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V. Khachatryan
Yerevan Physics Institute

Kenneth A. Bloom
University of Nebraska - Lincoln, kbloom2@unl.edu

S. Bose
University of Nebraska - Lincoln, sbose2@unl.edu

J. Butt
University of Nebraska - Lincoln

Daniel R. Claes
University of Nebraska - Lincoln, dclaes@unl.edu

See next page for additional authors

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Authors

V. Khachatryan, Kenneth A. Bloom, S. Bose, J. Butt, Daniel R. Claes, Aaron Dominguez, Michael Eads, J. Keller, T. Kelly, Ilya Kravchenko, J. Lazo-Flores, Carl Lundstedt, H. Malbouisson, Sudhir Malik, Gregory Snow, and CMS Collaboration



Transverse-Momentum and Pseudorapidity Distributions of Charged Hadrons in pp Collisions at $\sqrt{s} = 7$ TeV

V. Khachatryan *et al.**

(CMS Collaboration)

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Charged-hadron transverse-momentum and pseudorapidity distributions in proton-proton collisions at $\sqrt{s} = 7$ TeV are measured with the inner tracking system of the CMS detector at the LHC. The charged-hadron yield is obtained by counting the number of reconstructed hits, hit pairs, and fully reconstructed charged-particle tracks. The combination of the three methods gives a charged-particle multiplicity per unit of pseudorapidity $dN_{\text{ch}}/d\eta|_{|\eta|<0.5} = 5.78 \pm 0.01(\text{stat}) \pm 0.23(\text{syst})$ for non-single-diffractive events, higher than predicted by commonly used models. The relative increase in charged-particle multiplicity from $\sqrt{s} = 0.9$ to 7 TeV is $[66.1 \pm 1.0(\text{stat}) \pm 4.2(\text{syst})]\%$. The mean transverse momentum is measured to be $0.545 \pm 0.005(\text{stat}) \pm 0.015(\text{syst})$ GeV/ c . The results are compared with similar measurements at lower energies.

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Introduction.—Measurements of particle yields and kinematic distributions are an essential first step in exploring a new energy regime of particle collisions. Such studies contribute to our understanding of the physics of hadron production, including the relative roles of soft and hard scattering contributions, and help construct a solid foundation for other investigations. In the complicated environment of LHC pp collisions [1], firm knowledge of the rates and distributions of inclusive particle production is needed to distinguish rare signal events from the much larger backgrounds of soft hadronic interactions. They will also serve as points of reference for the measurement of nuclear-medium effects in Pb-Pb collisions in the LHC heavy ion program.

The bulk of the particles produced in pp collisions arise from soft interactions, which are modeled only phenomenologically. Experimental results provide the critical guidance for tuning these widely used models and event generators. Soft collisions are commonly classified as elastic scattering, inelastic single-diffractive (SD) dissociation, double-diffractive (DD) dissociation, and inelastic nondiffractive (ND) scattering [2]. (Double-Pomeron exchange is treated as DD in this Letter.) All results presented here refer to inelastic non-single-diffractive (NSD) interactions, and are based on an event selection that retains a large fraction of the ND and DD events, while disfavoring SD events.

The measurements focus on transverse-momentum p_T and pseudorapidity η distributions. The pseudorapidity,

commonly used to characterize the direction of particle emission, is defined as $\eta = -\ln \tan(\theta/2)$, where θ is the polar angle of the direction of the particle with respect to the anticlockwise beam direction. The count of primary charged hadrons N_{ch} is defined to include decay products of particles with proper lifetimes less than 1 cm. Products of secondary interactions are excluded, and a percent-level correction is applied for prompt leptons. The measurements reported here are of $dN_{\text{ch}}/d\eta$ and dN_{ch}/dp_T in the pseudorapidity range $|\eta| < 2.4$ and closely follow our previous analysis of minimum-bias data at lower center-of-mass energies of $\sqrt{s} = 0.9$ and 2.36 TeV as reported in Ref. [3].

The data for this study are drawn from an integrated luminosity of $1.1 \mu\text{b}^{-1}$ recorded with the Compact Muon Solenoid (CMS) experiment [4] on 30 March 2010, during the first hour of the LHC operation at $\sqrt{s} = 7$ TeV. These results are the highest center-of-mass energy measurements of the $dN_{\text{ch}}/d\eta$ and dN_{ch}/dp_T distributions conducted at a particle collider so far and complement the other recent measurements of the ALICE experiment at 7 TeV [5].

Experimental methods.—A detailed description of the CMS experiment can be found in Ref. [4]. The detectors used for the present analysis are the pixel and silicon-strip tracker, covering the region $|\eta| < 2.5$ and immersed in a 3.8 T axial magnetic field. The pixel tracker consists of three barrel layers and two end-cap disks at each barrel end. The forward calorimeter (HF), which covers the region $2.9 < |\eta| < 5.2$, was also used for event selection. The detailed Monte Carlo (MC) simulation of the CMS detector response is based on GEANT4 [6].

The event selection and analysis methods in this Letter are identical to those used in Ref. [3], where more details can be found. The inelastic pp collision rate was about 50 Hz. At these rates, the fraction of events in the data,

*Full author list given at the end of the article.

TABLE I. Numbers of events passing the selection cuts. The selection criteria are applied in sequence, i.e., each line includes the selection from the previous ones.

Selection	Number of events
Colliding bunches + one BSC signal	68 512
Reconstructed PV	61 551
HF coincidence	55 113
Beam-halo rejection	55 104
Other beam-background rejection	55 100

where two or more minimum-bias collisions occurred in the same bunch crossing, is estimated to be less than 0.3% and was neglected. Any hit in the beam scintillator counters (BSC, $3.23 < |\eta| < 4.65$) coinciding with colliding proton bunches was used for triggering the data acquisition. A sample mostly populated with NSD events was selected by requiring a primary vertex (PV) to be reconstructed with the tracker, together with at least one HF tower in each end with more than 3 GeV total energy. Beam-halo and other beam-background events were rejected as described in Ref. [3]. The remaining fraction of background events in the data was found to be less than 2×10^{-5} . The numbers of events satisfying the selection criteria are listed in Table I.

The event selection efficiency was estimated with simulated events using the PYTHIA [7,8] and PHOJET [9,10] event generators. The relative event fractions of SD, DD, and ND processes and their respective event selection efficiencies are listed in Table II. The fraction of diffractive events is predicted by the models to decrease as a function of collision energy, while the selection efficiency increases. At $\sqrt{s} = 7$ TeV, the fraction of SD (DD) events in the selected data sample, estimated with PYTHIA and PHOJET, are 6.8% (5.8%) and 5.0% (3.8%), respectively, somewhat higher than at $\sqrt{s} = 0.9$ and 2.36 TeV [3]. With PYTHIA, the overall correction for the selection efficiency of NSD processes and for the fraction of SD events remaining in the data sample lowers the measured charged-particle multiplicity by 6% compared with the uncorrected distribution.

The $dN_{\text{ch}}/d\eta$ distributions were obtained, as in Ref. [3], with three methods, based on counting the following quantities: (i) reconstructed clusters in the barrel part of the

pixel detector; (ii) pixel tracklets composed of pairs of clusters in different pixel barrel layers; and (iii) tracks reconstructed in the full tracker volume. The third method also allows a measurement of the dN_{ch}/dp_T distribution. All three methods rely on the reconstruction of a PV [11]. The PV reconstruction efficiency was found to be 98.3% (98.0%) in data (MC), evaluated after all other event selection cuts. In case of multiple PV candidates, the vertex with the largest track multiplicity was chosen. The three methods are sensitive to the measurement of particles down to p_T values of about 30, 50, and 100 MeV/ c , respectively. Only 0.5, 1.5, and 5% of all charged particles are estimated to be produced below these p_T values, respectively, and these fractions were corrected for.

The measurements were corrected for the geometrical acceptance ($\approx 2\%$), efficiency ($\approx 5\%–10\%$), fake ($< 1\%$) and duplicate tracks ($< 0.5\%$), low- p_T particles curling in the axial magnetic field ($< 1\%$), decay products of long-lived hadrons ($< 2\%$) and photon conversions ($< 1\%$), and inelastic hadronic interactions in the detector material ($\approx 1\%–2\%$), where the size of the corrections in parentheses refers to the tracking method. The PYTHIA parameter set from Ref. [8] was chosen to determine the corrections, because it reproduces the $dN_{\text{ch}}/d\eta$ and charged-particle multiplicity distributions, as well as other control distributions at 7 TeV, better than other available tuning parameter sets. Although the corrections do not depend significantly on the model used, it is indeed important that the simulated data set contains a sufficient number of high-multiplicity events to determine these corrections with the desired accuracy.

Results.—For the measurement of the dN_{ch}/dp_T distribution, charged-particle tracks with p_T in excess of 0.1 GeV/ c were used in 12 different $|\eta|$ bins, from 0 to 2.4. The average charged-hadron yields in NSD events are shown in Fig. 1 as a function of p_T and $|\eta|$. The Tsallis parametrization [12–14],

$$E \frac{d^3 N_{\text{ch}}}{dp^3} = \frac{1}{2\pi p_T} \frac{E}{p} \frac{d^2 N_{\text{ch}}}{d\eta dp_T} = C \frac{dN_{\text{ch}}}{dy} \left(1 + \frac{E_T}{nT}\right)^{-n}, \quad (1)$$

where $y = 0.5 \ln[(E + p_z)/(E - p_z)]$, $E_T = \sqrt{m^2 + p_T^2} - m$, and m is the charged pion mass, was fitted to the data. The p_T spectrum of charged hadrons, $1/(2\pi p_T) d^2 N_{\text{ch}}/d\eta dp_T$, measured in the range $|\eta| < 2.4$, is shown in Fig. 2 for data

TABLE II. Fractions of SD, DD, ND, and NSD processes obtained from the PYTHIA and PHOJET event generators before any selection, and the corresponding selection efficiencies determined from the MC simulation.

	PYTHIA		PHOJET	
	Fractions	Selection efficiencies	Fractions	Selection efficiencies
SD	19.2%	26.7%	13.8%	30.7%
DD	12.9%	33.6%	6.6%	48.3%
ND	67.9%	96.4%	79.6%	97.1%
NSD	80.8%	86.3%	86.2%	93.4%

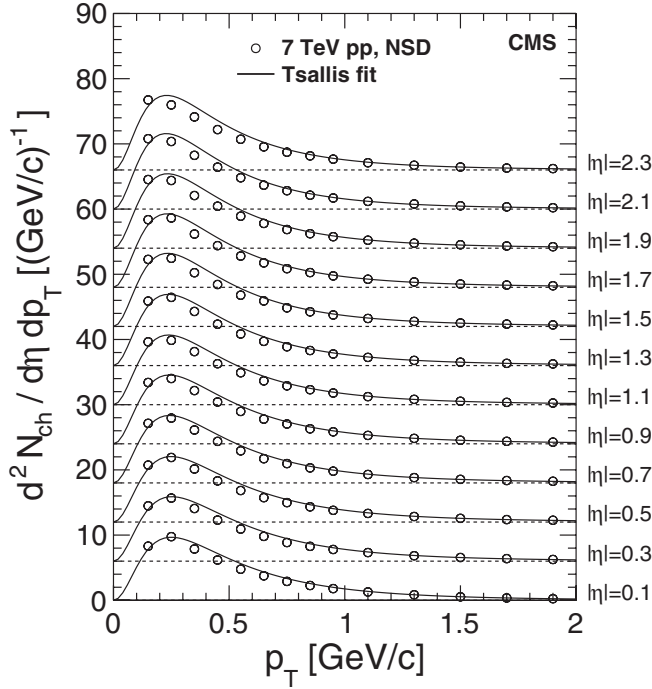


FIG. 1. Differential yield of charged hadrons in the range $|\eta| < 2.4$ in 0.2-unit-wide bins of $|\eta|$ in NSD events. The solid curves represent fits of Eq. (1) to the data. The measurements with increasing η are successively shifted by six units along the vertical axis.

at 0.9, 2.36, and 7 TeV. The high- p_T reach of the data is limited by the increase of systematic uncertainties with p_T . The fit to the data [Eq. (1)] is mainly used for extrapolations to $p_T = 0$, but is not expected to give a good description of the data in all η bins with only two parameters. The parameter T and the exponent n were found to be $T = 0.145 \pm 0.005(\text{syst})$ GeV and $n = 6.6 \pm 0.2(\text{syst})$. The average p_T , calculated from a combination of the measured data points and the low- and high- p_T contributions as determined from the fit, is $\langle p_T \rangle = 0.545 \pm 0.005(\text{stat}) \pm 0.015(\text{syst})$ GeV/c.

