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Characterizations of Cloud Droplet Shatter Artifacts in Two Airborne Aerosol Inlets

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Aircraft-based aerosol sampling in clouds is complicated by the generation of shatter artifact particles from aerodynamic or impaction breakup of cloud droplets and ice particles in and around the aerosol inlet. Aerodynamic breakup occurs when the Weber number of a droplet, which primarily depends on the droplet size and the magnitude of the relative motion of the droplet and the local air mass, exceeds a critical value. Impaction breakup of a droplet occurs when the droplet’s impaction breakup parameter, $K$, which is a combination of Weber and Ohnesorge numbers, exceeds a critical value. Considering these two mechanisms, the critical breakup diameters are estimated for two aerosol inlets of different designs—a conventional forward-facing solid diffuser inlet (SDI) and a cross-flow sampling sub-micron aerosol inlet (SMAI). From numerical simulations, it is determined that cloud droplets of all sizes will experience impaction breakup in SDI, while only droplets larger than $\sim 16 \, \mu m$ will experience impaction breakup in SMAI. The relatively better in-cloud sampling performance of SMAI is because of its cone design that slows the flow just upstream of the sample tube. The slowing upstream flow, however, causes aerodynamic breakup of drops larger than $\sim 100 \, \mu m$. The critical breakup diameters determined from analysis of field data largely validate numerical predictions. The cross-flow sampling design of SMAI is seen to ensure that shatter artifacts in the inlet are minimal even when there are a significant number of particles larger than the critical breakup size. The study results, thus, suggest that the SMAI design presents an effective approach to sample interstitial particles from aircraft.

1. INTRODUCTION

Aerosol role in cloud formation and the subsequent interaction of cloud droplets with the incoming solar radiation, i.e., the aerosol indirect effect, must be well understood to accurately determine the global radiative budget (Seinfeld and Pandis 2006; Denman et al. 2007; Andreae and Rosenfeld 2008). The aerosol indirect effect is modulated by the properties of the cloud system, i.e., the number concentration and the mean size of the cloud droplets formed, which in turn are dependent on the properties of the aerosol population initially present in the cloud-forming air mass (Rogers and Yau 1989; Seinfeld and Pandis 2006). Aerosol–cloud interaction can also determine the lifetime and fate of a cloud system. The contact of super-cooling aerosol particles with super-cooled cloud droplets can trigger ice formation through contact nucleation and initiate precipitation and, thus, alter the cloud system (Lohmann 2002). Also, as aerosols repeatedly activate and pass through cloud events, aqueous-phase processing and photochemical reactions can alter the nature of the aerosol population (Ervens et al. 2008). The ability to model the formation and evolution of cloud systems and the transformation of aerosol populations during cloud events is critical for large-scale global climate and regional air quality studies. The development and validation of such models requires data from a wide range of aerosol–cloud systems and such data can, often, only be obtained from aircraft-based measurements.

Complete characterization of aerosol particles and cloud droplets from aircraft require the deployment of a large number of instruments. Cloud droplet instruments are located outside the aircraft, where they sample passively from the freestream. The large sizes of the cloud droplets make in situ passive measurements both necessary, because these particles cannot be sampled efficiently through any length of tubing, and sufficient, because physical characterization is possible with relatively open instruments. Aerosol measurements, however, often require large instruments that can only be located inside the fuselage. Accurate aerosol measurements are, therefore, dependent on the ability to representatively sample particles from outside the aircraft.
and efficiently transport them to the cabin. While reasonably error-free sampling is possible in clear air, in-cloud sampling is complicated by the interaction of cloud droplets with aerosol inlets. The shatter of cloud droplets in and around an aerosol inlet can result in the generation of a large number of shatter artifact particles and contamination of the aerosol sample (Hudson and Frisbie 1991; Hudson 1993; Weber et al. 1998; Craig et al. in press). Thus, aircraft-based analysis of interstitial particle population properties has largely been impossible.

The extent of the shatter artifact problem has been seen to depend on droplet properties, inlet geometry, and aircraft conditions. By comparing aerosol measurements from two forward-facing aerosol inlets of different sizes, Weber et al. (1998) determined that shatter artifacts were larger for an inlet with a smaller leading-edge diameter. From wind tunnel tests, Twohy et al. (2003) established that the number of shatter particles generated increased with airspeed. With most aerosol inlets, the extent of the shatter problem is different in warm and cold clouds, with the shatter problem, generally, less significant with ice particles than liquid droplets (Weber et al. 1998; Craig et al. in press). Droplet shatter has also been a problem for cloud and precipitation measurements (Gardiner and Hallett 1985; Field et al. 2003; Korolev and Isaac 2005; Jensen et al. 2009), but advances in cloud probe tip designs and introduction of new high-speed data analyses approaches have helped alleviate some of these problems related to size distribution measurements of activated droplets (Korolev and Isaac 2005; Lawson 2011). There have, however, been very limited efforts to address the problems related to aircraft-based aerosol measurements in clouds.

