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Search for Doubly Charged Higgs Bosons Decaying to Dileptons in $p\bar{p}$ Collisions at $\sqrt{s}$ =1:96 TeV

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We present the results of a search for doubly charged Higgs bosons ($H^{\pm\pm}$) decaying to dileptons (ll') using $\approx 240\ \mathrm{pb}^{-1}$ of $p\bar{p}$ collision data collected by the CDF II experiment at the Fermilab Tevatron. In our search region, given by same-sign ll' mass $m_{ll'} > 80\ \mathrm{GeV}/c^2$ ($100\ \mathrm{GeV}/c^2$ for ee channel), we observe no evidence for $H^{\pm\pm}$ production. We set limits on $\sigma(p\bar{p} \to H^{++}H^{--} \to l^+l'^+l^-l'^-)$ as a function of the mass of the $H^{\pm\pm}$ and the chirality of its couplings. Assuming exclusive same-sign dilepton decays, we derive lower mass limits on $H_L^{\pm\pm}$ of 133, 136, and 115 GeV/c^2 in the ee, $\mu\mu$, and $e\mu$ channels, respectively, and a lower mass limit of 113 GeV/c^2 on $H_R^{\pm\pm}$ in the $\mu\mu$ channel, all at the 95% confidence level.

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The standard model (SM) gives a good description of the known fundamental particles, using the $\mathrm{SU}(3)_C \times \mathrm{SU}(2)_L \times \mathrm{U}(1)_Y$ gauge group to describe their nongravitational interactions. The $\mathrm{SU}(2)_L \times \mathrm{U}(1)_Y$ electroweak gauge symmetry is broken to $\mathrm{U}(1)_{\mathrm{EM}}$ by the Higgs mechanism, but a Higgs boson has yet to be observed. In addition to the SM $\mathrm{SU}(2)_L$ Higgs doublet, a number of models [1–3] predict new Higgs doublets or triplets containing doubly charged Higgs bosons ($H^{\pm\pm}$). For example, the left-right symmetric model [2], predicated on a right-handed version of the weak force $\mathrm{SU}(2)_R$, requires a Higgs triplet. The model predicts light neutrino masses by the seesaw mechanism [4], consistent with recent data on neutrino oscillations [5]. Furthermore, the left-right symmetric model suggests light [$\mathcal{O}(100\ \mathrm{GeV}/c^2)$] doubly charged Higgs particles if supersymmetry is a property of nature [3] and is therefore of interest for direct searches at high-energy colliders.

$H^{\pm\pm}$ bosons couple directly to leptons, photons, W and Z bosons, and singly charged Higgs bosons ($H^{\pm}$). The $H_L^{\pm\pm}$ and $H_R^{\pm\pm}$ bosons, respectively, couple to left- and right-handed particles and may have different fermionic couplings. Their coupling to a pair of W bosons is experimentally constrained to be small due to the small observed value of $|\rho_{EW} - 1|$ [6], resulting in a negligible cross section for the process $p\bar{p} \to W^{\pm} \to W^{\mp}H^{\pm\pm}$. $H^{\pm\pm}$ production would be dominated by the reaction $p\bar{p} \to Z/\gamma^* \to H^{++}H^{--}$, whose cross section is independent of the $H^{\pm\pm}$ fermionic couplings at tree level.

The $H^{\pm\pm}$ decays predominantly to charged leptons if $m_{H^{\pm\pm}} < 2m_{H^{\pm}}$ and $m_{H^{\pm\pm}} - m_{H^{\pm}} < m_{W^{\pm}}$ [7]. The leptonic decays conserve the quantum number $B - L$, where B is the baryon number and L is the lepton number. The $H^{\pm\pm}$ couplings $h_{ll'}$ to electrons and muons are experimentally constrained by the absence of $H^{\pm\pm}$ production in e^+e^- collisions ($h_{ee} < 0.05$) [8], and the nonobservation of the decays $\mu \to 3e$ ($h_{ee}h_{e\mu} < 3.2 \times 10^{-7}$) and $\mu \to e\gamma$ ($h_{\mu\mu}h_{e\mu} < 2 \times 10^{-6}$) [9]. The experimental constraints on the couplings (quoted here for $m_{H^{\pm\pm}} = 100$ GeV$/c^2$) weaken with increasing $H^{\pm\pm}$ mass. The $h_{\mu\mu}$ coupling is probed by measurements of the anomalous magnetic moment of the muon $(g-2)_\mu$; the previous limit $h_{\mu\mu} < 0.25$ [9] has not been reanalyzed using the most recent $(g-2)_\mu$ measurement [10].