Experimental uncertainties related to the trigger and event selection are common to all the analysis methods. The uncertainty related to the presence of SD (DD) events in the final sample was estimated to be 1.4% (1.1%), based on consistency checks between data and simulation for diffractive event candidates. The total event selection uncertainty, which also includes the selection efficiency of the BSC and HF, was found to be 3.5%. Based on studies similar to those presented in Ref. [3], additional 3% and 2% uncertainties were assigned to the tracklet and track reconstruction algorithm efficiencies, respectively. Corrections at the percent level were applied to the final results to extrapolate to $p_T = 0$. The uncertainty on these extrapolation corrections was found to be less than 1%. All other uncertainties are identical to those listed in Ref. [3]. The $dN_{\text{ch}}/d\eta$ measurements were repeated on a separate

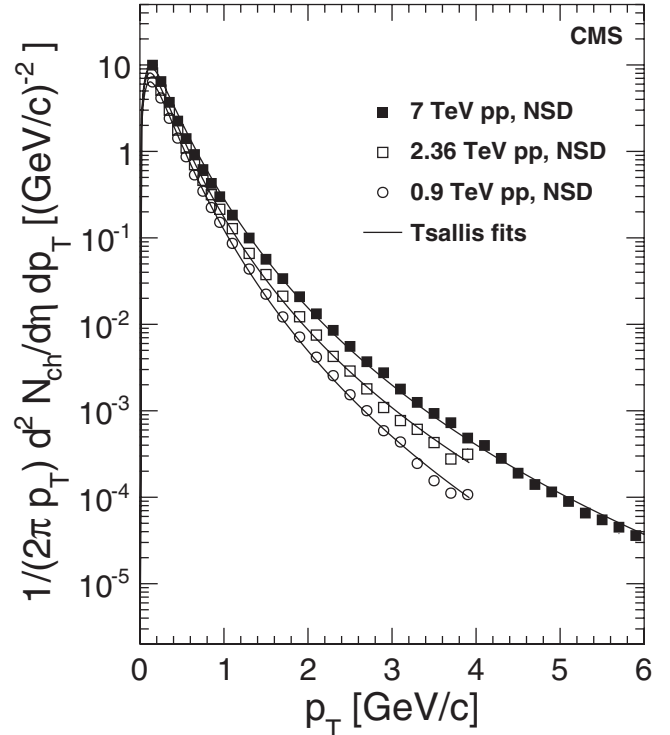


FIG. 2. Charged-hadron yield in the range $|\eta| < 2.4$ in NSD events as a function of p_T ; the systematic uncertainties are smaller than the symbols. The measurements at $\sqrt{s} = 0.9$ and 2.36 TeV [3] are also shown. The solid lines represent fits of Eq. (1) to the data.

data sample without any magnetic field, for which almost no p_T extrapolation is needed, and gave results consistent within 1.5%. The final systematic uncertainties for the pixel counting, tracklet, and track methods were found to be 5.7%, 4.6%, and 4.3%, respectively, and are strongly correlated.

For the $dN_{\text{ch}}/d\eta$ measurements, the results for the three individual layers within the cluster-counting method were found to be consistent within 1.2% and were combined. The three layer pairs in the pixel-tracklet method provided results that agreed within 0.6% and were also combined. Finally, the results from the three different measurement methods, which agree with the combined result within 1% to 4% depending on η , were averaged. The final $dN_{\text{ch}}/d\eta$ distributions are shown in Fig. 3 for $\sqrt{s} = 0.9, 2.36,$ and 7 TeV. The CMS results are compared with measurements made by other experiments. In the ATLAS Collaboration analysis [15], events and particles were selected in a different region of phase space, which makes a direct comparison difficult. Their results are therefore not included in the figure.

The results can also be compared to earlier experiments as a function of \sqrt{s} . The energy dependence of the average charged hadron p_T can be described by a quadratic function of $\ln s$ [16]. As shown in Fig. 4, the present measure-

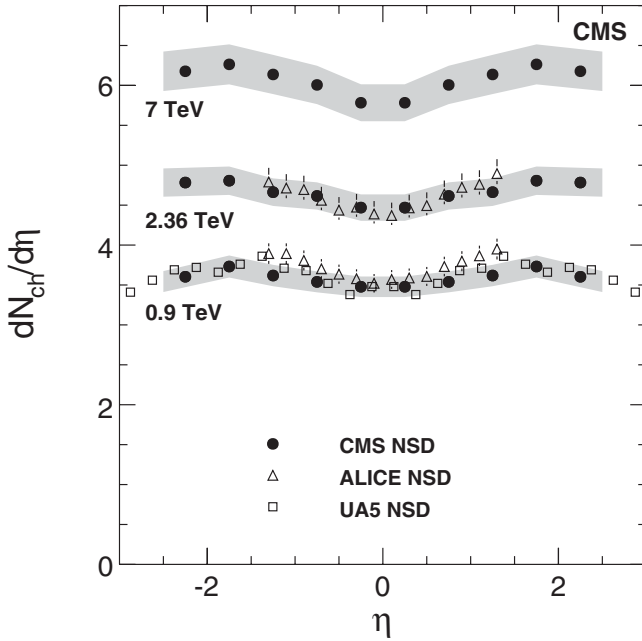


FIG. 3. Distributions of $dN_{\text{ch}}/d\eta$, averaged over the three measurement methods and compared with data from UA5 [23] ($p\bar{p}$, with statistical errors only) and ALICE [24] (with systematic uncertainties). The shaded band shows systematic uncertainties of the CMS data. The CMS and UA5 data are averaged over negative and positive values of η .

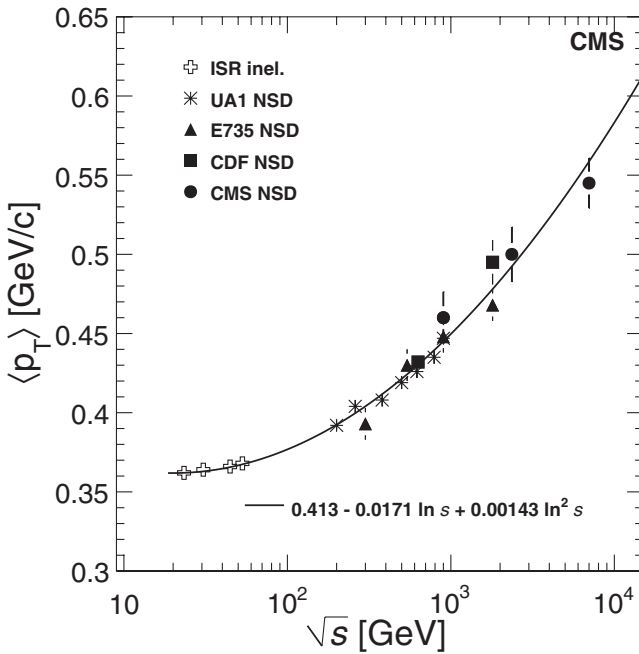


FIG. 4. Average p_T of charged hadrons as a function of the center-of-mass energy. The CMS measurements are for $|\eta| < 2.4$. Also shown are measurements from the ISR [25] (pp), E735 [26] ($p\bar{p}$), and CDF [27] ($p\bar{p}$) for $|\eta| < 0.5$, and from UA1 [16] ($p\bar{p}$) for $|\eta| < 2.5$. The solid line is a fit of the functional form $\langle p_T \rangle = 0.413 - 0.0171 \ln s + 0.00143 \ln^2 s$ to the data. The error bars on the CMS data include the systematic uncertainties.

ment follows this trend. The choice of the $|\eta|$ interval can influence the average p_T value by a few percent.

For $|\eta| < 0.5$, the average charged multiplicity density is $dN_{\text{ch}}/d\eta = 5.78 \pm 0.01(\text{stat}) \pm 0.23(\text{syst})$ for NSD events. The \sqrt{s} dependence of the measured $dN_{\text{ch}}/d\eta|_{\eta=0}$ is shown in Fig. 5, which includes data from various other experiments. The $dN_{\text{ch}}/d\eta$ results reported here show a rather steep increase between 0.9 and 7 TeV, which is measured to be $[66.1 \pm 1.0(\text{stat}) \pm 4.2(\text{syst})]\%$. Using a somewhat different event selection, the ALICE Collaboration has found a similar increase of $[57.6 \pm 0.4(\text{stat})^{+3.6}_{-1.8}(\text{syst})]\%$ [5]. The measured charged-particle multiplicity is accurate enough to distinguish among most sets of event-generator tuning parameter values and various models. The measured value at 7 TeV significantly exceeds the prediction of 4.57 from PHOJET [9,10], and the predictions of 3.99, 4.18, and 4.34 from the DW [17], PROQ20 [18], and Perugia0 [19] tuning parameter values of PYTHIA, respectively, while it is closer to the prediction of 5.48 from the PYTHIA parameter set from Ref. [8] and to the recent model predictions of 5.58 and 5.78 from Refs. [20,21]. The measured excess of the number of charged hadrons with respect to the event generators is independent of η and concentrated in the $p_T < 1 \text{ GeV}/c$

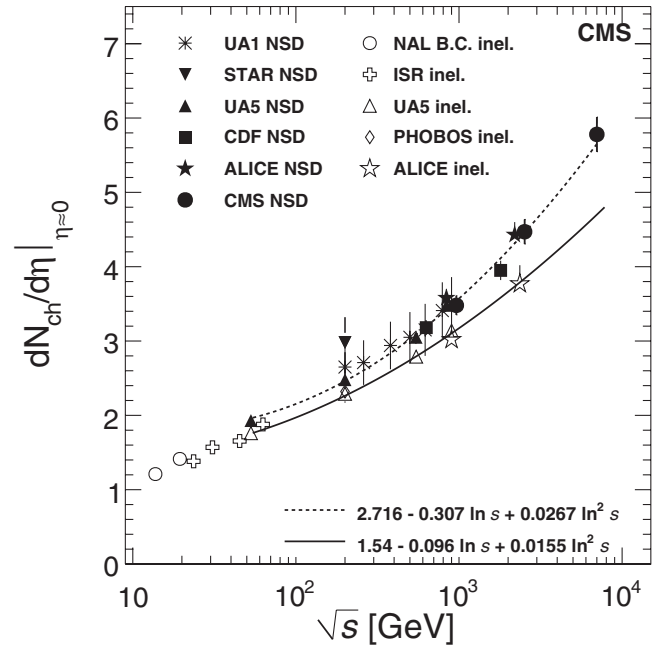


FIG. 5. Average value of $dN_{\text{ch}}/d\eta$ in the central η region as a function of center-of-mass energy in pp and $p\bar{p}$ collisions. Also shown are NSD and inelastic measurements from the NAL Bubble Chamber [28] ($p\bar{p}$), ISR [29] (pp), UA1 [16] ($p\bar{p}$), UA5 [23] ($p\bar{p}$), CDF [30] ($p\bar{p}$), STAR [31] (pp), PHOBOS [32] (pp), and ALICE [24] (pp). The curves are second-order polynomial fits for the inelastic (solid) and NSD event selections (dashed). The error bars include systematic uncertainties, when available. Data points at 0.9 and 2.36 TeV are slightly displaced horizontally for visibility.

range. These differences indicate the need for a continued model development and simulation tuning. Work on updated event generators based on LHC data is currently under way.

