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Structure and magnetic properties of sputtered hard/soft multilayer magnets

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The films with $\text{HM}=(\text{Pr,Dy})(\text{Fe,Co,Nb,B})_{5.5}$ and $\text{SM}=\text{Fe}$, FeCo were prepared by sputtering and subsequent heat treatment. The coercivity of Ti-buffered $(\text{Pr,Dy})(\text{Fe,Co,Nb,B})_{5.5}$ single-layer film with 320 nm thickness is as large as 18.8 kOe at room temperature. X-ray diffraction results reveal that the $\text{Pr}_2\text{Fe}_{14}\text{B}$ -type phase is randomly oriented in almost all the multilayer films. For the multilayers of $\text{Ti}(30\text{ nm})/[\text{HM}(16\text{ nm})\text{Fe}(x\text{ nm})]\times 20/\text{Ti}(30\text{ nm})/\text{Si}(\text{substrate})$, the remanence increases and the coercivity decreases with the addition of Fe content, in comparison with the results of the single-layer film and the maximum energy product of 14.8 MGOe is obtained at $x=3.0$. A noticeable shoulder on the demagnetization curve is observed at low temperatures. When $\text{Fe}_{65}\text{Co}_{35}$ is used as the SM component rather than Fe, similar results are found. The enhancement of the magnetic properties in the nanocomposite multilayer films is explained by means of the exchange coupling between the SM and/or HM nanograins of the intra- and interlayers. © 2003 American Institute of Physics. [DOI: 10.1063/1.1558663]

I. INTRODUCTION

Recently, some studies on the exchange coupling were carried out for nanostructured CoSm/FeCo and PrCo/Co multilayers prepared by sputtering and subsequent heat treatment.^{1,2} Magnetic properties of exchange-coupled $\alpha\text{-Fe}/\text{Nd-Fe-B}$ multilayer magnets were investigated by Shindo and Ishizone³ and the observations for $\text{Nd-Fe-B}/\text{Fe}/\text{Nd-Fe-B}$ trilayers were reported by Parhofer *et al.*^{4,5} and Yang *et al.*^{3,6} In our recent work,⁷ we investigated the magnetic properties of nanocomposite multilayer magnets of $(\text{Nd,Dy})(\text{Fe,Co,Nb,B})_{5.5}/\text{M}(\text{M}=\text{Co,Fe}_{65}\text{Co}_{35})$ multilayer on Ti-buffered Si substrates prepared by sputtering and subsequent heat treatments. In comparison with the case of the single layer, the remanence of the multilayer magnets increases noticeably. Although the behavior at room temperature of the Pr-based $\text{R}_2\text{Fe}_{14}\text{B}$ -type alloy is somewhat similar to that of the Nd-based counterpart, the former is still distinctive from the latter, which stimulates us to study the structural and magnetic properties of nanocomposite $(\text{Pr,Dy})(\text{Fe,Co,Nb,B})_{5.5}/(\text{Fe or Fe}_{65}\text{Co}_{35})$ multilayer magnets synthesized by sputtering and subsequent annealing.

II. EXPERIMENT

$(\text{Pr,Dy})(\text{Fe,Co,Nb,B})_{5.5}/\text{Fe}$ (or $\text{Fe}_{65}\text{Co}_{35}$) thin films were prepared with a multiple-gun dc- and rf-sputtering system by depositing the hard magnetic (HM) $(\text{Pr}_{0.9}\text{Dy}_{0.1})(\text{Fe}_{0.77}\text{Co}_{0.12}\text{Nb}_{0.03}\text{B}_{0.08})_{5.5}$ alloy and the soft magnetic (SM) Fe or $\text{Fe}_{65}\text{Co}_{35}$ targets onto silicon substrate, covered with a 30 nm Ti buffer. The alloy targets were home-made by sintering powdered compacts and others were commercial products. Purities of all the targets were higher than 99.9%. The base pressure of the sputtering system was $2-3\times 10^{-7}$ Torr, and the Ar pressure during the sputtering was 5×10^{-3} Torr. The thickness of the films was measured by weighing the mass of the films. The as-deposited films were annealed in a furnace with a vacuum of 2×10^{-7} Torr. The crystalline structure of the phases in the films was identified by x-ray diffraction (XRD) with $\text{Cu K}\alpha$ radiation. Magnetic properties of the films were measured by an alternating gradient force magnetometer and a superconducting quantum interference device magnetometer. The hysteresis loops, as well as the values for the magnetic properties, were recorded without the demagnetizing correction.