Direct searches by the OPAL, L3, and DELPHI Collaborations in e^+e^- collisions [11] have excluded $H^{\pm\pm}$ bosons below masses of about 100 GeV$/c^2$, assuming exclusive $H^{\pm\pm}$ decay to a given dilepton channel. A recent search by the D0 Collaboration in the $\mu\mu$ channel [12] has excluded $H_L^{\pm\pm}$ below a mass of 118 GeV$/c^2$. In this Letter, we describe a search for doubly charged resonances in the same-sign ee, $e\mu$, and $\mu\mu$ channels, using ≈ 240 pb^{-1} [13] of data collected at $\sqrt{s} = 1.96$ TeV by the CDF (Collider Detector at Fermilab) II experiment at the Fermilab Tevatron. We present our results using the $H^{\pm\pm}$ production model [4] and set the world's highest mass limits for a range of couplings to e and μ. We probe the range $10^{-5} < h_{ll'} < 0.5$, which corresponds to narrow resonances that decay promptly ($c\tau < 10$ μm, where τ is the lifetime).

The CDF II detector [14] consists of an inner tracking detector, a lead (iron) scintillator sampling calorimeter for measuring electromagnetic (hadronic) showers, and outer drift chambers for muon identification. The inner detector includes a high-resolution wire chamber [the central outer tracker (COT) [15]] which, along with the central calorimeter and muon system, covers the pseudorapidity interval $|\eta| < 1$ [16].

Our strategy is to search for one of the pair-produced $H^{\pm\pm}$ bosons to maximize the sensitivity and to permit detection of any singly produced doubly charged resonance. The event triggers can be classified by the requirements of (i) two energy clusters with $E_T > 18$ GeV in the electromagnetic calorimeter (2EM), (ii) a central electromagnetic cluster with $E_T > 18$ GeV and matching track $p_T > 9$ GeV$/c$ (1EM), or (iii) a COT track with $p_T > 18$ GeV$/c$ with an associated track segment ("stub") in the muon detectors.

The same-sign ee sample is selected primarily using the 2EM trigger. In the off-line analysis, we require two same-sign central electrons with calorimeter $E_T > 30$ GeV and COT track $p_T > 10$ GeV$/c$. Electrons are identified using the ratio of calorimeter energy (E) to track momentum (p) ($\frac{E}{pc} < 4$), longitudinal and lateral shower profiles, track-cluster matching, calorimeter isolation energy in a surrounding cone, and photon-conversion identification using the tracker. The same-sign ee sample corresponds to an integrated luminosity of (235 ± 13) pb^{-1}. The luminosity is determined by measuring the rate of inelastic collisions, and the uncertainty has equal contributions from the uncertainties on the inelastic cross section and on the acceptance of the luminosity counters.

The same-sign $\mu\mu$ sample is selected using the single-muon trigger, with a consistent off-line requirement of a matching stub. We select tracks with $p_T > 25$ GeV$/c$ that are minimum ionizing, i.e., have small electromagnetic and hadronic energy depositions in the calorimeters. The cosmic-ray muon background is suppressed by requiring the muons to originate from the beam line, to be coincident in time with each other and with a $p\bar{p}$ collision, and to be consistent with a pair of outgoing particles [17]. Track-quality requirements and calorimeter isolation suppress hadronic-jet backgrounds. The integrated luminosity of the same-sign $\mu\mu$ sample is (242 ± 14) pb^{-1}.

The same-sign $e\mu$ sample is selected mainly using the 1EM trigger. We require a central electron and a track matched to a muon stub. The stub requirement significantly reduces background, but also reduces the fiducial acceptance of $H^{\pm\pm} \to e\mu$ relative to the $\mu\mu$ and ee samples. The integrated luminosity of the same-sign $e\mu$ sample is (240 ± 14) pb^{-1}. All electron and muon tracks are constrained to the transverse position of the beam to improve their momentum resolution.