Summary.—Charged-hadron transverse-momentum and pseudorapidity distributions have been measured in proton-proton collisions at $\sqrt{s} = 7$ TeV. The numerical values of the data presented in this Letter can be found in the HEPDATA database [22]. The combined result for the central pseudorapidity density, from three mutually consistent methods of measurement, is $dN_{\text{ch}}/d\eta|_{|\eta|<0.5} = 5.78 \pm 0.01(\text{stat}) \pm 0.23(\text{syst})$ for non-single-diffractive events. This value is higher than most predictions and provides new information to constrain ongoing improvements of soft particle production models and event generators. The mean transverse momentum has been measured to be $0.545 \pm 0.005(\text{stat}) \pm 0.015(\text{syst})$ GeV/ c . These studies are the first steps in the exploration of particle production at the new center-of-mass energy frontier, and contribute to the understanding of the dynamics in soft hadronic interactions.

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V. Khachatryan,¹ A. M. Sirunyan,¹ A. Tumasyan,¹ W. Adam,² T. Bergauer,² M. Dragicevic,² J. Erö,² C. Fabjan,² M. Friedl,² R. Frühwirth,² V. M. Ghete,² J. Hammer,^{2,b} S. Häsnel,² M. Hoch,² N. Hörmann,² J. Hrubec,² M. Jeitler,² G. Kasieczka,² W. Kiesenhofer,² M. Krammer,² D. Liko,² I. Mikulec,² M. Pernicka,² H. Rohringer,² R. Schöfbeck,²

J. Strauss,² A. Taurok,² F. Teischinger,² W. Waltenberger,² G. Walzel,² E. Widl,² C.-E. Wulz,² V. Mossolov,³ N. Shumeiko,³ J. Suarez Gonzalez,³ L. Benucci,⁴ L. Ceard,⁴ E. A. De Wolf,⁴ M. Hashemi,⁴ X. Janssen,⁴ T. Maes,⁴ L. Mucibello,⁴ S. Ochesanu,⁴ B. Roland,⁴ R. Rougny,⁴ M. Selvaggi,⁴ H. Van Haevermaet,⁴ P. Van Mechelen,⁴ N. Van Remortel,⁴ V. Adler,⁵ S. Beauceron,⁵ S. Blyweert,⁵ J. D'Hondt,⁵ O. Devroede,⁵ A. Kalogeropoulos,⁵ J. Maes,⁵ M. Maes,⁵ S. Tavernier,⁵ W. Van Doninck,⁵ P. Van Mulders,⁵ I. Villella,⁵ E. C. Chabert,⁶ O. Charaf,⁶ B. Clerbaux,⁶ G. De Lentdecker,⁶ V. Dero,⁶ A. P. R. Gay,⁶ G. H. Hammad,⁶ P. E. Marage,⁶ C. Vander Velde,⁶ P. Vanlaer,⁶ J. Wickens,⁶ S. Costantini,⁷ M. Grunewald,⁷ B. Klein,⁷ A. Marinov,⁷ D. Ryckbosch,⁷ F. Thyssen,⁷ M. Tytgat,⁷ L. Vanelderden,⁷ P. Verwilligen,⁷ S. Walsh,⁷ N. Zaganidis,⁷ S. Basegmez,⁸ G. Bruno,⁸ J. Caudron,⁸ J. De Favereau De Jeneret,⁸ C. Delaere,⁸ P. Demin,⁸ D. Favart,⁸ A. Giammanco,⁸ G. Grégoire,⁸ J. Hollar,⁸ V. Lemaitre,⁸ O. Militaru,⁸ S. Olyn,⁸ D. Pagano,⁸ A. Pin,⁸ K. Piotrkowski,^{8,b} L. Quertenmont,⁸ N. Schul,⁸ N. Belyi,⁹ T. Caeberts,⁹ E. Daubie,⁹ G. A. Alves,¹⁰ M. E. Pol,¹⁰ M. H. G. Souza,¹⁰ W. Carvalho,¹¹ E. M. Da Costa,¹¹ D. De Jesus Damiao,¹¹ C. De Oliveira Martins,¹¹ S. Fonseca De Souza,¹¹ L. Mundim,¹¹ V. Oguri,¹¹ A. Santoro,¹¹ S. M. Silva Do Amaral,¹¹ A. Sznajder,¹¹ F. Torres Da Silva De Araujo,¹¹ F. A. Dias,¹² M. A. F. Dias,¹² T. R. Fernandez Perez Tomei,¹² E. M. Gregores,¹² F. Marinho,¹² S. F. Novaes,¹² Sandra S. Padula,¹² N. Dargatzis,^{13,b} L. Dimitrov,¹³ V. Genchev,^{13,b} P. Iaydjiev,¹³ S. Piperov,¹³ S. Stoykova,¹³ G. Sultanov,¹³ R. Trayanov,¹³ I. Vankov,¹³ M. Dyulendarova,¹⁴ R. Hadjiiska,¹⁴ V. Kozhuharov,¹⁴ L. Litov,¹⁴ E. Marinova,¹⁴ M. Mateev,¹⁴ B. Pavlov,¹⁴ P. Petkov,¹⁴ J. G. Bian,¹⁵ G. M. Chen,¹⁵ H. S. Chen,¹⁵ C. H. Jiang,¹⁵ D. Liang,¹⁵ S. Liang,¹⁵ J. Wang,¹⁵ J. Wang,¹⁵ X. Wang,¹⁵ Z. Wang,¹⁵ M. Yang,¹⁵ J. Zang,¹⁵ Z. Zhang,¹⁵ Y. Ban,¹⁶ S. Guo,¹⁶ Z. Hu,¹⁶ Y. Mao,¹⁶ S. J. Qian,¹⁶ H. Teng,¹⁶ B. Zhu,¹⁶ A. Cabrera,¹⁷ C. A. Carrillo Montoya,¹⁷ B. Gomez Moreno,¹⁷ A. A. Ocampo Rios,¹⁷ A. F. Osorio Oliveros,¹⁷ J. C. Sanabria,¹⁷ N. Godinovic,¹⁸ D. Lelas,¹⁸ K. Lelas,¹⁸ R. Plestina,^{18,c} D. Polic,¹⁸ I. Puljak,¹⁸ Z. Antunovic,¹⁹ M. Dzelalija,¹⁹ V. Brigljevic,²⁰ S. Duric,²⁰ K. Kadija,²⁰ S. Morovic,²⁰ A. Attikis,²¹ R. Fereos,²¹ M. Galanti,²¹ J. Mousa,²¹ C. Nicolaou,²¹ A. Papadakis,²¹ F. Ptochos,²¹ P. A. Razis,²¹ H. Rykaczewski,²¹ D. Tsiakkouri,²¹ Z. Zinonos,²¹ M. Mahmoud,²² A. Hektor,²³ M. Kadastik,²³ K. Kannike,²³ M. Müntel,²³ M. Raidal,²³ L. Rebane,²³ V. Azzolini,²⁴ P. Eerola,²⁴ S. Czellar,²⁵ J. Härkönen,²⁵ A. Heikkinen,²⁵ V. Karimäki,²⁵ R. Kinnunen,²⁵ J. Klem,²⁵ M. J. Kortelainen,²⁵ T. Lampén,²⁵ K. Lassila-Perini,²⁵ S. Lehti,²⁵ T. Lindén,²⁵ P. Luukka,²⁵ T. Mäenpää,²⁵ E. Tuominen,²⁵ J. Tuominiemi,²⁵ E. Tuovinen,²⁵ D. Ungaro,²⁵ L. Wendland,²⁵ K. Banzuzi,²⁶ A. Korpela,²⁶ T. Tuuva,²⁶ D. Sillou,²⁷ M. Besancon,²⁸ M. Dejardin,²⁸ D. Denegri,²⁸ J. Descamps,²⁸ B. Fabbro,²⁸ J. L. Faure,²⁸ F. Ferri,²⁸ S. Ganjour,²⁸ F. X. Gentit,²⁸ A. Givernaud,²⁸ P. Gras,²⁸ G. Hamel de Monchenault,²⁸ P. Jarry,²⁸ E. Locci,²⁸ J. Malcles,²⁸ M. Marionneau,²⁸ L. Millischer,²⁸ J. Rander,²⁸ A. Rosowsky,²⁸ D. Rousseau,²⁸ M. Titov,²⁸ P. Verrecchia,²⁸ S. Baffioni,²⁹ L. Bianchini,²⁹ M. Bluj,^{29,d} C. Brouin,²⁹ P. Busson,²⁹ C. Charlot,²⁹ L. Dobrzynski,²⁹ S. Elgammal,²⁹ R. Granier de Cassagnac,²⁹ M. Hagnauer,²⁹ A. Kalinowski,²⁹ P. Miné,²⁹ P. Paganini,²⁹ D. Sabes,²⁹ Y. Sirois,²⁹ C. Thiebaux,²⁹ A. Zabi,²⁹ J.-L. Agram,³⁰ A. Besson,³⁰ D. Bloch,³⁰ D. Bodin,³⁰ J.-M. Brom,³⁰ M. Cardaci,³⁰ E. Conte,³⁰ F. Drouhin,³⁰ C. Ferro,³⁰ J.-C. Fontaine,³⁰ D. Gelé,³⁰ U. Goerlach,³⁰ S. Greder,³⁰ P. Juillot,³⁰ M. Karim,³⁰ A.-C. Le Bihan,³⁰ Y. Mikami,³⁰ J. Speck,³⁰ P. Van Hove,³⁰ F. Fassi,³¹ D. Mercier,³¹ C. Baty,³² N. Beaupere,³² M. Bedjidian,³² O. Bondu,³² G. Boudoul,³² D. Boumediene,³² H. Brun,³² N. Chanon,³² R. Chierici,³² D. Contardo,³² P. Depasse,³² H. El Mamouni,³² J. Fay,³² S. Gascon,³² B. Ille,³² T. Kurca,³² T. Le Grand,³² M. Lethuillier,³² L. Mirabito,³² S. Perries,³² S. Tosi,³² Y. Tschudi,³² P. Verdier,³² H. Xiao,³² V. Roinishvili,³³ G. Anagnostou,³⁴ M. Edelhoff,³⁴ L. Feld,³⁴ N. Heracleous,³⁴ O. Hindrichs,³⁴ R. Jussen,³⁴ K. Klein,³⁴ J. Merz,³⁴ N. Mohr,³⁴ A. Ostapchuk,³⁴ A. Perieanu,³⁴ F. Raupach,³⁴ J. Sammet,³⁴ S. Schael,³⁴ D. Sprenger,³⁴ H. Weber,³⁴ M. Weber,³⁴ B. Wittmer,³⁴ O. Actis,³⁵ M. Ata,³⁵ W. Bender,³⁵ P. Biallass,³⁵ M. Erdmann,³⁵ J. Frangenheim,³⁵ T. Hebbeker,³⁵ A. Hinzmann,³⁵ K. Hoepfner,³⁵ C. Hof,³⁵ M. Kirsch,³⁵ T. Klimkovich,³⁵ P. Kreuzer,^{35,b} D. Lanske,^{35,a} C. Magass,³⁵ M. Merschmeyer,³⁵ A. Meyer,³⁵ P. Papacz,³⁵ H. Pieta,³⁵ H. Reithler,³⁵ S. A. Schmitz,³⁵ L. Sonnenschein,³⁵ M. Sowa,³⁵ J. Steggemann,³⁵ D. Teyssier,³⁵ C. Zeidler,³⁵ M. Bontenackels,³⁶ M. Davids,³⁶ M. Duda,³⁶ G. Flügge,³⁶ H. Geenen,³⁶ M. Giffels,³⁶ W. Haj Ahmad,³⁶ D. Heydhausen,³⁶ T. Kress,³⁶ Y. Kuessel,³⁶ A. Linn,³⁶ A. Nowack,³⁶ L. Perchalla,³⁶ O. Pooth,³⁶ P. Sauerland,³⁶ A. Stahl,³⁶ M. Thomas,³⁶ D. Tornier,³⁶ M. H. Zoeller,³⁶ M. Aldaya Martin,³⁷ W. Behrenhoff,³⁷ U. Behrens,³⁷ M. Bergholz,³⁷ K. Borras,³⁷ A. Campbell,³⁷ E. Castro,³⁷ D. Dammann,³⁷ G. Eckerlin,³⁷ A. Flossdorf,³⁷ G. Flucke,³⁷ A. Geiser,³⁷ J. Hauk,³⁷ H. Jung,³⁷ M. Kasemann,³⁷ I. Katkov,³⁷ C. Kleinwort,³⁷ H. Kluge,³⁷ A. Knutsson,³⁷ E. Kuznetsova,³⁷ W. Lange,³⁷ W. Lohmann,³⁷ R. Mankel,³⁷ M. Marienfeld,³⁷ I.-A. Melzer-Pellmann,³⁷ A. B. Meyer,³⁷ J. Mnich,³⁷ A. Mussgiller,³⁷ J. Olzem,³⁷ A. Parenti,³⁷ A. Raspereza,³⁷ R. Schmidt,³⁷ T. Schoerner-Sadenius,³⁷ N. Sen,³⁷ M. Stein,³⁷ J. Tomaszewska,³⁷ D. Volyansky,³⁷ C. Wissing,³⁷ C. Autermann,³⁸