An illustration of the different liquid cloud droplet shatter mechanisms relevant to aerosol inlet sampling is shown in Figure 1. Droplets can breakup upon impaction on solid surfaces or because of adverse aerodynamic forces in the freestream. When a liquid droplet impacts on a surface, it may rebound, spread over the surface, or shatter (Rein 1993; Mao et al. 1997; Rioboo et al. 2002; Yarin 2006). The exact nature of the droplet response upon impaction depends on droplet properties (e.g., surface tension, viscosity, and size), impaction velocity, and surface conditions (Mundo et al. 1995; Yarin and Weiss 1995; Mao et al. 1997; Rioboo et al. 2001; Yarin 2006). A droplet impacting on a surface will begin to spread out radially from the point of impaction (Mao et al. 1997; Rioboo et al. 2002) and, depending on the contact angle, the liquid layer may recoil back to the impaction point and rebound (nonwetting surfaces), or for low contact angles (wetting surfaces), stay on the surface as a film (Mao et al. 1997; Rioboo et al. 2001). The size of the film increases with increasing impaction velocity and droplet size and decreases with increasing droplet viscosity (Mao et al. 1997; Rioboo et al. 2002). For surfaces with finite dimensions, the film could move toward the edge of the surface, where it can eventually become unattached and spray into secondary droplets (Figure 1a; Gardiner and Hallett 1985; Lawson and Cooper 1990; Emery et al. 2004). Additionally, when a droplet impacts with high momentum, the film may become unstable as it spreads out (Lin and Reitz 1998) and break up, producing secondary droplets (Figure 1b; Mundo et al. 1995; Yarin and Weiss 1995; Rioboo et al. 2001; Yarin 2006). For aircraft-based measurements, this impaction breakup case is the most relevant.

The impaction breakup of liquid droplets is governed by the droplet Weber and Reynolds numbers. If the impacting surface is wet, the presence of the film suppresses secondary droplet generation relative to a dry surface case (Cossali et al. 1997; Wang and Chen 2000; Rioboo et al. 2003). A surface can be considered dry, if the ratio of the surface film thickness to the droplet diameter ($H$) is small ($<\sim0.1$; Wang and Chen 2000). For aerosol inlets on aircraft, because of the large droplet impaction velocities and the low concentrations of the impacting activated droplets, the surface is expected to remain dry after droplet impaction (Povarov and Rastorguev 1986). From experiments of individual liquid droplets impacting on a dry surface, Mundo et al. (1995) determined that the fate of impacting droplets depended on the value of their impaction parameter, $K$, defined as

$$K = \text{OhRe}^{1.25},$$  \hspace{1cm} [1]

where Oh is the Ohnesorge number, defined as

$$\text{Oh} = \frac{\sqrt{\text{We}_{\text{wall}}}}{\text{Re}},$$  \hspace{1cm} [2]

where $\text{We}_{\text{wall}}$ and Re are the Weber and Reynolds number of the impacting droplet, defined as

$$\text{We}_{\text{wall}} = \frac{\rho_d D_d U_d^2}{\sigma_d},$$  \hspace{1cm} [3]

$$\text{Re} = \frac{\rho_d D_d U_d}{\mu_d},$$  \hspace{1cm} [4]
where \( \rho_d \) is the droplet density, \( D_d \) is the droplet diameter, \( U_d \) is the droplet velocity (normal to the surface), \( \sigma_d \) is the surface tension of the droplet, and \( \mu_d \) is the dynamic viscosity of the droplet.

If the \( K \) value of droplets impacting on a surface is more than the critical value \( (K_{\text{crit}}) \) of 57.7, then the droplets will shatter; else they will just deposit on the surface and form a liquid film (Mundo et al. 1995). Corresponding to the \( K_{\text{crit}} \) value, a critical wall impaction breakup diameter, \( D_{\text{crit,wall}} \), i.e., the critical diameter for wall impaction breakup, can be calculated for known inlet dimensions and sampling conditions. Only droplets larger than this critical size will shatter on impaction. For \( K > K_{\text{crit}} \), the number of secondary particles generated from droplet shatter increases highly nonlinearly with the value of \( K \) (O’Rourke and Amsden 2000), and the mean size and standard deviation of the size distribution of the secondary droplets decreases (Mundo et al. 1995).