We calculate trigger efficiencies using separate unbiased triggers, and the tracking and lepton-identification efficiencies using $Z \to ee/\mu\mu$ events. We obtain $(96.6 \pm 0.4)\%$ and$(100.00^{+0.00}_{-0.02})\%$ as the efficiencies of the 1EM and 2EM triggers, respectively. The muon trigger efficiencies, including the matching-stub requirements, are $(77.1 \pm 1.3)\%$ and $(93.9 \pm 0.8)\%$ for $|\eta| < 0.6$ and $0.6 < |\eta| < 1$, respectively, each corresponding to a separate detector subsystem. The tracking efficiency is high (>99%) for isolated particles within the COT fiducial volume. The lepton-identification efficiencies are $(92.7 \pm 0.3)\%$ and $(90.8 \pm 0.2)\%$ for electrons and muons, respectively. The corresponding efficiencies measured in simulated [18] Z events are $(89.3 \pm 0.1)\%$ and $(91.3 \pm 0.1)\%$. The simulated $H^{\pm\pm}$ detection efficiency is corrected by the ratio of data to simulated Z boson efficiencies.

The potential backgrounds from SM processes are (i) hadrons that decay to leptons or are misidentified as such, (ii) leptonic decays of W bosons, produced in association with hadronic jet(s) (W + jet), (iii) Z/γ^* decays (Drell-Yan), where the same-sign track comes from a photon conversion, (iv) WZ production, where both the W and Z decay leptonically, and (v) cosmic rays.

The hadronic background is estimated using lepton-triggered events with two same-sign lepton candidates [21], each failing the identification requirements ("failing lepton candidate"). The ratio of the number of lepton candidates passing to the number failing the requirements (the "pass-fail ratio") is measured using jet data samples. These samples are selected either using $E_T > 100$ GeV or $E_T > 20$ GeV jet triggers, or using single-lepton triggers and excluding leptonic W and Z decays. The pass-fail ratio is $\mathcal{O}(0.05)$, with a systematic uncertainty of $\approx 80\%$ arising from its sample dependence. It is used to apply a weight to each candidate lepton (as a function of E_T) in events with two failing lepton candidates to obtain the dilepton mass distribution.

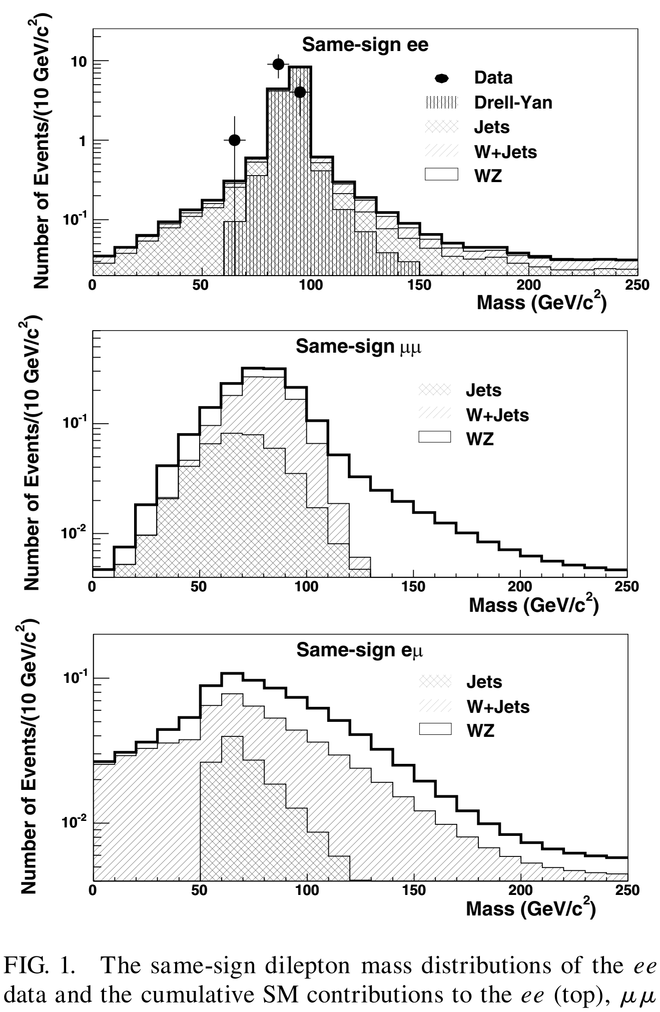

FIG. 1. The same-sign dilepton mass distributions of the ee data and the cumulative SM contributions to the ee (top), $\mu\mu$ (middle), and $e\mu$ (bottom) samples. The solid line is the overall sum of the indicated areas. No same-sign $\mu\mu$ or $e\mu$ events are observed.

The W + jet background is determined by applying the pass-fail ratio as a weight to W data events which have a second failing lepton and $25 < \not{E}_T < 60$ GeV. The expected misidentified-W contribution (from jets) is subtracted to prevent double counting. We use simulated [18] W + jet events to correct for the acceptance of the $\not{E}_T$ requirement. Background from $W\gamma$ production, where the photon converts to an e^+e^- pair, is implicitly included in this estimate. It is studied explicitly using the simulation and found to be negligible.