J. Draeger,³⁸ D. Eckstein,³⁸ H. Enderle,³⁸ U. Gebbert,³⁸ K. Kaschube,³⁸ G. Kaussen,³⁸ R. Klanner,³⁸ B. Mura,³⁸
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 J. Schwandt,³⁸ H. Stadie,³⁸ G. Steinbrück,³⁸ J. Thomsen,³⁸ R. Wolf,³⁸ J. Bauer,³⁹ V. Buege,³⁹ A. Cakir,³⁹
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 Th. Müller,³⁹ M. Niegel,³⁹ O. Oberst,³⁹ A. Oehler,³⁹ J. Ott,³⁹ T. Peiffer,³⁹ D. Piparo,³⁹ G. Quast,³⁹ K. Rabbertz,³⁹
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 G. Schott,³⁹ H. J. Simonis,³⁹ F. M. Stober,³⁹ D. Troendle,³⁹ J. Wagner-Kuhr,³⁹ M. Zeise,³⁹ V. Zhukov,^{39,e}
 E. B. Ziebarth,³⁹ G. Daskalakis,⁴⁰ T. Geralis,⁴⁰ A. Kyriakis,⁴⁰ D. Loukas,⁴⁰ I. Manolakis,⁴⁰ A. Markou,⁴⁰
 C. Markou,⁴⁰ C. Mavrommatis,⁴⁰ E. Petrakou,⁴⁰ L. Gouskos,⁴¹ P. Katsas,⁴¹ A. Panagiotou,^{41,b} I. Evangelou,⁴²
 P. Kokkas,⁴² N. Manthos,⁴² I. Papadopoulos,⁴² V. Patras,⁴² F. A. Triantis,⁴² A. Aranyi,⁴³ G. Bencze,⁴³ L. Boldizsar,⁴³
 G. Debreczeni,⁴³ C. Hajdu,^{43,b} D. Horvath,^{43,f} A. Kapusi,⁴³ K. Krajczar,⁴³ A. Laszlo,⁴³ F. Sikler,⁴³
 G. Vesztergombi,⁴³ N. Beni,⁴⁴ J. Molnar,⁴⁴ J. Palinkas,⁴⁴ Z. Szillasi,^{44,b} V. Veszpremi,⁴⁴ P. Raics,⁴⁵ Z. L. Trocsanyi,⁴⁵
 B. Ujvari,⁴⁵ S. Bansal,⁴⁶ S. B. Beri,⁴⁶ V. Bhatnagar,⁴⁶ M. Jindal,⁴⁶ M. Kaur,⁴⁶ J. M. Kohli,⁴⁶ M. Z. Mehta,⁴⁶
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 S. Bhattacharya,^{47,g} S. Chauhan,⁴⁷ B. C. Choudhary,⁴⁷ P. Gupta,⁴⁷ S. Jain,⁴⁷ S. Jain,⁴⁷ A. Kumar,⁴⁷ K. Ranjan,⁴⁷
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 P. Shukla,⁴⁸ P. Suggiseti,⁴⁸ T. Aziz,⁴⁹ M. Guchait,^{49,h} A. Gurtu,⁴⁹ M. Maity,⁴⁹ D. Majumder,⁴⁹ G. Majumder,⁴⁹
 K. Mazumdar,⁴⁹ G. B. Mohanty,⁴⁹ A. Saha,⁴⁹ K. Sudhakar,⁴⁹ N. Wickramage,⁴⁹ S. Banerjee,⁵⁰ S. Dugad,⁵⁰
 N. K. Mondal,⁵⁰ H. Arfaei,⁵¹ H. Bakhshiansohi,⁵¹ A. Fahim,⁵¹ A. Jafari,⁵¹ M. Mohammadi Najafabadi,⁵¹
 S. Paktinat Mehdiabadi,⁵¹ B. Safarzadeh,⁵¹ M. Zeinali,⁵¹ M. Abbrescia,^{52a,52b} L. Barbone,^{52a} A. Colaleo,^{52a}
 D. Creanza,^{52a,52c} N. De Filippis,^{52a} M. De Palma,^{52a,52b} A. Dimitrov,^{52a} F. Fedele,^{52a} L. Fiore,^{52a} G. Iaselli,^{52a,52c}
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 T. Rovelli,^{53a,53b} G. Siroli,^{53a,53b} R. Travaglini,^{53a,53b} S. Albergo,^{54a,54b} G. Cappello,^{54a,54b} M. Chiorboli,^{54a,54b}
 S. Costa,^{54a,54b} A. Tricomi,^{54a,54b} C. Tuve,^{54a} G. Barbagli,^{55a} G. Broccolo,^{55a,55b} V. Ciulli,^{55a,55b} C. Civinini,^{55a}
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 M. Meschini,^{55a} S. Paoletti,^{55a} G. Sguazzoni,^{55a} A. Tropiano,^{55a} L. Benussi,⁵⁶ S. Bianco,⁵⁶ S. Colafranceschi,⁵⁶
 F. Fabbri,⁵⁶ D. Piccolo,⁵⁶ P. Fabbriatore,⁵⁷ R. Musenich,⁵⁷ A. Benaglia,^{58a,58b} G. B. Cerati,^{58a,58b,b} F. De Guio,^{58a,58b}
 L. Di Matteo,^{58a,58b} A. Ghezzi,^{58a,58b,b} P. Govoni,^{58a,58b} M. Malberti,^{58a,58b,b} S. Malvezzi,^{58a} A. Martelli,^{58a,58b,c}
 A. Massironi,^{58a,58b} D. Menasce,^{58a} V. Miccio,^{58a,58b} L. Moroni,^{58a} P. Negri,^{58a,58b} M. Paganoni,^{58a,58b} D. Pedrini,^{58a}
 S. Ragazzi,^{58a,58b} N. Redaelli,^{58a} S. Sala,^{58a} R. Salerno,^{58a,58b} T. Tabarelli de Fatis,^{58a,58b} V. Tancini,^{58a,58b}
 S. Taroni,^{58a,58b} S. Buontempo,^{59a} A. Cimmino,^{59a,59b} A. De Cosa,^{59a,59b,b} M. De Gruttola,^{59a,59b,b} F. Fabozzi,^{59a}
 A. O. M. Iorio,^{59a} L. Lista,^{59a} P. Noli,^{59a,59b} P. Paolucci,^{59a} P. Azzi,^{60a} N. Bacchetta,^{60a} P. Bellan,^{60a,60b,b}
 M. Bellato,^{60a} M. Biasotto,^{60a} D. Bisello,^{60a,60b} R. Carlin,^{60a,60b} P. Checchia,^{60a} M. De Mattia,^{60a,60b} T. Dorigo,^{60a}
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 M. Margoni,^{60a,60b} G. Maron,^{60a} A. T. Meneguzzo,^{60a,60b} M. Nespolo,^{60a} L. Perrozzi,^{60a} N. Pozzobon,^{60a,60b}
 P. Ronchese,^{60a,60b} F. Simonetto,^{60a,60b} E. Torassa,^{60a} M. Tosi,^{60a,60b} A. Triossi,^{60a} S. Vanini,^{60a,60b} G. Zumerle,^{60a,60b}
 P. Baesso,^{61a,61b} U. Berzano,^{61a} C. Riccardi,^{61a,61b} P. Torre,^{61a,61b} P. Vitulo,^{61a,61b} C. Viviani,^{61a,61b} M. Biasini,^{62a,62b}
 G. M. Bilei,^{62a} B. Caponeri,^{62a,62b} L. Fanò,^{62a} P. Lariccia,^{62a,62b} A. Lucaroni,^{62a,62b} G. Mantovani,^{62a,62b}
 M. Menichelli,^{62a} A. Nappi,^{62a,62b} A. Santocchia,^{62a,62b} L. Servoli,^{62a} M. Valdata,^{62a} R. Volpe,^{62a,62b,b}
 P. Azzurri,^{63a,63c} G. Bagliesi,^{63a} J. Bernardini,^{63a,63b,b} T. Boccali,^{63a} R. Castaldi,^{63a} R. T. Dagnolo,^{63a,63c}
 R. Dell'Orso,^{63a} F. Fiori,^{63a,63b} L. Foà,^{63a,63c} A. Giassi,^{63a} A. Kraan,^{63a} F. Ligabue,^{63a,63c} T. Lomtadze,^{63a}
 L. Martini,^{63a} A. Messineo,^{63a,63b} F. Palla,^{63a} F. Palmonari,^{63a} G. Segneri,^{63a} A. T. Serban,^{63a} P. Spagnolo,^{63a,b}
 R. Tenchini,^{63a,b} G. Tonelli,^{63a,63b,b} A. Venturi,^{63a} P. G. Verdini,^{63a} L. Barone,^{64a,64b} F. Cavallari,^{64a,b}
 D. Del Re,^{64a,64b} E. Di Marco,^{64a,64b} M. Diemmoz,^{64a} D. Franci,^{64a,64b} M. Grassi,^{64a} E. Longo,^{64a,64b}
 G. Organtini,^{64a,64b} A. Palma,^{64a,64b} F. Pandolfi,^{64a,64b} R. Paramatti,^{64a,b} S. Rahatlou,^{64a,64b,b} N. Amapane,^{65a,65b}