Cloud droplets traveling toward an inlet may shatter even without impacting on a surface. When the velocities of a droplet and the local air flow are mismatched, the resulting shear on the droplet surface could generate instabilities (Pilch and Erdman 1987; Hsiang and Faeth 1992) causing droplet breakup and secondary particle generation, as illustrated in Figure 1c. This aerodynamic breakup process is governed by the droplet’s aerodynamic Weber number defined as

\[
\text{We}_{\text{aero}} = \frac{\rho_g D_d (U_g - U_d)^2}{\sigma_d},
\]

where \( \rho_g \) is the density of the gas (i.e., air), \( U_g \) is the gas velocity, and \( U_d \) is the droplet velocity magnitude. Particles with Weber numbers greater than a critical aerodynamic breakup Weber number \( (\text{We}_{\text{crit,aero}}) \) of 12 (Wierzba 1990) are expected to shatter, with the exact nature of the droplet breakup mechanism (e.g., vibrational, bag, rim, stripping, etc.; Pilch and Erdman 1987) dependent on the droplet’s \( \text{We}_{\text{aero}} \) number. For selected inlet geometry and operating conditions, a critical aerodynamic breakup diameter, \( D_{\text{crit,aero}} \), can be calculated corresponding to the critical Weber number. Droplets larger than this critical size will shatter aerodynamically.

In addition to impaction and aerodynamic breakup, other possible cloud droplet breakup mechanisms are multi-stage breakup process in which in-coming cloud droplets collide with shattered droplets, which are then subject to further breakup; and aerodynamic breakup after the cloud drop has impacted and shattered on the surface. These shatter mechanisms are likely to be important only under conditions of very high cloud concentrations or when aircraft speeds are much higher, and are not relevant for the current analysis.

In this study, critical cloud droplet breakup sizes, considering aerodynamic and wall impaction breakup mechanisms, are determined from numerical simulations for two aerosol inlets of different designs—a conventional forward-facing aerosol inlet, the Solid Diffuser Inlet (SDI), and a cross-flow inlet, the Sub-Micron Aerosol Inlet (SMAI), which has recently shown promise for interstitial aerosol sampling. The model results are validated with aircraft data obtained from simultaneous operation of the two inlets.

2. FIELD DATA FOR ANALYSIS

The SMAI and SDI aerosol inlets were both simultaneously deployed on the NSF/NCAR C-130 aircraft during the VAMOS Ocean-Cloud-Atmosphere-Land Study (VOCALS) field project during Summer 2008, off the coast of Chile, South America. Analysis of the aerosol number concentration measurements made with the two inlets during this project provides a good basis to assess the relative performance of the two inlets. The schematic diagrams of the SMAI and SDI inlets are shown in Figure 2. The SDI (Figure 2a) is a forward-facing, one-stage diffusing aerosol inlet, with a shroud to align the flow entering the inlet. The diameter of the shroud is 3.14 cm and the entrance diameter of the inlet is 0.51 cm, which then expands to a diameter of 4.47 cm at an angle of 4.5° (Huebert et al. 2004; McNaughton et al. 2007). The sampling efficiency of this inlet has been characterized by Huebert et al. (2004) and McNaughton et al. (2007), and inlet cut size was determined to be ~3–5 \( \mu \)m, aerodynamic diameter. During VOCALS, the SDI was mounted on the side of the aircraft, extending 30.5 cm from the surface of the aircraft skin. The inlet flow was actively controlled to ensure isokinetic-sampling conditions.

The main feature of the SMAI (Figure 2b) design is a flow-through cone facing the freestream flow, with the cone entrance diameter (6.35 cm) being significantly larger (~12×) than that

![Size comparison of the two aerosol inlets flown on the C-130 during VOCALS: (a) SDI and (b) SMAI. The drawings are not to scale.](image-url)
of the SDI. Inside the cone, a perpendicular sampling tube of inner diameter of 1.27 cm is located at a distance of 6.86 cm from the entrance. The perpendicular sampling tube contains a flat-plate sample tube lip that is oriented largely parallel to the bulk flow in the cone. The sample tube lip has an outer diameter of 1.59 cm and a radius of curvature of 0.875 cm. A sample flow rate of 20 lpm is drawn through the perpendicular sample tube, of which 5 lpm is sub-sampled to the instruments in the cabin and the excess flow is expelled out through a scarf tube (Beaton and Spowart 2012; Craig et al. in press). The SMAI sampling characteristics were numerically determined by Craig et al. (in press) and the inlet cut size was determined to be \( \sim 3 \mu m \). The SMAI was placed on the belly of the aircraft, just ahead of the back-bay doors, with the center of the cone located \( \sim 46 \text{ cm} \) from the aircraft skin. The inlet tip dimensions were comparable to those of the SDI, with a cone angle of \( \sim 1^{\circ} \) and a large size of the cone entrance of \( 4.3 \text{ cm} \). The SMAI cone diameter of \( 4.3 \text{ cm} \) is approximately \( \sim 1.3 \text{ million} \) tetrahedral cells. The boundary conditions far from the SMAI were typically chosen to match typical C-130 flight conditions during VOCALS: i.e., an ambient pressure of 875 mb, a static temperature of 283 K, and Mach number of 0.325 (108 m s\(^{-1}\)). A large, external hemisphere of radius 51 cm was selected as the extent of the modeling domain around the SMAI. At the finest grid resolution, the domain was divided into a total of \( \sim 1.3 \text{ million} \) tetrahedral cells. The boundary conditions far from the SMAI were typically chosen to match typical C-130 flight conditions during VOCALS: i.e., an ambient pressure of 875 mb, a static temperature of 283 K, and Mach number of 0.325 (108 m s\(^{-1}\)). In addition to these conditions, to determine the effect of aircraft velocities on the inlet performance, the simulations were repeated for Mach numbers of 0.3 (100 m s\(^{-1}\)) and 0.4 (135 m s\(^{-1}\)).