Background from $Z/\gamma^* \to e^+e^-$ occurs when one electron radiates a photon which subsequently converts to an e^+e^- pair. When a same-sign conversion electron has higher momentum than the prompt electron and is associated with the cluster, the event is reconstructed with two same-sign electrons. The mass dependence is obtained from simulated [18] Drell-Yan events. The simulated sample is normalized using the number of same-sign candidates in the Z mass region ($80\ \text{GeV}/c^2 < m_{ee} < 100\ \text{GeV}/c^2$), after subtracting jet and W + jet contributions.

Background from $WZ \to l\nu ll$ production is estimated using simulation [18]. We use the production cross section of 4.0 pb [22] and apply the trigger, tracking, and lepton-identification efficiencies to the events that pass the kinematic and geometric selection.

The cosmic-ray background is estimated using COT timing information. We use an independently identified sample of cosmic rays to estimate the residual contribution surviving the timing requirements made in the $\mu\mu$ analysis. The expected cosmic-ray background is found to be 0.02 ± 0.02 events, which we take to be negligible.

Figure 1 shows the total background and the data as a function of $m_{ll'}$ for each sample. The predominantly back-to-back lepton topologies, the kinematic thresholds, and

TABLE I. Integrated backgrounds in the low-mass ($<80\ \text{GeV}/c^2$) and high-mass (100–$300\ \text{GeV}/c^2$ for ee, 80–$300\ \text{GeV}/c^2$ for $\mu\mu$ and $e\mu$) regions.

Background	Low-mass region	High-mass region
$Z/\gamma^* \to ee$	0.46 ± 0.13	0.37 ± 0.11
Jets $\to ee$	$0.47^{+0.23}_{-0.19}$	$0.62^{+0.71}_{-0.44}$
W + jet $\to ee$	0.14 ± 0.08	0.36 ± 0.21
$WZ \to ee$	0.07 ± 0.02	0.11 ± 0.03
Total ee	1.1 ± 0.4	$1.5^{+0.9}_{-0.6}$
Jets $\to \mu\mu$	$0.30^{+0.24}_{-0.16}$	$0.19^{+0.35}_{-0.17}$
W + jet $\to \mu\mu$	0.32 ± 0.22	0.40 ± 0.27
$WZ \to \mu\mu$	0.21 ± 0.04	0.19 ± 0.03
Total $\mu\mu$	0.8 ± 0.4	$0.8^{+0.5}_{-0.4}$
Jets $\to e\mu$	0.09 ± 0.05	0.06 ± 0.05
W + jet $\to e\mu$	$0.22^{+0.24}_{-0.15}$	0.25 ± 0.17
$WZ \to e\mu$	0.12 ± 0.02	0.12 ± 0.03
Total $e\mu$	0.4 ± 0.2	0.4 ± 0.2

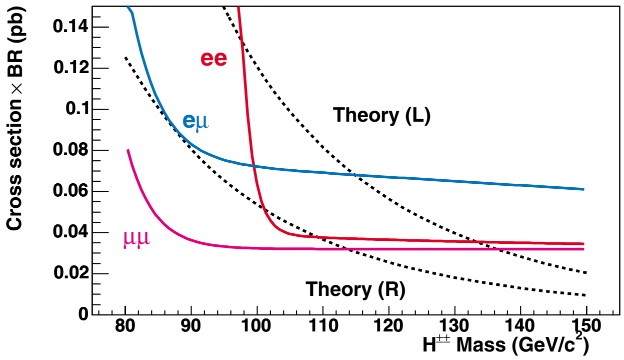

FIG. 2 (color). Experimental limits on cross section × branching ratio (BR) at 95% C.L. as a function of $H^{\pm\pm}$ mass (solid curves). Dotted curves show the theoretical next-to-leading order total cross sections [25] for $H_L^{\pm\pm}$ and $H_R^{\pm\pm}$.