R. Arcidiacono,^{65a,65b} S. Argiro,^{65a,65b} M. Arneodo,^{65a,65c} C. Biino,^{65a} C. Botta,^{65a,65b} N. Cartiglia,^{65a}
R. Castello,^{65a,65b} M. Costa,^{65a,65b} N. Demaria,^{65a} A. Graziano,^{65a,65b} C. Mariotti,^{65a} M. Marone,^{65a,65b} S. Maselli,^{65a}
E. Migliore,^{65a,65b} G. Mila,^{65a,65b} V. Monaco,^{65a,65b} M. Musich,^{65a,65b} M. M. Obertino,^{65a,65c} N. Pastrone,^{65a}
M. Pelliccioni,^{65a,65b,b} A. Romero,^{65a,65b} M. Ruspa,^{65a,65c} R. Sacchi,^{65a,65b} A. Solano,^{65a,65b} A. Staiano,^{65a}
D. Trocino,^{65a,65b} A. Vilela Pereira,^{65a,65b,b} F. Ambroglini,^{66a,66b} S. Belforte,^{66a} F. Cossutti,^{66a} G. Della Ricca,^{66a,66b}
B. Gobbo,^{66a} D. Montanino,^{66a} A. Penzo,^{66a} S. Chang,⁶⁷ J. Chung,⁶⁷ D. H. Kim,⁶⁷ G. N. Kim,⁶⁷ J. E. Kim,⁶⁷
D. J. Kong,⁶⁷ H. Park,⁶⁷ D. C. Son,⁶⁷ Zero Kim,⁶⁸ J. Y. Kim,⁶⁸ S. Song,⁶⁸ B. Hong,⁶⁹ H. Kim,⁶⁹ J. H. Kim,⁶⁹
T. J. Kim,⁶⁹ K. S. Lee,⁶⁹ D. H. Moon,⁶⁹ S. K. Park,⁶⁹ H. B. Rhee,⁶⁹ K. S. Sim,⁶⁹ M. Choi,⁷⁰ S. Kang,⁷⁰ H. Kim,⁷⁰
C. Park,⁷⁰ I. C. Park,⁷⁰ S. Park,⁷⁰ S. Choi,⁷¹ Y. Choi,⁷¹ Y. K. Choi,⁷¹ J. Goh,⁷¹ J. Lee,⁷¹ S. Lee,⁷¹ H. Seo,⁷¹ I. Yu,⁷¹
M. Janulis,⁷² D. Martisiute,⁷² P. Petrov,⁷² T. Sabonis,⁷² H. Castilla Valdez,^{73,b} E. De La Cruz Burelo,⁷³
R. Lopez-Fernandez,⁷³ A. Sánchez Hernández,⁷³ L. M. Villaseñor-Cendejas,⁷³ S. Carrillo Moreno,⁷⁴
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J. Varela,^{82,b} H. K. Wöhri,⁸² I. Altsybeev,⁸³ I. Belotelov,⁸³ P. Bunin,⁸³ M. Finger,⁸³ M. Finger, Jr.,⁸³ I. Golutvin,⁸³
A. Kamenev,⁸³ V. Karjavin,⁸³ G. Kozlov,⁸³ A. Lanev,⁸³ P. Moisezenz,⁸³ V. Palichik,⁸³ V. Perelygin,⁸³ S. Shmatov,⁸³
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I. Smirnov,⁸⁴ V. Sulimov,⁸⁴ L. Uvarov,⁸⁴ S. Vavilov,⁸⁴ A. Vorobyev,⁸⁴ Yu. Andreev,⁸⁵ S. Gninenko,⁸⁵ N. Golubev,⁸⁵
M. Kirsanov,⁸⁵ N. Krasnikov,⁸⁵ V. Matveev,⁸⁵ A. Pashenkov,⁸⁵ A. Toropin,⁸⁵ S. Troitsky,⁸⁵ V. Epshteyn,⁸⁶
V. Gavrilov,⁸⁶ N. Ilina,⁸⁶ V. Kaftanov,^{86,a} M. Kossov,^{86,b} A. Krokhotin,⁸⁶ S. Kuleshov,⁸⁶ A. Oulianov,⁸⁶
G. Safronov,⁸⁶ S. Semenov,⁸⁶ I. Shreyber,⁸⁶ V. Stolin,⁸⁶ E. Vlasov,⁸⁶ A. Zhokin,⁸⁶ E. Boos,⁸⁷ M. Dubinin,^{87,i}
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A. Vinogradov,⁸⁸ I. Azhgirey,⁸⁹ S. Bitiukov,⁸⁹ K. Datsko,⁸⁹ V. Grishin,^{89,b} V. Kachanov,⁸⁹ D. Konstantinov,⁸⁹
V. Krychkin,⁸⁹ V. Petrov,⁸⁹ R. Ryutin,⁸⁹ S. Slabospitsky,⁸⁹ A. Sobol,⁸⁹ A. Sytine,⁸⁹ L. Tourtchanovitch,⁸⁹
S. Troshin,⁸⁹ N. Tyurin,⁸⁹ A. Uzunian,⁸⁹ A. Volkov,⁸⁹ P. Adzic,⁹⁰ M. Djordjevic,⁹⁰ D. Krpic,⁹⁰ D. Maletic,⁹⁰
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G. Gomez,⁹⁴ J. Gonzalez Sanchez,⁹⁴ R. Gonzalez Suarez,⁹⁴ C. Jorda,⁹⁴ P. Lobelle Pardo,⁹⁴ A. Lopez Virto,⁹⁴
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A. Ruiz Jimeno,⁹⁴ L. Scodellaro,⁹⁴ M. Sobron Sanudo,⁹⁴ I. Vila,⁹⁴ R. Vilar Cortabitarte,⁹⁴ D. Abbaneo,⁹⁵
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Kaya,¹⁰³ M. Özbek,¹⁰³ S. Ozkorucuklu,¹⁰³ N. Sonmez,¹⁰³ L. Levchuk,¹⁰⁴ P. Bell,¹⁰⁵ F. Bostock,¹⁰⁵ J. J. Brooke,¹⁰⁵ T. L. Cheng,¹⁰⁵ D. Cussans,¹⁰⁵ R. Frazier,¹⁰⁵ J. Goldstein,¹⁰⁵ M. Hansen,¹⁰⁵ G. P. Heath,¹⁰⁵ H. F. Heath,¹⁰⁵ C. Hill,¹⁰⁵ B. Huckvale,¹⁰⁵ J. Jackson,¹⁰⁵ L. Kreczko,¹⁰⁵ C. K. Mackay,¹⁰⁵ S. Metson,¹⁰⁵ D. M. Newbold,^{105,p} K. Nirunpong,¹⁰⁵ V. J. Smith,¹⁰⁵ S. Ward,¹⁰⁵ L. Basso,¹⁰⁶ K. W. Bell,¹⁰⁶ A. Belyaev,¹⁰⁶ C. Brew,¹⁰⁶ R. M. Brown,¹⁰⁶ B. Camanzi,¹⁰⁶ D. J. A. Cockerill,¹⁰⁶ J. A. Coughlan,¹⁰⁶ K. Harder,¹⁰⁶ S. Harper,¹⁰⁶ B. W. Kennedy,¹⁰⁶ E. Olaiya,¹⁰⁶ B. C. Radburn-Smith,¹⁰⁶ C. H. Shepherd-Themistocleous,¹⁰⁶ I. R. Tomalin,¹⁰⁶ W. J. Womersley,¹⁰⁶ S. D. Worm,¹⁰⁶ R. Bainbridge,¹⁰⁷ G. Ball,¹⁰⁷ J. Ballin,¹⁰⁷ R. Beuselinck,¹⁰⁷ O. Buchmuller,¹⁰⁷ D. Colling,¹⁰⁷ N. Cripps,¹⁰⁷ M. Cutajar,¹⁰⁷ G. Davies,¹⁰⁷ M. Della Negra,¹⁰⁷ C. Foudas,¹⁰⁷ J. Fulcher,¹⁰⁷ D. Futyan,¹⁰⁷ A. Guneratne Bryer,¹⁰⁷ G. Hall,¹⁰⁷ Z. Hatherell,¹⁰⁷ J. Hays,¹⁰⁷ G. Iles,¹⁰⁷ G. Karapostoli,¹⁰⁷ L. Lyons,¹⁰⁷ A.-M. Magnan,¹⁰⁷ J. Marrouche,¹⁰⁷ R. Nandi,¹⁰⁷ J. Nash,¹⁰⁷ A. Nikitenko,^{107,m} A. Papageorgiou,¹⁰⁷ M. Pesaresi,¹⁰⁷ K. Petridis,¹⁰⁷ M. Pioppi,^{107,q} D. M. Raymond,¹⁰⁷ N. Rompotis,¹⁰⁷ A. Rose,¹⁰⁷ M. J. Ryan,¹⁰⁷ C. Seez,¹⁰⁷ P. Sharp,¹⁰⁷ A. Sparrow,¹⁰⁷ M. Stoye,¹⁰⁷ A. Tapper,¹⁰⁷ S. Tourneur,¹⁰⁷ M. Vazquez Acosta,¹⁰⁷ T. Virdee,^{107,b} S. Wakefield,¹⁰⁷ D. Wardrope,¹⁰⁷ T. Whyntie,¹⁰⁷ M. Barrett,¹⁰⁸ M. Chadwick,¹⁰⁸ J. E. Cole,¹⁰⁸ P. R. Hobson,¹⁰⁸ A. Khan,¹⁰⁸ P. Kyberd,¹⁰⁸ D. Leslie,¹⁰⁸ I. D. Reid,¹⁰⁸ L. Teodorescu,¹⁰⁸ T. Bose,¹⁰⁹ A. Clough,¹⁰⁹ A. Heister,¹⁰⁹ J. St. John,¹⁰⁹ P. Lawson,¹⁰⁹ D. Lazic,¹⁰⁹ J. Rohlf,¹⁰⁹ L. Sulak,¹⁰⁹ J. Andrea,¹¹⁰ A. Avetisyan,¹¹⁰ S. Bhattacharya,¹¹⁰ J. P. Chou,¹¹⁰ D. Cutts,¹¹⁰ S. Esen,¹¹⁰ U. Heintz,¹¹⁰ S. Jabeen,¹¹⁰ G. Kukartsev,¹¹⁰ G. Landsberg,¹¹⁰ M. Narain,¹¹⁰ D. Nguyen,¹¹⁰ T. Speer,¹¹⁰ K. V. Tsang,¹¹⁰ M. A. Borgia,¹¹¹ R. Breedon,¹¹¹ M. Calderon De La Barca Sanchez,¹¹¹ D. Cebra,¹¹¹ M. Chertok,¹¹¹ J. Conway,¹¹¹ P. T. Cox,¹¹¹ J. Dolen,¹¹¹ R. Erbacher,¹¹¹ E. Friis,¹¹¹ W. Ko,¹¹¹ A. Kopecky,¹¹¹ R. Lander,¹¹¹ H. Liu,¹¹¹ S. Maruyama,¹¹¹ T. Miceli,¹¹¹ M. Nikolic,¹¹¹ D. Pellett,¹¹¹ J. Robles,¹¹¹ T. Schwarz,¹¹¹ M. Searle,¹¹¹ J. Smith,¹¹¹ M. Squires,¹¹¹ M. Tripathi,¹¹¹ R. Vasquez Sierra,¹¹¹ C. Veelken,¹¹¹ V. Andreev,¹¹² K. Arisaka,¹¹² D. Cline,¹¹² R. Cousins,¹¹² A. Deisher,¹¹² S. Erhan,^{112,b} C. Farrell,¹¹² M. Felcini,¹¹² J. Hauser,¹¹² M. Ignatenko,¹¹² C. Jarvis,¹¹² C. Plager,¹¹² G. Rakness,¹¹² P. Schlein,^{112,a} J. Tucker,¹¹² V. Valuev,¹¹² R. Wallny,¹¹² J. Babb,¹¹³ R. Clare,¹¹³ J. Ellison,¹¹³ J. W. Gary,¹¹³ G. Hanson,¹¹³ G. Y. Jeng,¹¹³ S. C. Kao,¹¹³ F. Liu,¹¹³ H. Liu,¹¹³ A. Luthra,¹¹³ H. Nguyen,¹¹³ G. Pasztor,^{113,r} A. Satpathy,¹¹³ B. C. Shen,^{113,a} R. Stringer,¹¹³ J. Sturdy,¹¹³ S. Sumowidagdo,¹¹³ R. Wilken,¹¹³ S. Wimpenny,¹¹³ W. Andrews,¹¹⁴ J. G. Branson,¹¹⁴ E. Dusinger,¹¹⁴ D. Evans,¹¹⁴ F. Golf,¹¹⁴ A. Holzner,¹¹⁴ R. Kelley,¹¹⁴ M. Lebourgeois,¹¹⁴ J. Letts,¹¹⁴ B. Mangano,¹¹⁴ J. Muelmenstaedt,¹¹⁴ S. Padhi,¹¹⁴ C. Palmer,¹¹⁴ G. Petrucciani,¹¹⁴ H. Pi,¹¹⁴ M. Pieri,¹¹⁴ R. Ranieri,¹¹⁴ M. Sani,¹¹⁴ V. Sharma,^{114,b} S. Simon,¹¹⁴ Y. Tu,¹¹⁴ A. Vartak,¹¹⁴ F. Würthwein,¹¹⁴ A. Yagil,¹¹⁴ D. Barge,¹¹⁵ M. Blume,¹¹⁵ C. Campagnari,¹¹⁵ M. D'Alfonso,¹¹⁵ T. Danielson,¹¹⁵ J. Garberson,¹¹⁵ J. Incandela,¹¹⁵ C. Justus,¹¹⁵ P. Kalavase,¹¹⁵ S. A. Koay,¹¹⁵ D. Kovalskyi,¹¹⁵ V. Krutelyov,¹¹⁵ J. Lamb,¹¹⁵ S. Lowette,¹¹⁵ V. Pavlunin,¹¹⁵ F. Rebassoo,¹¹⁵ J. Ribnik,¹¹⁵ J. Richman,¹¹⁵ R. Rossin,¹¹⁵ D. Stuart,¹¹⁵ W. To,¹¹⁵ J. R. Vlimant,¹¹⁵ M. Witherell,¹¹⁵ A. Bornheim,¹¹⁶ J. Bunn,¹¹⁶ M. Gataullin,¹¹⁶ D. Kcira,¹¹⁶ V. Litvine,¹¹⁶ Y. Ma,¹¹⁶ H. B. Newman,¹¹⁶ C. Rogan,¹¹⁶ K. Shin,¹¹⁶ V. Timciuc,¹¹⁶ J. Veverka,¹¹⁶ R. Wilkinson,¹¹⁶ Y. Yang,¹¹⁶ R. Y. Zhu,¹¹⁶ B. Akgun,¹¹⁷ R. Carroll,¹¹⁷ T. Ferguson,¹¹⁷ D. W. Jang,¹¹⁷ S. Y. Jun,¹¹⁷ M. Paulini,¹¹⁷

