 1.2 mm thick glass slide. Assuming a simple oscillating dipole model this implies that the near-field of the antenna may be used as an electron switch that responds faster than 20 fs. Alternatively, ultrafast electron diffraction may be used to investigate the near-field dynamics of the plasmonic antenna.

Light incident on a metallic material is strongly attenuated at the surface. For metallic films of nanometer-scale thickness, all electrons in the material can interact with light at optical frequencies. When this thin metallic film is patterned with nanometer scale boundaries to form a metallic nanoparticle, electron resonance behavior can occur. This phenomenon can take place when the dominant spatial mode of the charge oscillation, i.e. the dipole oscillation mode, corresponds to the boundary shape of the particle [1]. Consequently, metallic nanoparticles that are made of the same material but have different shapes reflect different colors under white light illumination [2, 3].

As the size of metallic particles shrinks down to subwavelength dimensions, a strong near-field can arise at the poles of the oscillating dipole. The magnitude of this near field can exceed ten times that of the incident field which drives the particle, a phenomena known as near-field enhancement [4, 5]. The near-field enhancement can be even greater if several metallic particles are positioned close to each other [4]. Additionally, the resonant response of a particle assembly can be tuned through variation of the inter-particle distances [6, 7]. Thus, an array of metallic nanoparticles can be used as an antenna, known as a plasmonic antenna [6]. The periodic structure and strong near-field enhancement of a plasmonic antenna may enable it to act as an electron diffraction grating. In order to utilize a plasmonic antenna as an electron diffraction/deflection device in the ultrafast regime, it is important to know how a plasmonic antenna responds to ultrafast light. In this Letter, we investigate the temporal response of a plasmonic antenna upon excitation by a femtosecond light pulse.

The temporal response of regular arrays of nanoparticles have been previously studied with methods sensitive to a variety of nanoparticle shapes [8–11]. The results from autocorrelation measurements of transmitted second- and third- harmonic generation signals indicated a fast relaxation time of 6 fs. The arrays used had periodicities and gapsizes of hundreds of nanometers. In this study, we use a reflective cross-correlation method, which does not require higher-harmonic generation. We also propose that this arrangement could be used for electron manipulation. Periodic structures with a period of about 100 nm and gaps of tens of nanometers are of particular interest for electron diffraction. The former provides practical diffraction angles, while the latter yields an intensity enhancement that may be suited for femtosecond electron switching. In view of our proposed application, the main objective of the present work is to establish that the enhanced near-fields in the small gaps between elements of the array studied do not significantly lengthen the temporal response of the reflected pulse.

The near-field of nanoscale structures has been probed time-independently with tip-enhanced electron emission microscopy [12] and time-dependently with photon-induced near-field microscopy [13]. As pointed out previously in [13], these techniques may also be used to investigate plasmonic antennas. The simple all-optical technique presented in this letter may be useful for selecting plasmonic antennas with an interesting response in the far-field that warrant a more involved study with pulsed electron techniques.

When excited by a light pulse, the plasmonic antenna generates a strong localized near-field. The far-field is radiated as a reflected pulse. For an array of dipole radiators, the relation between the antenna’s near- and far-fields can be deduced by inspecting the fields of each

* email: mgbecker3@gmail.com
† email: u910134@alumni.nthu.edu.tw
‡ email: lhatelaan@unl.edu
§ Current affiliation: Schlumberger-Doll Research in Cambridge, MA.
single oscillating dipole, \( p(t) = p_\omega e^{-i(\omega t + \phi_\omega)} \) [14],

\[
E_{\text{disp},\omega}(r, t) = \frac{1}{4\pi\epsilon_0} \left\{ \frac{k^2}{r} \left( \hat{r} \times \hat{z} \right) \times \hat{r} \\
+ \left( \frac{1}{r^3} - \frac{i k}{r^2} \right) \left[ 3(\hat{r} \cdot \hat{z}) \hat{r} - \hat{z} \right] \right\} p_\omega e^{-i(\omega t + \phi_\omega)} e^{ikr} \\
= E_{\text{far}}(r) p_\omega \omega^2 e^{i\phi_\omega} e^{-i\omega(t-r/c)} \\
+ E_{\text{mid}}(r) p_\omega \omega e^{i\phi_\omega} e^{-i\omega(t-r/c)} \\
+ E_{\text{near}}(r) p_\omega e^{i\phi_\omega} e^{-i\omega(t-r/c)},
\] (1)