To characterize the performance of the two aerosol inlets in clouds, specific cloud events were identified from measurements of two-wing-mounted in situ cloud probes: the Cloud Droplet Probe (CDP; Droplet Measurement Technologies (DMT), Inc., Boulder, CO, USA) and the 2-Dimensional Cloud probe (2D-C; Particle Measuring Systems [PMS], Inc., Boulder, CO, USA). The CDP cloud droplet number concentration (CDNC\(_{\text{CDP}}\)) was used to identify small cloud droplet events with sizes between 2 and 50 \( \mu m \), while the 2D-C concentration (CDNC\(_{\text{2DC}}\)) was used to derive large cloud droplet events with droplet sizes between 62.5 and 1537.5 \( \mu m \). During VOCALS, the CDP measurements were in good agreement with that of the Forward Scattering Spectrometer Probe (FSSP-100; PMS, Inc.). Full range bead calibrations of the in situ cloud probes and flow calibrations of the TSI CN counters were conducted pre- and post-project. A summary of the aerosol and cloud instrumentation discussed in this article is provided in Table 1.

### Table 1: Aerosol inlets and cloud probes flown on the C-130 during VOCALS

<table>
<thead>
<tr>
<th>Inlet/cloud probe</th>
<th>Measurement diameter size range</th>
<th>Comment</th>
</tr>
</thead>
<tbody>
<tr>
<td>SDI</td>
<td>( &gt; 10 \text{ nm} )</td>
<td>Instrument: TSI 3760a</td>
</tr>
<tr>
<td>SMAI</td>
<td>( &gt; 10 \text{ nm} )</td>
<td>Instrument: TSI 3760a</td>
</tr>
<tr>
<td>CDP</td>
<td>( 2-50 \text{ ( \mu m ))}</td>
<td>Derived cloud: CDNC(_{\text{CDP}}) ( &gt; 1 \text{ cm}^{-3} )</td>
</tr>
<tr>
<td>2D-C</td>
<td>( 62.5-1537.5 \text{ ( \mu m ))}</td>
<td>Derived cloud: CDNC(_{\text{2DC}}) ( &gt; 0.005 \text{ L}^{-1} )</td>
</tr>
</tbody>
</table>

3. **Computational Fluid Dynamic (CFD) Droplet Breakup Predictions**

To evaluate the critical breakup diameters for aerosol inlets considering aerodynamic and wall impaction mechanisms, knowledge of the flow field in and around the inlets is required. As the SDI operates isokinetically and isoaxially, the streamlines can be assumed to be straight lines headed straight toward the inlet and the droplet velocities can be assumed to be the same as that of the flow. Droplets headed toward the inlet can impact on the shroud and the leading edge of the inlet (i.e., the inlet tip). The impaction of droplets on the shroud can be ignored, because any secondary particles generated there will largely remain close to the shroud, and away from the sample flow streamlines. Only shatter particles generated from droplet impaction breakup at the inlet tip will likely contaminate the aerosol sample. The critical impaction droplet diameter will, thus, be determined based on the entrance tip dimensions. In the isokinetic-sampling SDI, the particle and the flow velocities will be matched and hence no aerodynamic droplet breakup is expected upstream of the inlet. It is, however, possible for large droplets that enter the SDI to breakup aerodynamically when the flow slows down inside the inlet. The critical aerodynamic breakup diameters for SDI will, thus, be estimated considering the inlet diffuser section.

### 3.1. SMAI Flow Pattern

The SMAI flow field is more complicated than the SDI’s and the knowledge of the exact flow field is critical for accurate determination of critical breakup diameters. The CFD software, FLUENT® (ANSYS, Inc., Lebanon, NH, USA) is used to calculate the three-dimensional flow field in and around the SMAI, with the flow treated as steady, compressible, and turbulent. The turbulence in the flow was modeled using the two-equation standard \( k-\varepsilon \) turbulence model (the Reynolds number based on the average SMAI cone diameter of \( 4.3 \text{ cm} \) is \( \sim 10^{5} \)). A large, external hemisphere of radius \( 51 \text{ cm} \) was selected as the extent of the modeling domain around the SMAI. At the finest grid resolution, the domain was divided into a total of \( \sim 1.3 \text{ million} \) tetrahedral cells. The boundary conditions far from the SMAI were typically chosen to match typical C-130 flight conditions during VOCALS: i.e., an ambient pressure of 875 mb, a static temperature of 283 K, and Mach number of 0.325 (108 m s\(^{-1}\)). In addition to these conditions, to determine the effect of aircraft velocities on the inlet performance, the simulations were repeated for Mach numbers of 0.3 (100 m s\(^{-1}\)) and 0.4 (135 m s\(^{-1}\)).