J. Russ,¹¹⁷ N. Terentyev,¹¹⁷ H. Vogel,¹¹⁷ I. Vorobiev,¹¹⁷ J. P. Cumalat,¹¹⁸ M. E. Dinardo,¹¹⁸ B. R. Dressel,¹¹⁸ W. T. Ford,¹¹⁸ B. Heyburn,¹¹⁸ E. Luiggi Lopez,¹¹⁸ U. Nauenberg,¹¹⁸ J. G. Smith,¹¹⁸ K. Stenson,¹¹⁸ K. A. Ulmer,¹¹⁸ S. R. Wagner,¹¹⁸ S. L. Zang,¹¹⁸ L. Agostino,¹¹⁹ J. Alexander,¹¹⁹ F. Blekman,¹¹⁹ A. Chatterjee,¹¹⁹ S. Das,¹¹⁹ N. Eggert,¹¹⁹ L. J. Fields,¹¹⁹ L. K. Gibbons,¹¹⁹ B. Heltsley,¹¹⁹ W. Hopkins,¹¹⁹ A. Khukhunaishvili,¹¹⁹ B. Kreis,¹¹⁹ V. Kuznetsov,¹¹⁹ G. Nicolas Kaufman,¹¹⁹ J. R. Patterson,¹¹⁹ D. Puigh,¹¹⁹ D. Riley,¹¹⁹ A. Ryd,¹¹⁹ X. Shi,¹¹⁹ W. Sun,¹¹⁹ W. D. Teo,¹¹⁹ J. Thom,¹¹⁹ J. Thompson,¹¹⁹ J. Vaughan,¹¹⁹ Y. Weng,¹¹⁹ P. Wittich,¹¹⁹ A. Biselli,¹²⁰ G. Cirino,¹²⁰ D. Winn,¹²⁰ S. Abdullin,¹²¹ M. Albrow,¹²¹ J. Anderson,¹²¹ G. Apollinari,¹²¹ M. Atac,¹²¹ J. A. Bakken,¹²¹ S. Banerjee,¹²¹ L. A. T. Bauerdick,¹²¹ A. Beretvas,¹²¹ J. Berryhill,¹²¹ P. C. Bhat,¹²¹ I. Bloch,¹²¹ F. Borchering,¹²¹ K. Burkett,¹²¹ J. N. Butler,¹²¹ V. Chetluru,¹²¹ H. W. K. Cheung,¹²¹ F. Chlebana,¹²¹ S. Cihangir,¹²¹ M. Demarteau,¹²¹ D. P. Eartly,¹²¹ V. D. Elvira,¹²¹ I. Fisk,¹²¹ J. Freeman,¹²¹ Y. Gao,¹²¹ E. Gottschalk,¹²¹ D. Green,¹²¹ O. Gutsche,¹²¹ A. Hahn,¹²¹ J. Hanlon,¹²¹ R. M. Harris,¹²¹ E. James,¹²¹ H. Jensen,¹²¹ M. Johnson,¹²¹ U. Joshi,¹²¹ R. Khatiwada,¹²¹ B. Kilminster,¹²¹ B. Klima,¹²¹ K. Kousouris,¹²¹ S. Kunori,¹²¹ S. Kwan,¹²¹ P. Limon,¹²¹ R. Lipton,¹²¹ J. Lykken,¹²¹ K. Maeshima,¹²¹ J. M. Marraffino,¹²¹ D. Mason,¹²¹ P. McBride,¹²¹ T. McCauley,¹²¹ T. Miao,¹²¹ K. Mishra,¹²¹ S. Mrenna,¹²¹ Y. Musienko,^{121,s} C. Newman-Holmes,¹²¹ V. O'Dell,¹²¹ S. Popescu,¹²¹ R. Pordes,¹²¹ O. Prokofyev,¹²¹ N. Saoulidou,¹²¹ E. Sexton-Kennedy,¹²¹ S. Sharma,¹²¹ R. P. Smith,^{121,a} A. Soha,¹²¹ W. J. Spalding,¹²¹ L. Spiegel,¹²¹ P. Tan,¹²¹ L. Taylor,¹²¹ S. Tkaczyk,¹²¹ L. Uplegger,¹²¹ E. W. Vaandering,¹²¹ R. Vidal,¹²¹ J. Whitmore,¹²¹ W. Wu,¹²¹ F. Yumiceva,¹²¹ J. C. Yun,¹²¹ D. Acosta,¹²² P. Avery,¹²² D. Bourilkov,¹²² M. Chen,¹²² G. P. Di Giovanni,¹²² D. Dobur,¹²² A. Drozdetskiy,¹²² R. D. Field,¹²² Y. Fu,¹²² I. K. Furic,¹²² J. Gartner,¹²² B. Kim,¹²² S. Klimentenko,¹²² J. Konigsberg,¹²² A. Korytov,¹²² K. Kotov,¹²² A. Kropivnitskaya,¹²² T. Kypreos,¹²² K. Matchev,¹²² G. Mitselmakher,¹²² Y. Pakhotin,¹²² J. Piedra Gomez,¹²² C. Prescott,¹²² R. Remington,¹²² M. Schmitt,¹²² B. Scurlock,¹²² P. Sellers,¹²² D. Wang,¹²² J. Yelton,¹²² M. Zakaria,¹²² C. Ceron,¹²³ V. Gaultney,¹²³ L. Kramer,¹²³ L. M. Lebolo,¹²³ S. Linn,¹²³ P. Markowitz,¹²³ G. Martinez,¹²³ D. Mesa,¹²³ J. L. Rodriguez,¹²³ T. Adams,¹²⁴ A. Askew,¹²⁴ J. Chen,¹²⁴ B. Diamond,¹²⁴ S. V. Gleyzer,¹²⁴ J. Haas,¹²⁴ S. Hagopian,¹²⁴ V. Hagopian,¹²⁴ M. Jenkins,¹²⁴ K. F. Johnson,¹²⁴ H. Prosper,¹²⁴ S. Sekmen,¹²⁴ V. Veeraraghavan,¹²⁴ M. M. Baarmand,¹²⁵ S. Guragain,¹²⁵ M. Hohlmann,¹²⁵ H. Kalakhety,¹²⁵ H. Mermerkaya,¹²⁵ R. Ralich,¹²⁵ I. Vodopiyarov,¹²⁵ M. R. Adams,¹²⁶ I. M. Anghel,¹²⁶ L. Apanasevich,¹²⁶ V. E. Bazterra,¹²⁶ R. R. Betts,¹²⁶ J. Callner,¹²⁶ R. Cavanaugh,¹²⁶ C. Dragoiu,¹²⁶ E. J. Garcia-Solis,¹²⁶ C. E. Gerber,¹²⁶ D. J. Hofman,¹²⁶ S. Khalatian,¹²⁶ F. Lacroix,¹²⁶ E. Shabalina,¹²⁶ A. Smoron,¹²⁶ D. Strom,¹²⁶ N. Varelas,¹²⁶ U. Akgun,¹²⁷ E. A. Albayrak,¹²⁷ B. Bilki,¹²⁷ K. Cankocak,¹²⁷ W. Clarida,¹²⁷ F. Duru,¹²⁷ C. K. Lae,¹²⁷ E. McCliment,¹²⁷ J.-P. Merlo,¹²⁷ A. Mestvirishvili,¹²⁷ A. Moeller,¹²⁷ J. Nachtman,¹²⁷ C. R. Newsom,¹²⁷ E. Norbeck,¹²⁷ J. Olson,¹²⁷ Y. Onel,¹²⁷ F. Ozok,¹²⁷ S. Sen,¹²⁷ J. Wetzel,¹²⁷ T. Yetkin,¹²⁷ K. Yi,¹²⁷ B. A. Barnett,¹²⁸ B. Blumenfeld,¹²⁸ A. Bonato,¹²⁸ C. Eskew,¹²⁸ D. Fehling,¹²⁸ G. Giurgiu,¹²⁸ A. V. Gritsan,¹²⁸ Z. J. Guo,¹²⁸ G. Hu,¹²⁸ P. Maksimovic,¹²⁸ S. Rappoccio,¹²⁸ M. Swartz,¹²⁸ N. V. Tran,¹²⁸ A. Whitbeck,¹²⁸ P. Baringer,¹²⁹ A. Bean,¹²⁹ G. Benelli,¹²⁹ O. Grachov,¹²⁹ M. Murray,¹²⁹ V. Radicci,¹²⁹ S. Sanders,¹²⁹ J. S. Wood,¹²⁹ V. Zhukova,¹²⁹ D. Bandurin,¹³⁰ T. Bolton,¹³⁰ I. Chakaberia,¹³⁰ A. Ivanov,¹³⁰ K. Kaadze,¹³⁰ Y. Maravin,¹³⁰ S. Shrestha,¹³⁰ I. Svintradze,¹³⁰ Z. Wan,¹³⁰ J. Gronberg,¹³¹ D. Lange,¹³¹ D. Wright,¹³¹ D. Baden,¹³² M. Boutemur,¹³² S. C. Eno,¹³² D. Ferencek,¹³² N. J. Hadley,¹³² R. G. Kellogg,¹³² M. Kim,¹³² A. Mignerey,¹³² K. Rossato,¹³² P. Rumerio,¹³² F. Santanastasio,¹³² A. Skuja,¹³² J. Temple,¹³² M. B. Tonjes,¹³² S. C. Tonwar,¹³² E. Twedt,¹³² B. Alver,¹³³ G. Bauer,¹³³ J. Bendavid,¹³³ W. Busza,¹³³ E. Butz,¹³³ I. A. Cali,¹³³ M. Chan,¹³³ D. D'Enterria,¹³³ P. Everaerts,¹³³ G. Gomez Ceballos,¹³³ M. Goncharov,¹³³ K. A. Hahn,¹³³ P. Harris,¹³³ Y. Kim,¹³³ M. Klute,¹³³ Y.-J. Lee,¹³³ W. Li,¹³³ C. Loizides,¹³³ P. D. Luckey,¹³³ T. Ma,¹³³ S. Nahn,¹³³ C. Paus,¹³³ C. Roland,¹³³ G. Roland,¹³³ M. Rudolph,¹³³ G. S. F. Stephans,¹³³ K. Sumorok,¹³³ K. Sung,¹³³ E. A. Wenger,¹³³ B. Wyslouck,¹³³ S. Xie,¹³³ Y. Yilmaz,¹³³ A. S. Yoon,¹³³ M. Zanetti,¹³³ P. Cole,¹³⁴ S. I. Cooper,¹³⁴ P. Cushman,¹³⁴ B. Dahmes,¹³⁴ A. De Benedetti,¹³⁴ P. R. Duderov,¹³⁴ G. Franzoni,¹³⁴ J. Haupt,¹³⁴ K. Klapoetke,¹³⁴ Y. Kubota,¹³⁴ J. Mans,¹³⁴ D. Petyt,¹³⁴ V. Rekovic,¹³⁴ R. Rusack,¹³⁴ M. Sasseville,¹³⁴ A. Singovsky,¹³⁴ L. M. Cremaldi,¹³⁵ R. Godang,¹³⁵ R. Kroeger,¹³⁵ L. Perera,¹³⁵ R. Rahmat,¹³⁵ D. A. Sanders,¹³⁵ P. Sonnek,¹³⁵ D. Summers,¹³⁵ K. Bloom,¹³⁶ S. Bose,¹³⁶ J. Butt,¹³⁶ D. R. Claes,¹³⁶ A. Dominguez,¹³⁶ M. Eads,¹³⁶ J. Keller,¹³⁶ T. Kelly,¹³⁶ I. Kravchenko,¹³⁶ J. Lazo-Flores,¹³⁶ C. Lundstedt,¹³⁶ H. Malbouisson,¹³⁶ S. Malik,¹³⁶ G. R. Snow,¹³⁶ U. Baur,¹³⁷ I. Iashvili,¹³⁷ A. Kharchilava,¹³⁷ A. Kumar,¹³⁷ K. Smith,¹³⁷ M. Strang,¹³⁷ J. Zennaro,¹³⁷ G. Alverson,¹³⁸ E. Barberis,¹³⁸ D. Baumgartel,¹³⁸ O. Boeriu,¹³⁸ S. Reucroft,¹³⁸ J. Swain,¹³⁸ D. Wood,¹³⁸ J. Zhang,¹³⁸ A. Anastassov,¹³⁹ A. Kubik,¹³⁹ R. A. Ofierzynski,¹³⁹ A. Pozdnyakov,¹³⁹ M. Schmitt,¹³⁹ S. Stoynev,¹³⁹ M. Velasco,¹³⁹ S. Won,¹³⁹ L. Antonelli,¹⁴⁰ D. Berry,¹⁴⁰ M. Hildreth,¹⁴⁰ C. Jessop,¹⁴⁰ D. J. Karmgard,¹⁴⁰ J. Kolb,¹⁴⁰ T. Kolberg,¹⁴⁰