where \( k = \omega/c, \omega \) is the frequency of the incident field, \( p_\omega \) is the dipole strength of the nanoparticle, \( \phi_\omega \) is the relative phase between the incident field and the induced dipole, \( \hat{z} \) is the unit vector along the direction of the dipole, and \( \hat{r} \) is the unit vector in the radial direction from the center of the dipole. \( E_{\text{far}}(r) \), \( E_{\text{mid}}(r) \), and \( E_{\text{near}}(r) \) are the spatial patterns of the far-, intermediate-, and near-fields, respectively. It is clear from Eq. (1) that the near- and far-fields have the same phase spectrum. Although the power spectra of the near- and far-fields differ by a factor of \( \omega^2 \), this difference is negligible for a plasmonic resonance curve that is about 100nm wide and centered at 800nm. Thus, the near- and far-fields of a nanoparticle can have an identical temporal pulse shape. The field reflected from the plasmonic antenna is the sum of the fields from many dipole radiators. When the plasmonic antenna is excited by a pulse incident at an angle, the far-fields from all the radiators have the same time delay in the direction specified by the law of reflection. In other directions, the far-fields interfere destructively. Thus, the far-field of the plasmonic antenna has the same temporal shape as the far-field of the individual radiators. The near-field close to the surface of the antenna has contributions from the incident field and the nearest radiators of the antenna. The field in the immediate vicinity of one particular radiator has also a contribution from a neighboring radiator. However that field has a time delay that depends on the distance between the next neighbor pairs. For an array of nanoparticles in an antenna, these distances can be in the subwavelength region. Then the contribution of a next neighbor radiator is delayed in time at most by their distance divided by the speed of light, and therefore, within the sub-cycle oscillation. This contribution would be in phase and not reduce the local field. The near-fields from more distant neighbors can have a significant time delay and may potentially cause the temporal duration of the near-field to broaden. However, the strength of the near-field of any radiator decays quickly over the distance of one wavelength, so the contribution from the far neighbors can be ignored. Consequently, the near-field of the plasmonic antenna has the same temporal shape as the near-field of the individual radiator. Assuming that the antenna is an array of independent dipole radiators, the temporal response of the antenna’s far-field equals the temporal response of the antenna’s near-field.

Information about the temporal width of the antenna’s far-field may be obtained from comparing the cross-correlation signal of the reflected pulse to the auto-correlation signal of the incident pulse. Comparison between the cross-correlation signal and the auto-correlation signal shows the contrast between the reflected pulse and the incident pulse, revealing any broadening that may have occurred. However, without knowing the specific shape of the pulse, its temporal width cannot be inferred directly from the cross-correlation signal. In this study, the cross-correlation signal is modeled by reconstructing the pulse fields from the experimentally obtained power spectrum and the theoretically computed phase spectrum. The phase spectrum of the reflected pulse is modeled by calculating the phase dispersion of a pulse which passes through a glass slide (fused silica) twice,

\[
\phi_{\text{disp}}(\omega) = \frac{2\omega d}{c} n(\omega) + \phi_{bs}(\omega),
\] (2)