III. RESULTS AND DISCUSSION

For comparison with the results of multilayers, initially, a hard phase single-layer film with composition of $\text{Ti}(30\text{ nm})/\text{HM}(320\text{ nm})/\text{Ti}(30\text{ nm})/(\text{Si substrate})$ was investigated.

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FIG. 1. XRD patterns for the Ti (30 nm)/HM(320 nm)/Ti(30 nm)/(Si substrate) single layer and Ti(30 nm)/[HM(16 nm)Fe(9 nm)] $\times$ 20/Ti(30 nm)/(Si substrate) multilayer films annealed at 600 °C for 5 min.

Similar to the case of (Nd,Dy)(Fe,Co,Nb,B)_{5.5} single-layer films,⁷ the as-deposited (Pr,Dy)(Fe,Co,Nb,B)_{5.5} single-layer films are amorphous. After annealing at 600 °C for 5 min, the main phase of the films is of Pr₂Fe₁₄B type, accompanied by a Pr-rich phase and some Pr₂O₃. XRD patterns for the Ti(30 nm)/HM (320 nm)/Ti(30 nm)/(Si substrate) single-layer film and Ti(30 nm)/[HM (16 nm) Fe(9 nm)] $\times$ 20/Ti(30 nm)/(Si substrate) multilayer film annealed at 600 °C for 5 min are shown in Fig. 1. In comparison with the result of the single-layer film, it is clear that adding a 10 nm thick Fe layer in the system with a 16 nm thick hard-phase layer results in the appearance of a large amount of α -Fe and the disappearance of the Pr-rich phase in the sample after annealing. Almost all of the XRD peaks of the films correspond to randomly oriented Pr₂Fe₁₄B-type phase, except for α -Fe.

Figure 2 gives hysteresis loops at room temperature for the single-layer and multilayer films, whose structures are shown in Fig. 1. The magnetic properties of $JH_c = 18.8$ kOe, $4\pi M_r = 6.1$ kG, and $(BH)_{\max} = 8.5$ MGOe are achieved for the single-layer film. For the Ti(30 nm)/[HM(16 nm)Fe(9 nm)] $\times$ 20/Ti(30 nm)/(Si substrate) multilayer film, although the coercivity decreases, the remanence is enhanced greatly due to the effective exchange coupling between the nanograins of SM and HM phases in the multilayer film.

To understand the effect of the thickness of the Fe layer, the magnetic properties at room temperature of the thin films of Ti(30 nm)/[HM(16 nm)Fe(x nm)] $\times$ 20/Ti(30 nm)/(Si substrate) multilayers annealed at 600 °C for 5 min are given in Fig. 3. Compared to the result of the single-layer film ($x=0$), by increasing the thickness of the Fe layer, the intrinsic coercivity decreases, the remanence increases clearly, and the energy products reach the maximum of 14.8 MGOe for $x=3$.