K. Lannon,¹⁴⁰ S. Lynch,¹⁴⁰ N. Marinelli,¹⁴⁰ D. M. Morse,¹⁴⁰ R. Ruchti,¹⁴⁰ J. Slaunwhite,¹⁴⁰ N. Valls,¹⁴⁰ J. Warchol,¹⁴⁰ M. Wayne,¹⁴⁰ J. Ziegler,¹⁴⁰ B. Bylsma,¹⁴¹ L. S. Durkin,¹⁴¹ J. Gu,¹⁴¹ P. Killewald,¹⁴¹ T. Y. Ling,¹⁴¹ G. Williams,¹⁴¹ N. Adam,¹⁴² E. Berry,¹⁴² P. Elmer,¹⁴² D. Gerbaudo,¹⁴² V. Halyo,¹⁴² A. Hunt,¹⁴² J. Jones,¹⁴² E. Laird,¹⁴² D. Lopes Pegna,¹⁴² D. Marlow,¹⁴² T. Medvedeva,¹⁴² M. Mooney,¹⁴² J. Olsen,¹⁴² P. Piroué,¹⁴² D. Stickland,¹⁴² C. Tully,¹⁴² J. S. Werner,¹⁴² A. Zuranski,¹⁴² J. G. Acosta,¹⁴³ X. T. Huang,¹⁴³ A. Lopez,¹⁴³ H. Mendez,¹⁴³ S. Oliveros,¹⁴³ J. E. Ramirez Vargas,¹⁴³ A. Zatzerklyaniy,¹⁴³ E. Alagoz,¹⁴⁴ V. E. Barnes,¹⁴⁴ G. Bolla,¹⁴⁴ L. Borrello,¹⁴⁴ D. Bortoletto,¹⁴⁴ A. Everett,¹⁴⁴ A. F. Garfinkel,¹⁴⁴ Z. Gece,¹⁴⁴ L. Gutay,¹⁴⁴ M. Jones,¹⁴⁴ O. Koybasi,¹⁴⁴ A. T. Laasanen,¹⁴⁴ N. Leonardo,¹⁴⁴ C. Liu,¹⁴⁴ V. Maroussov,¹⁴⁴ P. Merkel,¹⁴⁴ D. H. Miller,¹⁴⁴ N. Neumeister,¹⁴⁴ K. Potamianos,¹⁴⁴ I. Shipsey,¹⁴⁴ D. Silvers,¹⁴⁴ H. D. Yoo,¹⁴⁴ J. Zablocki,¹⁴⁴ Y. Zheng,¹⁴⁴ P. Jindal,¹⁴⁵ N. Parashar,¹⁴⁵ V. Cuplov,¹⁴⁶ K. M. Ecklund,¹⁴⁶ F. J. M. Geurts,¹⁴⁶ J. H. Liu,¹⁴⁶ J. Morales,¹⁴⁶ B. P. Padley,¹⁴⁶ R. Redjimi,¹⁴⁶ J. Roberts,¹⁴⁶ B. Betchart,¹⁴⁷ A. Bodek,¹⁴⁷ Y. S. Chung,¹⁴⁷ P. de Barbaro,¹⁴⁷ R. Demina,¹⁴⁷ H. Flacher,¹⁴⁷ A. Garcia-Bellido,¹⁴⁷ Y. Gotra,¹⁴⁷ J. Han,¹⁴⁷ A. Harel,¹⁴⁷ D. C. Miner,¹⁴⁷ D. Orbaker,¹⁴⁷ G. Petrillo,¹⁴⁷ D. Vishnevskiy,¹⁴⁷ M. Zielinski,¹⁴⁷ A. Bhatti,¹⁴⁸ L. Demortier,¹⁴⁸ K. Goulianos,¹⁴⁸ K. Hatakeyama,¹⁴⁸ G. Lungu,¹⁴⁸ C. Mesropian,¹⁴⁸ M. Yan,¹⁴⁸ O. Atramentov,¹⁴⁹ Y. Gershtein,¹⁴⁹ R. Gray,¹⁴⁹ E. Halkiadakis,¹⁴⁹ D. Hidas,¹⁴⁹ D. Hits,¹⁴⁹ A. Lath,¹⁴⁹ K. Rose,¹⁴⁹ S. Schnetzer,¹⁴⁹ S. Somalwar,¹⁴⁹ R. Stone,¹⁴⁹ S. Thomas,¹⁴⁹ G. Cerizza,¹⁵⁰ M. Hollingsworth,¹⁵⁰ S. Spanier,¹⁵⁰ Z. C. Yang,¹⁵⁰ A. York,¹⁵⁰ J. Asaadi,¹⁵¹ R. Eusebi,¹⁵¹ J. Gilmore,¹⁵¹ A. Gurrola,¹⁵¹ T. Kamon,¹⁵¹ V. Khotilovich,¹⁵¹ R. Montalvo,¹⁵¹ C. N. Nguyen,¹⁵¹ J. Pivarski,¹⁵¹ A. Safonov,¹⁵¹ S. Sengupta,¹⁵¹ D. Toback,¹⁵¹ M. Weinberger,¹⁵¹ N. Akchurin,¹⁵² C. Bardak,¹⁵² J. Damgov,¹⁵² C. Jeong,¹⁵² K. Kovitanggoon,¹⁵² S. W. Lee,¹⁵² P. Mane,¹⁵² Y. Roh,¹⁵² A. Sill,¹⁵² I. Volobouev,¹⁵² R. Wigmans,¹⁵² E. Yazgan,¹⁵² E. Appel,¹⁵³ E. Brownson,¹⁵³ D. Engh,¹⁵³ C. Florez,¹⁵³ W. Gabella,¹⁵³ W. Johns,¹⁵³ P. Kurt,¹⁵³ C. Maguire,¹⁵³ A. Melo,¹⁵³ P. Sheldon,¹⁵³ J. Velkovska,¹⁵³ M. W. Arenton,¹⁵⁴ M. Balazs,¹⁵⁴ M. Buehler,¹⁵⁴ S. Conetti,¹⁵⁴ B. Cox,¹⁵⁴ R. Hirosky,¹⁵⁴ A. Ledovskoy,¹⁵⁴ C. Neu,¹⁵⁴ R. Yohay,¹⁵⁴ S. Gollapinni,¹⁵⁵ K. Gunthoti,¹⁵⁵ R. Harr,¹⁵⁵ P. E. Karchin,¹⁵⁵ M. Mattson,¹⁵⁵ C. Milstène,¹⁵⁵ A. Sakharov,¹⁵⁵ M. Anderson,¹⁵⁶ M. Bachtis,¹⁵⁶ J. N. Bellinger,¹⁵⁶ D. Carlsmith,¹⁵⁶ S. Dasu,¹⁵⁶ S. Dutta,¹⁵⁶ J. Efron,¹⁵⁶ L. Gray,¹⁵⁶ K. S. Grogg,¹⁵⁶ M. Grothe,¹⁵⁶ R. Hall-Wilton,^{156,b} M. Herndon,¹⁵⁶ P. Klabbers,¹⁵⁶ J. Klukas,¹⁵⁶ A. Lanaro,¹⁵⁶ C. Lazaridis,¹⁵⁶ J. Leonard,¹⁵⁶ D. Lomidze,¹⁵⁶ R. Loveless,¹⁵⁶ A. Mohapatra,¹⁵⁶ G. Polese,¹⁵⁶ D. Reeder,¹⁵⁶ A. Savin,¹⁵⁶ W. H. Smith,¹⁵⁶ J. Swanson,¹⁵⁶ and M. Weinberg¹⁵⁶

(CMS Collaboration)

¹*Yerevan Physics Institute, Yerevan, Armenia*²*Institut für Hochenergiephysik der OeAW, Wien, Austria*³*National Centre for Particle and High Energy Physics, Minsk, Belarus*⁴*Universiteit Antwerpen, Antwerpen, Belgium*⁵*Vrije Universiteit Brussel, Brussel, Belgium*⁶*Université Libre de Bruxelles, Bruxelles, Belgium*⁷*Ghent University, Ghent, Belgium*⁸*Université Catholique de Louvain, Louvain-la-Neuve, Belgium*⁹*Université de Mons, Mons, Belgium*¹⁰*Centro Brasileiro de Pesquisas Fisicas, Rio de Janeiro, Brazil*¹¹*Universidade do Estado do Rio de Janeiro, Rio de Janeiro, Brazil*¹²*Instituto de Fisica Teorica, Universidade Estadual Paulista, Sao Paulo, Brazil*¹³*Institute for Nuclear Research and Nuclear Energy, Sofia, Bulgaria*¹⁴*University of Sofia, Sofia, Bulgaria*¹⁵*Institute of High Energy Physics, Beijing, China*¹⁶*State Key Laboratory of Nuclear Physics and Technology, Peking University, Beijing, China*¹⁷*Universidad de Los Andes, Bogota, Colombia*¹⁸*Technical University of Split, Split, Croatia*¹⁹*University of Split, Split, Croatia*²⁰*Institute Rudjer Boskovic, Zagreb, Croatia*²¹*University of Cyprus, Nicosia, Cyprus*²²*Academy of Scientific Research and Technology of the Arab Republic of Egypt, Egyptian Network of High Energy Physics, Cairo, Egypt*²³*National Institute of Chemical Physics and Biophysics, Tallinn, Estonia*²⁴*Department of Physics, University of Helsinki, Helsinki, Finland*