where \( d \) is the thickness of the glass slide, \( \omega \) is the frequency, \( \phi_{bs}(\omega) \) is the phase dispersion due to the beamsplitter (BK7), and \( n(\omega) \) is the index of refraction of the glass (fused silica). The thickness of the glass, \( d \), is the only fitting parameter in this “glass-model” and is determined by comparing theoretical and experimental cross-correlation signals. To test the validity of this glass-model, experiments were performed in the antenna-configuration and the glass-configuration, shown in FIG. 1. After the pulse passed through a dispersion compensator [15], a beamsplitter splits the pulse to two
The plasmonic antenna used in the experiment was fabricated using electron-beam lithography as described in [6]. Images obtained with a scanning-electron microscope (see FIG. 2) indicate rod dimensions of 170nm × 80nm. The rod array has a period of 160nm in the direction perpendicular to the rod length and a spacing of 10 – 20nm along their length. The entire array has a size of 100µm × 100µm. Defects, such as missing rods and scratches, were present and are likely a result of the fabrication process and/or damage incurred during handling the device. The broadband reflectance spectrum was measured by focusing white light from a xenon lamp with a 50× microscope objective onto the plasmonic antenna (FIG. 2). Reflected light was collected into a fiber-coupled spectrometer and the resulting spectrum was normalized with the spectrum reflected by a silver mirror. The result is consistent with the relations between resonance and rod dimensions and spacing reported in [6]. Comparison of the input laser spectrum with the plasmonic antenna reflectance spectrum shows that the laser pulse contained frequencies appropriate for generating near-field enhancement in the antenna. The reflection of the array was also measured using a focused beam of linearly polarized light with a wavelength of 800nm (FIG. 2). The polarization-dependence of the reflected power is indicative of the resonance behavior of the antenna.

The experimental results are shown in FIG. 3. The autocorrelation signal is obtained by taking a Fourier transform of the measured power spectrum according to the Wiener-Khinchin theorem. The width of the autocorrelation signal is about 30fs. The cross-correlation widths measured in the antenna- and glass-configuration are about 40fs and 50fs, respectively. To interpret this result, the glass-model is used. Agreement between the model and the experimental data is found when the fitting parameter is set at 1.145nm, as shown in FIG. 4. This is indeed close to the measured thickness of the glass slide used in the experiment, which is 1.14nm. When this model is applied to the cross-correlation spectrum acquired in the antenna-configuration, the effect of reflection from the antenna can be modeled well by the glass-model when the fitting parameter is set to 0.25mm.

Use of this model enables the reconstruction of the reflected pulse by the nanorod array without inclusion of dispersion from other optical components in the system. The temporal width of the pulse reflected by the antenna is about 40 widths measured in the antenna- and glass-configuration.

In our experiment, we use a light pulse which has a full-width-half-maximum spectral bandwidth of ∆λ = 63nm and a central frequency of λc = 800nm. Assuming a transform-limited pulse shape, the full-width-half-maximum temporal width of the pulse field is ∆t ≃ 20fs.
amplified femtosecond laser oscillator operated at 800 nm, the dermotive potential drops to a small value. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800$, a significant enhancement of $\kappa = 800$. For a non-intensity of $I$, near-field, and $\Delta m$ estimated to be $\kappa$. Given an enhancement factor of $\kappa = 800
FIG. 5. Proposed femtosecond electron switch. A laser pulse affects the scattering of an electron pulse through the ponderomotive potential provided by the enhanced near field of a plasmonic antenna. The intersection of a laser pulse and electron pulse can be directed onto a planar structure. Following an idea of Zewail [19], the electron velocity and interaction angle can be chosen in such a way that the pulses remain synchronous as they sweep over the planar plasmonic structure. The consequence is that the temporal resolution of the switched electron beams (dashed blue lines) is set by the laser pulse duration.

justifies use of the glass-model to model the phase spectrum of the laser pulse reflected from the plasmonic antenna. According to this model, the temporal width of the antenna’s far field is shown to be close to that of the excitation pulse. Under the assumption that the plasmonic antenna is described by a collection of dipole radiators, the near field of the plasmonic antenna responds on the femtosecond scale. This ultrafast enhanced near field may affect the motion of incident electron beams. Femtosecond nanoscale pulsed electron sources have been developed [20, 21]. Electron pulse compression techniques applied to electron beams extracted from pulsed sources have been shown to deliver electron pulses of about 100 fs [22]. Electron pulse compression techniques have been proposed to deliver sub-femtosecond electron pulses [23–25]. A plasmonic antenna functioning as an electron switch may be an ideal detection method for these subfemtosecond electron pulses.

W. Huang, M. Becker, and H. Batelaan gratefully acknowledge funding from NSF Grant No. 0969506. The work of E. Smythe and F. Capasso was supported by the Air Force Office of Scientific Research (AFOSR) under a Multidisciplinary University Research Initiative (MURI) Program (Contract FA9550-05-1-0435) and by the Center for Nanoscale Systems (CNS) at Harvard University, a member of the National Nanotechnology Infrastructure Network (NNIN).