The CFD results of the velocity field in and around the SMAI for the Mach number case of 0.325 are shown in Figure 3. The combination of a small flow through the SMAI (relative to that directed toward the inlet) and the large size of the cone entrance results in a region of high pressure just upstream of the inlet. The cone, therefore, acts as a virtual blunt body that decelerates the flow upstream of the inlet, with the gas velocity reaching a minimal value of \( \sim 10 \text{ m s}^{-1} \) just after the cone entrance. Inside
the cone, the velocity increases and the flow reaches a velocity of \( \sim 40 \text{ m s}^{-1} \) at the location of the sample tube.

### 3.1.1. Breakup Predictions

In order to determine the critical breakup diameters for each mechanism, the trajectories of liquid water droplets over a size range of 1–500 \( \mu \text{m} \) were tracked from the freestream to the sample tube. The droplets were injected far upstream of the inlet with velocities matched to the local flow conditions. Because of their large sizes, cloud droplets traversing through regions of large gas velocity gradients have finite Reynolds numbers and thus, a Reynolds-number-dependent particle drag coefficient is used for the droplet trajectory calculations (Hinds 1999; Baron and Willeke 2005). The combination of small residence time of the droplets in the vicinity of the inlet surface and their large sizes makes turbulent dispersion unimportant for the current simulations. The trajectories are, therefore, calculated only considering the mean-flow velocities. The density of the tracked droplets was increased until the results remained invariant with the number of injections.

#### 3.1.1.1. Aerodynamic breakup.

To characterize aerodynamic breakup in SMAI, the aerodynamic Weber numbers, \( W_{\text{aero}} \), were calculated along the droplet trajectories, starting from a location well upstream of the cone entrance (–30 cm) to downstream of the sample tube (+8 cm) and these values were compared against the critical breakup value. Upstream of the cone, the \( W_{\text{aero}} \) numbers were calculated within the area bounded by the diameter of the cone entrance, while inside the inlet, the spatial extent of the calculations were restricted by the cone geometry.

At a selected axial location (X) within the bounded area, drops of a given size will have a range of \( W_{\text{aero}} \) numbers corresponding to the radial variation in the flow velocities. The average, minimum, and maximum \( W_{\text{aero}} \) numbers at an axial location are calculated for different drop sizes and are shown in Figure 4. For a selected droplet diameter, \( W_{\text{aero}} \) increases as the droplet approaches the inlet entrance. This is because, as the gas decelerates upstream of the inlet, the finite relaxation times of the large droplets will result in significant relative velocities for the particles compared to that of the local air mass. Proportional to the particle relaxation times, the \( W_{\text{aero}} \) numbers increase with particle diameters. The maximum droplet \( W_{\text{aero}} \) number is at the entrance to the cone, where the gas velocity is the lowest. Thus, the smallest droplet diameter that will shatter, i.e., \( D_{\text{crit,aero}} \), must be determined from the \( W_{\text{aero}} \) values at the cone entrance. From the CFD flow field calculations, the \( D_{\text{crit,aero}} \) value for SMAI is determined to be \( \sim 100 \mu \text{m} \), with droplets smaller than this size able to enter the inlet without experiencing aerodynamic breakup. The Weber number of droplets larger than 100 \( \mu \text{m} \) will exceed the critical value upstream of the cone entrance and these particles will, thus, shatter prior to entering the inlet and contribute to shatter artifacts in the sample flow. At higher aircraft velocities, the critical droplet sizes decrease, and making it more likely for even smaller droplets to shatter aerodynamically. For SDI, the aerodynamic breakup region is inside the inlet where the flow slows down and the critical size calculated assuming freestream droplet velocities and an average sample flow velocity inside the inlet is \( \sim 100 \mu \text{m} \), similar to the SMAI.

#### 3.1.1.2. Wall impingement.

While droplets smaller than 100 \( \mu \text{m} \) will enter the SMAI without shattering aerodynamically, some of these droplets may break up when they impact...
on the inlet walls, if their $K$ value exceeds $K_{\text{crit}}$. Possible impaction surfaces include the outer rim of the cone, inside walls of the cone, and the sample tube lip. The only shatter-generated secondary droplets likely to be aspirated into the perpendicular sample tube are those generated from impaction on the sample tube lip. The $D_{\text{crit,wall}}$ values for SMAI are, therefore, calculated based on conditions at this location. For SDI, considering isoaxial- and isokinetic-sampling conditions, the $K$ values are calculated assuming that the droplets impact with the freestream velocity at the critical impaction area, the leading edge of the inlet (i.e., the inlet tip). The $K$ values as a function of droplet diameter for the SMAI and SDI are shown in Figure 5.