the typical lepton p_T from W or Z decays lead to the observed peaked shapes of the background distributions. The search is performed in the region of $m_{ll'} > 80\ \mathrm{GeV}/c^2$ for the $\mu\mu$ and $e\mu$ samples, and in the region of $m_{ee} > 100\ \mathrm{GeV}/c^2$ for the ee sample. The low-mass regions ($m_{ll'} < 80\ \mathrm{GeV}/c^2$) are used to check our background predictions. Table I summarizes the total background predictions. We estimate 1.1 ± 0.4 (ee), 0.8 ± 0.4 ($\mu\mu$), and 0.4 ± 0.2 ($e\mu$) events in the low-mass regions and observe one ee event ($m_{ee} = 70\ \mathrm{GeV}/c^2$) and no $\mu\mu$ or $e\mu$ events. As an additional check, we compare the predicted and observed backgrounds for same-sign dilepton events with one failing lepton candidate and $\not{E}_T < 15$ GeV. The expectations of 54 ± 21 (ee), 7.6 ± 3.1 ($\mu\mu$), and 2.4 ± 0.8 ($e\mu$) events are consistent with the observed numbers of 63 (ee), 8 ($\mu\mu$), and 2 ($e\mu$) events.

The same-sign dilepton mass resolution is ≈ 3.5% of the mass. The intrinsic $H^{\pm\pm}$ width is equal to $\sum_{l,l'} h_{ll'}^2 m_{H^{\pm\pm}}/8\pi$ [6] and contributes negligibly to the reconstructed mass if $\sum_{l,l'} h_{ll'}^2 < 0.5$. We define search windows of ±10% of a given $H^{\pm\pm}$ mass, corresponding to a $\pm 3\sigma$ window. We predict the acceptances as a function of $H^{\pm\pm}$ mass using the simulation [18], including the efficiency scale factors. The acceptance systematic uncertainty is dominated by the parton distribution function uncertainty, which we estimate to be 4% using the Martin-Roberts-Stirling-Thorne prescription [23]. In the mass range of interest, the acceptances are ≈ 34% for the ee and $\mu\mu$ channels and ≈ 18% for the $e\mu$ channel.

No events are found in the high-mass regions. This null result yields a 95% C.L. upper limit on the cross section as a function of $H^{\pm\pm}$ mass (Fig. 2). We calculate the limit using a Bayesian method [24] with a flat prior for the signal and Gaussian priors for background and acceptance uncertainties. Through comparison with the theoretical cross sections [25], we obtain mass limits of 133, 136, and 115 GeV/c^2, for exclusive $H_L^{\pm\pm}$ decays to ee, $\mu\mu$, and $e\mu$, respectively, and 113 GeV/c^2 for exclusive $H_R^{\pm\pm}$ decays to $\mu\mu$. Figure 3 shows these results in the mass-coupling plane, along with the current world limits.

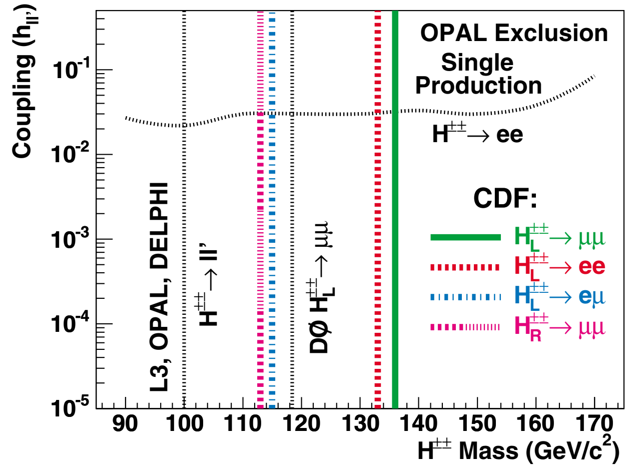

FIG. 3 (color). The doubly charged Higgs lower mass limits versus lepton coupling ($h_{ll'}$) from this analysis, assuming exclusive decay to a given dilepton pair. Our limits are valid for $h_{ll'} > 10^{-5}$. Previous limits [8,11,12] are also shown.

In summary, we have performed an inclusive search for doubly charged resonances in same-sign ee data with $m_{ee} > 100\ \mathrm{GeV}/c^2$, and same-sign $\mu\mu$ and $e\mu$ data with $m_{ll'} > 80\ \mathrm{GeV}/c^2$. We have found no evidence for new doubly charged resonances and have significantly extended the existing mass limits on doubly charged Higgs bosons decaying exclusively to ee ($m_{H_L^{\pm\pm}} > 133\ \mathrm{GeV}/c^2$), $\mu\mu$ ($m_{H_L^{\pm\pm}} > 136\ \mathrm{GeV}/c^2$ and $m_{H_R^{\pm\pm}} > 113\ \mathrm{GeV}/c^2$), or $e\mu$ ($m_{H_L^{\pm\pm}} > 115\ \mathrm{GeV}/c^2$) final states.

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