Figure 4 shows a hysteresis loop at room temperature and a demagnetization curve at 10 K for Ti(30 nm)/[HM(13

FIG. 2. Hysteresis loops at room temperature for the single layer and multilayer films shown in Fig. 1.

nm)Fe(5 nm)] $\times$ 20/Ti(30 nm)/(Si substrate) multilayer film annealed at 625 °C for 1 min. The magnetic properties, $4\pi M_r = 9.5$ kG, $JH_c = 9.9$ kOe, and $(BH)_{\max} = 14.7$ MGOe are achieved in the multilayer film at room temperature. In comparison with the result of the single-layer film mentioned herein, the remanence and the maximum energy product increase clearly due to the exchange coupling between the SM and/or HM nanograins of the intra- and interlayers. Coercivity up to 37.7 kOe is observed when measured at 10 K for Ti(30 nm)/[HM(13 nm)Fe(5 nm)] $\times$ 20/Ti(30 nm)/(Si substrate). However, a noticeable shoulder on the demagnetiza-

FIG. 3. Magnetic properties at room temperature of the thin films of Ti(30 nm)/[HM(16 nm)Fe(x nm)] $\times$ 20/Ti(30 nm)/(Si substrate) multilayer films annealed at 600 °C for 5 min.

FIG. 4. Hysteresis loop at room temperature and demagnetization curve at 10 K for Ti(30 nm)/[HM(13 nm)Fe(5 nm)] $\times$ 20/Ti(30 nm)/(Si substrate) multilayer film annealed at 625 °C for 1 min.

tion curve is observed because the enhancement of the anisotropy of the hard phase at low temperatures results in a significantly decrease of exchange length. This means that the exchange fields from the HM grains cannot completely cover the SM grains in the nearest-neighbor regions of the latter, and thus the magnetization reversal easily nucleates in the free region of the SM grains.

Because the saturation magnetization of Fe₆₅Co₃₅ alloy is higher than that of pure Fe, the Fe₆₅Co₃₅ alloy is also chosen as the SM layer component of the multilayer films. Room-temperature hysteresis loops of Ti(30 nm) buffered multilayer films with different thicknesses of HM and SM Fe₆₅Co₃₅ layers on a Si substrate annealed at 600 °C for 5 min are given in Fig. 5, where sample 1: Ti(30 nm)/[HM(18 nm)FeCo(9 nm)] $\times$ 20/Ti(30 nm)/(Si substrate), sample 2: Ti(30 nm)/[HM(15 nm)FeCo(7.5 nm)] $\times$ 20/Ti(30 nm)/(Si substrate), and sample 3: Ti(30 nm)/[HM(20 nm)FeCo(10 nm)] $\times$ 20/Ti(30 nm)/(Si substrate), respectively. It can be seen that all of the remanences and coercive forces of the multilayer films are higher and lower than that of the single layer film, respectively. Although the ratio of the thicknesses of hard and soft layers is same for the three samples, their effects on magnetic properties of the multilayer films are very different. Particularly in sample 3, the hard and soft layers are so thick that the squareness of the hysteresis loop is relatively small. Because the ratio of the thicknesses of the

FIG. 5. Room-temperature hysteresis loops of Ti(30 nm) buffered multilayer films with different thicknesses of HM and SM=Fe₆₅Co₃₅ layers on Si substrate annealed at 600 °C for 5 min, where sample 1: Ti(30 nm)/[HM(18 nm)FeCo(9 nm)] $\times$ 20/Ti(30 nm)/(Si substrate), sample 2: Ti(30 nm)/[HM(15 nm)FeCo(7.5 nm)] $\times$ 20/Ti(30 nm)/(Si substrate), and sample 3: Ti(30 nm)/[HM(20 nm)FeCo(10 nm)] $\times$ 20/Ti(30 nm)/(Si substrate).

hard and soft layers is constrained, the mutually dispersed soft and hard phases formed in the multilayer film after annealing may result in incomplete exchange coupling between some hard and soft nanograins in the multilayer films. The maximum energy product of 14.6 MGOe is achieved in Ti(30 nm)/[HM(18 nm)Fe(9 nm)] $\times$ 20/Ti(30 nm)/(Si substrate) multilayer film. It is concluded that the proper thickness of hard and soft layers is necessary for the multilayer films to have the complete exchange coupling.

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