The impaction $K$ values increase with droplet diameter because larger droplets impact with greater kinetic energies and, thus, break up more easily. The presence of the high pressure region upstream of the SMAI helps droplets relax to the slowing flow conditions, and lowers their momentum prior to reaching the sample tube lip, i.e., lowers the droplet $K$ values. Thus, compared to the SDI, the critical impaction diameters for the SMAI are much larger. The $D_{\text{crit,wall}}$ value for SDI is $\sim 2 \mu m$, while for SMAI it is $\sim 10$–$20 \mu m$. For the SMAI, a range of critical droplet diameters exists, because the impaction velocities of the droplets vary over the different sample tube lip locations.

As almost all cloud droplets will impact and shatter on the SDI, the SDI sample will always be contaminated in clouds. For SMAI, CFD predictions suggest droplets as large as $100 \mu m$ can enter the inlet without experiencing aerodynamic breakup, though droplets in the diameter range of $20$–$100 \mu m$ may impact on the sample tube lip and generate shatter artifacts.

4. EXPERIMENTAL RESULTS

The CFD predictions of critical breakup diameters for the two inlets can be validated from analysis of aircraft measurements of aerosol number concentrations made with the two inlets during the VOCALS campaign. The in-cloud sampling performance of the inlets will be characterized using an enhancement factor, $E$, defined as (Craig et al. in press)

$$E = \frac{\text{CN}_{\text{inlet}}}{\text{CN}_{\text{bg}}},$$

where $\text{CN}_{\text{inlet}}$ is the in-cloud CN concentration measured by an inlet and $\text{CN}_{\text{bg}}$ is the estimated background CN concentration for that same air mass that would have existed in the absence of cloud formation. Following the approach of Craig et al. (in press), the $\text{CN}_{\text{bg}}$ values are calculated as 15-s averages of the clear-air CN concentrations in the vicinity of an analyzed cloud system. The only data sets considered for the enhancement factor analysis are those with near-cloud clear-air masses having the same potential temperatures ($\pm 1$ K) as the analyzed cloud system. This definition of $\text{CN}_{\text{bg}}$ ensures that the enhancement factors are calculated considering CN concentrations from air masses of similar origin.

In clouds, because of the activation of some background particles to form droplets, the interstitial aerosol concentration should be lower than the background CN concentration, and, thus, in the absence of shatter artifacts, the measured enhancement factors should be $<1$. The enhancement factors are used to determine the in-cloud sampling performance of SMAI and SDI for two distinct cloud cases: (1) small cloud droplet events, when only droplets smaller than $50 \mu m$ were present and (2) large cloud drop events, when the only drops present were larger than $50 \mu m$.

4.1. Small Cloud Droplet Events

The performance of the SDI in cloud systems where the mean droplet sizes are only in the range of $2$–$50 \mu m$ (i.e., when no drops are recorded by 2D-C) is shown in Figure 6. For all CDP cloud concentrations, the SDI enhancements, $E_{\text{SDI}}$, are greater than 1 and enhancements are seen to increase with the mean cloud droplet diameter, $D_{\text{CDP,mean}}$ (obtained from CDP measurements). Consistent with CFD predictions, aerosol concentrations measured by the SDI inside clouds are significantly enhanced for all cloud droplet sizes. Thus, in-cloud aerosol measurements are not possible with the SDI and other similar forward-facing isokinetic aerosol inlets.

For the case of small cloud droplet events, the SMAI enhancements, $E_{\text{SMAI}}$, are shown in Figure 7. The SMAI enhancements are significantly reduced compared to that of the SDI. The $E_{\text{SMAI}}$ values decrease with increasing cloud concentrations, consistent with the expectation of reduced interstitial aerosol in clouds as the concentration of activated particles is increased. The $E_{\text{SMAI}}$ values are seen to increase with increasing $D_{\text{CDP,mean}}$. 

![Figure 5](image-url)
FIG. 6. The calculated \( E_{\text{SDI}} \) for cloud systems containing only small cloud droplets (i.e., no drizzle). The enhancements are shown as a function of (a) cloud droplet number concentration from the CDP, \( \text{CDNC}_{\text{CDP}} \), and (b) mean CDP cloud droplet diameter, \( D_{\text{CDP,mean}} \). The points represent 1 s data.

At a selected mean cloud droplet diameter, a significant variability in the \( E_{\text{SMAI}} \) values is observed (Figure 7) and this can be attributed to (1) possible inaccuracies in the estimation of \( \text{CN}_{\text{bg}} \) values, and/or (2) statistical fluctuations in incident drop sizes. Accurate estimate of \( \text{CN}_{\text{bg}} \) will require knowledge of cloud base CN concentrations and consideration of possible dilution of CN concentrations because of mixing of in-cloud and out-of-cloud air masses and entrainment of descending air mass. Thus, \( E \) values greater than 1 may represent an under prediction of \( \text{CN}_{\text{bg}} \) concentrations rather than shatter from the inlet. In addition, for the same mean size, a range of different cloud droplet size distributions are possible, particularly with variable number of large particles in the tail of the distribution. Thus, in an inlet sample volume, the likelihood of encountering large sized droplets that could shatter will differ with the cloudy air mass, resulting in variable inlet enhancement values. Relative to the differences in the magnitude of the \( E \) values for the two inlets, the variability in the values is insignificant and, hence, can be ignored.

To predict the critical breakup diameter from experimental data, time periods corresponding to a narrow range of total CDP concentration (say, 54 cm\(^{-3} \pm 10\%\)) were identified. During these time periods, while the total CDP concentrations are nearly constant, the size distributions could be quite variable. The variation in SMAI enhancements during these time periods was correlated with CDP concentrations in different size bins (mid-point bin diameters, \( D_{\text{mid}} \), of 8.88, 11.25, 13.63, 16.00, 18.38, and 21.30 \( \mu \text{m} \)). The resultant slopes of a linear fit to this correlation (\( E_{\text{SMAI}} \) vs. CDP bin concentration) were determined and are shown in Figure 8. The uncertainty bars in the plot correspond to the range of slope values obtained at the different total cloud concentrations (54, 110, 130, 200, 230, 300, and 340 cm\(^{-3} \pm 10\%)\). When the mean cloud droplet sizes are smaller than \( \sim 14 \mu \text{m} \), the expected relationship of decreasing SMAI CN concentration with increasing cloud concentration (i.e., negative slope) is observed. When the mean cloud droplet sizes are larger than \( \sim 18 \mu \text{m} \), increasing the cloud concentrations is seen to result in increasing SMAI CN concentration for all cloud concentrations (i.e., positive slope). This analysis suggests that the critical cloud droplet diameter for breakup is \( \sim 16 \mu \text{m} \), and with bigger droplets likely to contribute to shatter artifacts, largely consistent with CFD simulations.

4.2. Large Cloud Drop Events

The performance of SMAI when only drizzle was present (i.e., drops were detected by 2D-C, but no droplets were detected by the CDP) is shown in Figure 9. In general, SMAI CN enhancements are seen to increase with cloud concentration, but
FIG. 9. The calculated $E_{\text{SMAI}}$ for cloud systems containing only drizzle drops. The enhancements are shown as a function of (a) 2D-C cloud droplet number concentration, CDNC$_{2D-C}$, and (b) 2D-C mean cloud diameter, $D_{2D-C,\text{mean}}$. Each point represents 1 s data.

no clear trend with cloud drop size is observed, and no obvious evidence of aerodynamic breakup is visible. Even in the presence of drops much larger than the critical impaction and aerodynamic size, enhancement values less than 1 are observed. To explain these observations, a Monte Carlo-based simulation was conducted to determine the number of cloud drops intercepted by the critical impaction and aerodynamic breakup zones as a function of cloud number concentration. For impaction breakup, the critical zone was considered as the area of the central part of the sample tube lip ($1.6e^{-2}$ cm$^2$), over which most of the sample flow passed through. For the aerodynamic breakup, the critical zone was considered as the area at the cone entrance through which the sample flow streamlines passed through ($\sim 1.3e^{-1}$ cm$^2$ for a sample flow of $\sim 20$ lpm; Beaton and Spowart 2012).

To calculate the flux of drops through the two critical breakup areas, the number of drops in the volume swept by SMAI in 1 s was determined, and these particles were then distributed randomly over the SMAI cone entrance area. The number of drops in the two critical areas was then calculated. The random distribution of particles was repeated several times, and the mean particle flux rate was determined as a function of cloud drop concentration, as shown in Figure 10. As expected, the number of drops intercepted by the aerodynamic breakup zone is significantly higher than that by the impaction breakup zone. Thus, aerodynamic breakup should become important at much lower concentrations than impaction breakup.

In the 1 Hz CN concentration data set, observation of a consistent enhancement due to aerodynamic breakup will require at least 1 drop per second being present in the critical flow volume. The cloud concentration corresponding to a freestream velocity of 108 m s$^{-1}$ and the projected area of the sample tube lip (relevant for impaction) is $\sim 7$ L$^{-1}$, while the concentration corresponding to the aerodynamic breakup zone is 0.8 L$^{-1}$. Based on these critical concentrations, it can be hypothesized that when the cloud drops present are all smaller than $D_{\text{crit,aero}}$, enhancement in CN concentrations requires impaction breakup of these drops and, thus, cloud drop concentrations must be $\geq 7$ L$^{-1}$; and when cloud drops larger than $D_{\text{crit,aero}}$ are present, CN enhancements should become apparent at lower cloud concentrations ($\geq 0.8$ L$^{-1}$).

To identify aerodynamic breakup in the SMAI, sampling times were identified when only drizzle drops were present (i.e., no droplets detected by CDP) and the drop number concentrations ($N_{\text{bin}}$) were calculated in two size bins: diameters, $D \leq 112.5$ μm and $D \geq 212.5$ μm. The CN enhancements as a function of cloud drop number concentrations for the two size ranges are shown in Figure 11. For $D \leq 112.5$ μm, the SMAI enhancements start to increase above 1 only for large concentrations ($> 10$ L$^{-1}$), largely consistent with the predicted critical impaction number concentration for SMAI. Thus, these drops do not shatter aerodynamically. For $D \geq 212.5$ μm, CN enhancements increase significantly at concentrations $> 1$ L$^{-1}$, consistent with the prediction of aerodynamic breakup. These results provide the first direct evidence of aerodynamic breakup in airborne inlets.

4.2.1. Uncertainties

As the SMAI was located in the belly, for a cruise pitch of $\sim 4^\circ$, the SMAI is not expected to be in a droplet shadow region (King 1984). The SDI is at a distance of 30.5 cm from the skin of the aircraft and is located on the side of the aircraft, upstream of the wing. Following the analysis of King (1984), it can be predicted that at the SDI location, a shadow region for droplets in the size range of 100–500 μm might exist. This shadow region, however, does not affect our analysis of SDI.
TABLE 2
Critical breakup diameters for SDI and SMAI

<table>
<thead>
<tr>
<th>Aerosol inlet</th>
<th>Inlet type</th>
<th>Sample method</th>
<th>Sampling cut size (μm)</th>
<th>CFD $D_{\text{crit, aero}}$ (μm)</th>
<th>CFD $D_{\text{crit, wall}}$ (μm)</th>
<th>Experimental critical diameter (μm)</th>
</tr>
</thead>
<tbody>
<tr>
<td>SDI</td>
<td>Diffuser</td>
<td>Isokinetic/parallel</td>
<td>3–5</td>
<td>~100</td>
<td>2</td>
<td>2</td>
</tr>
<tr>
<td>SMAI</td>
<td>Forward-facing cone</td>
<td>Sub-isokinetic/anisoaxial (90°)</td>
<td>3</td>
<td>~100</td>
<td>~20</td>
<td>~16</td>
</tr>
</tbody>
</table>

FIG. 11. SMAI enhancement as a function of cloud drop concentration ($N_{\text{bin}}$) in different size bins, $D \leq 112.5$ μm and $D \geq 212.5$ μm. In selecting data points for $D \leq 112.5$ μm, a constant concentration of $D > 112.5$ μm was maintained (1 ± 0.25 L$^{-1}$). Similarly for $D > 212.5$ μm, the data points were selected such that the concentrations of $D < 212.5$ μm were constant (1 ± 0.25 L$^{-1}$). Note that the exact choice of the bin concentrations did not affect the analysis.

Analysis of the performance of the new SMAI inlet design suggests that this inlet is largely unaffected by the presence of cloud droplets when the mean sizes are smaller than ~16 μm. Even when significant concentrations of droplets larger than the critical size are present, the cross-flow design of the inlet ensures that the enhancements in SMAI CN concentrations are minimal. The SMAI design provides two advantages for in-cloud sampling: its flow field is such that the critical impaction diameter is significantly increased over that of the SDI, and using a perpendicularly sub-sampling tube minimizes the critical impaction area of the inlet, and thus minimizes the number of droplets impacting on the inlet. The critical impaction breakup diameter was determined from CFD simulations to be ~20 μm, and the experimental results largely match predictions. The CFD calculations predict that the SMAI design will induce aerodynamic breakup of cloud drops larger than 100 μm and preliminary analysis of the field data confirms the role of aerodynamic breakup for droplets larger than 100 μm. The critical sampling area for aerodynamic breakup is large enough that this breakup mechanism becomes important even for low drizzle concentrations.

As the SMAI samples are seen to be shatter free or minimally impacted by shatter in the absence of drizzle, SMAI measurements can be analyzed for size distribution and composition measurements of the nonvolatile component of interstitial aerosol in a wide range of nonprecipitating cloud systems. Thus, sensitive aerosol–cloud closure studies may finally be possible. To fully establish the in-cloud sampling performance of SMAI, its measurements must be compared against other inlet designs and under different cloud conditions. Toward that goal, the SMAI was flown during a recent Ice in Clouds Experiment–Tropics (ICE-T; July 2011) campaign and its performance during those tests are compared against that of a new interstitial aerosol sampler, called the Blunt-body Aerosol Sampler (BASE), and those results will be published elsewhere (Moharreri et al. in press).

REFERENCES