- ²⁵*Helsinki Institute of Physics, Helsinki, Finland*
- ²⁶*Lappeenranta University of Technology, Lappeenranta, Finland*
- ²⁷*Laboratoire d'Annecy-le-Vieux de Physique des Particules, IN2P3-CNRS, Annecy-le-Vieux, France*
- ²⁸*DSM/IRFU, CEA/Saclay, Gif-sur-Yvette, France*
- ²⁹*Laboratoire Leprince-Ringuet, Ecole Polytechnique, IN2P3-CNRS, Palaiseau, France*
- ³⁰*Institut Pluridisciplinaire Hubert Curien, Université de Strasbourg, Université de Haute Alsace Mulhouse, CNRS/IN2P3, Strasbourg, France*
- ³¹*Centre de Calcul de l'Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France*
- ³²*Université de Lyon, Université Claude Bernard Lyon 1, CNRS-IN2P3, Institut de Physique Nucléaire de Lyon, Villeurbanne, France*
- ³³*E. Andronikashvili Institute of Physics, Academy of Science, Tbilisi, Georgia*
- ³⁴*RWTH Aachen University, I. Physikalisches Institut, Aachen, Germany*
- ³⁵*RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany*
- ³⁶*RWTH Aachen University, III. Physikalisches Institut B, Aachen, Germany*
- ³⁷*Deutsches Elektronen-Synchrotron, Hamburg, Germany*
- ³⁸*University of Hamburg, Hamburg, Germany*
- ³⁹*Institut für Experimentelle Kernphysik, Karlsruhe, Germany*
- ⁴⁰*Institute of Nuclear Physics "Demokritos," Aghia Paraskevi, Greece*
- ⁴¹*University of Athens, Athens, Greece*
- ⁴²*University of Ioánnina, Ioánnina, Greece*
- ⁴³*KFKI Research Institute for Particle and Nuclear Physics, Budapest, Hungary*
- ⁴⁴*Institute of Nuclear Research ATOMKI, Debrecen, Hungary*
- ⁴⁵*University of Debrecen, Debrecen, Hungary*
- ⁴⁶*Panjab University, Chandigarh, India*
- ⁴⁷*University of Delhi, Delhi, India*
- ⁴⁸*Bhabha Atomic Research Centre, Mumbai, India*
- ⁴⁹*Tata Institute of Fundamental Research–EHEP, Mumbai, India*
- ⁵⁰*Tata Institute of Fundamental Research–HECR, Mumbai, India*
- ⁵¹*Institute for Studies in Theoretical Physics & Mathematics (IPM), Tehran, Iran*
- ^{52a}*INFN Sezione di Bari, Bari, Italy*
- ^{52b}*Università di Bari, Bari, Italy*
- ^{52c}*Politecnico di Bari, Bari, Italy*
- ^{53a}*INFN Sezione di Bologna, Bologna, Italy*
- ^{53b}*Università di Bologna, Bologna, Italy*
- ^{54a}*INFN Sezione di Catania, Catania, Italy*
- ^{54b}*Università di Catania, Catania, Italy*
- ^{55a}*INFN Sezione di Firenze, Firenze, Italy*
- ^{55b}*Università di Firenze, Firenze, Italy*
- ⁵⁶*INFN Laboratori Nazionali di Frascati, Frascati, Italy*
- ⁵⁷*INFN Sezione di Genova, Genova, Italy*
- ^{58a}*INFN Sezione di Milano-Bicocca, Milano, Italy*
- ^{58b}*Università di Milano-Bicocca, Milano, Italy*
- ^{59a}*INFN Sezione di Napoli, Napoli, Italy*
- ^{59b}*Università di Napoli "Federico II," Napoli, Italy*
- ^{60a}*INFN Sezione di Padova, Padova, Italy*
- ^{60b}*Università di Padova, Padova, Italy*
- ^{60c}*Università di Trento (Trento), Padova, Italy*
- ^{61a}*INFN Sezione di Pavia, Pavia, Italy*
- ^{61b}*Università di Pavia, Pavia, Italy*
- ^{62a}*INFN Sezione di Perugia, Perugia, Italy*
- ^{62b}*Università di Perugia, Perugia, Italy*
- ^{63a}*INFN Sezione di Pisa, Pisa, Italy*
- ^{63b}*Università di Pisa, Pisa, Italy*
- ^{63c}*Scuola Normale Superiore di Pisa, Pisa, Italy*
- ^{64a}*INFN Sezione di Roma, Roma, Italy*
- ^{64b}*Università di Roma "La Sapienza," Roma, Italy*
- ^{65a}*INFN Sezione di Torino, Torino, Italy*
- ^{65b}*Università di Torino, Torino, Italy*
- ^{65c}*Università del Piemonte Orientale (Novara), Torino, Italy*
- ^{66a}*INFN Sezione di Trieste, Trieste, Italy*
- ^{66b}*Università di Trieste, Trieste, Italy*
- ⁶⁷*Kyungpook National University, Daegu, Korea*

- ⁶⁸*Chonnam National University, Institute for Universe and Elementary Particles, Kwangju, Korea*
⁶⁹*Korea University, Seoul, Korea*
⁷⁰*University of Seoul, Seoul, Korea*
⁷¹*Sungkyunkwan University, Suwon, Korea*
⁷²*Vilnius University, Vilnius, Lithuania*
⁷³*Centro de Investigacion y de Estudios Avanzados del IPN, Mexico City, Mexico*
⁷⁴*Universidad Iberoamericana, Mexico City, Mexico*
⁷⁵*Benemerita Universidad Autonoma de Puebla, Puebla, Mexico*
⁷⁶*Universidad Autónoma de San Luis Potosí, San Luis Potosí, Mexico*
⁷⁷*University of Auckland, Auckland, New Zealand*
⁷⁸*University of Canterbury, Christchurch, New Zealand*
⁷⁹*National Centre for Physics, Quaid-I-Azam University, Islamabad, Pakistan*
⁸⁰*Institute of Experimental Physics, Warsaw, Poland*
⁸¹*Soltan Institute for Nuclear Studies, Warsaw, Poland*
⁸²*Laboratório de Instrumentação e Física Experimental de Partículas, Lisboa, Portugal*
⁸³*Joint Institute for Nuclear Research, Dubna, Russia*
⁸⁴*Petersburg Nuclear Physics Institute, Gatchina (St. Petersburg), Russia*
⁸⁵*Institute for Nuclear Research, Moscow, Russia*
⁸⁶*Institute for Theoretical and Experimental Physics, Moscow, Russia*
⁸⁷*Moscow State University, Moscow, Russia*
⁸⁸*P. N. Lebedev Physical Institute, Moscow, Russia*
⁸⁹*State Research Center of Russian Federation, Institute for High Energy Physics, Protvino, Russia*
⁹⁰*Vinca Institute of Nuclear Sciences, Belgrade, Serbia*
⁹¹*Centro de Investigaciones Energéticas Medioambientales y Tecnológicas (CIEMAT), Madrid, Spain*
⁹²*Universidad Autónoma de Madrid, Madrid, Spain*
⁹³*Universidad de Oviedo, Oviedo, Spain*
⁹⁴*Instituto de Física de Cantabria (IFCA), CSIC-Universidad de Cantabria, Santander, Spain*
⁹⁵*CERN, European Organization for Nuclear Research, Geneva, Switzerland*
⁹⁶*Paul Scherrer Institut, Villigen, Switzerland*
⁹⁷*Institute for Particle Physics, ETH Zurich, Zurich, Switzerland*
⁹⁸*Universität Zürich, Zurich, Switzerland*
⁹⁹*National Central University, Chung-Li, Taiwan*
¹⁰⁰*National Taiwan University (NTU), Taipei, Taiwan*
¹⁰¹*Cukurova University, Adana, Turkey*
¹⁰²*Physics Department, Middle East Technical University, Ankara, Turkey*
¹⁰³*Department of Physics, Bogaziçi University, Istanbul, Turkey*
¹⁰⁴*National Scientific Center, Kharkov Institute of Physics and Technology, Kharkov, Ukraine*
¹⁰⁵*University of Bristol, Bristol, United Kingdom*
¹⁰⁶*Rutherford Appleton Laboratory, Didcot, United Kingdom*
¹⁰⁷*Imperial College, University of London, London, United Kingdom*
¹⁰⁸*Brunel University, Uxbridge, United Kingdom*
¹⁰⁹*Boston University, Boston, Massachusetts 02215, USA*
¹¹⁰*Brown University, Providence, Rhode Island 02912, USA*
¹¹¹*University of California, Davis, Davis, California 95616, USA*
¹¹²*University of California, Los Angeles, Los Angeles, California 90095, USA*
¹¹³*University of California, Riverside, Riverside, California 92521, USA*
¹¹⁴*University of California, San Diego, La Jolla, California 92093, USA*
¹¹⁵*University of California, Santa Barbara, Santa Barbara, California 93106, USA*
¹¹⁶*California Institute of Technology, Pasadena, California 91125, USA*
¹¹⁷*Carnegie Mellon University, Pittsburgh, Pennsylvania 15213, USA*
¹¹⁸*University of Colorado at Boulder, Boulder, Colorado 80309, USA*
¹¹⁹*Cornell University, Ithaca, New York 14853-5001, USA*
¹²⁰*Fairfield University, Fairfield, Connecticut 06824, USA*
¹²¹*Fermi National Accelerator Laboratory, Batavia, Illinois 60510-0500, USA*
¹²²*University of Florida, Gainesville, Florida 32611-8440, USA*
¹²³*Florida International University, Miami, Florida 33199, USA*
¹²⁴*Florida State University, Tallahassee, Florida 32306-4350, USA*
¹²⁵*Florida Institute of Technology, Melbourne, Florida 32901, USA*
¹²⁶*University of Illinois at Chicago (UIC), Chicago, Illinois 60607-7059, USA*
¹²⁷*The University of Iowa, Iowa City, Iowa 52242-1479, USA*
¹²⁸*Johns Hopkins University, Baltimore, Maryland 21218, USA*

- ¹²⁹*The University of Kansas, Lawrence, Kansas 66045, USA*
¹³⁰*Kansas State University, Manhattan, Kansas 66506, USA*
¹³¹*Lawrence Livermore National Laboratory, Livermore, California 94720, USA*
¹³²*University of Maryland, College Park, Maryland 20742, USA*
¹³³*Massachusetts Institute of Technology, Cambridge, Massachusetts 02139, USA*
¹³⁴*University of Minnesota, Minneapolis, Minnesota 55455, USA*
¹³⁵*University of Mississippi, University, Mississippi 38677, USA*
¹³⁶*University of Nebraska–Lincoln, Lincoln, Nebraska 68588-0111, USA*
¹³⁷*State University of New York at Buffalo, Buffalo, New York 14260-1500, USA*
¹³⁸*Northeastern University, Boston, Massachusetts 02115, USA*
¹³⁹*Northwestern University, Evanston, Illinois 60208-3112, USA*
¹⁴⁰*University of Notre Dame, Notre Dame, Indiana 46556, USA*
¹⁴¹*The Ohio State University, Columbus, Ohio 43210, USA*
¹⁴²*Princeton University, Princeton, New Jersey 08544-0708, USA*
¹⁴³*University of Puerto Rico, Mayaguez, Puerto Rico 00680*
¹⁴⁴*Purdue University, West Lafayette, Indiana 47907-1396, USA*
¹⁴⁵*Purdue University Calumet, Hammond, Indiana 46323, USA*
¹⁴⁶*Rice University, Houston, Texas 77251-1892, USA*
¹⁴⁷*University of Rochester, Rochester, New York 14627-0171, USA*
¹⁴⁸*The Rockefeller University, New York, New York 10021-6399, USA*
¹⁴⁹*Rutgers, The State University of New Jersey, Piscataway, New Jersey 08854-8019, USA*
¹⁵⁰*University of Tennessee, Knoxville, Tennessee 37996-1200, USA*
¹⁵¹*Texas A&M University, College Station, Texas 77843-4242, USA*
¹⁵²*Texas Tech University, Lubbock, Texas 79409-1051, USA*
¹⁵³*Vanderbilt University, Nashville, Tennessee 37235, USA*
¹⁵⁴*University of Virginia, Charlottesville, Virginia 22901, USA*
¹⁵⁵*Wayne State University, Detroit, Michigan 48202, USA*
¹⁵⁶*University of Wisconsin, Madison, Wisconsin 53706, USA*

^aDeceased.

^bAlso at CERN, European Organization for Nuclear Research, Geneva, Switzerland.

^cAlso at Laboratoire Leprince-Ringuet, Ecole Polytechnique, IN2P3-CNRS, Palaiseau, France.

^dAlso at Soltan Institute for Nuclear Studies, Warsaw, Poland.

^eAlso at Moscow State University, Moscow, Russia.

^fAlso at Institute of Nuclear Research ATOMKI, Debrecen, Hungary.

^gAlso at University of California, San Diego, La Jolla, CA, USA.

^hAlso at Tata Institute of Fundamental Research–HECR, Mumbai, India.

ⁱAlso at California Institute of Technology, Pasadena, CA, USA.

^jAlso at INFN Sezione di Roma, Università di Roma “La Sapienza,” Roma, Italy.

^kAlso at University of Athens, Athens, Greece.

^lAlso at The University of Kansas, Lawrence, KS, USA.

^mAlso at Institute for Theoretical and Experimental Physics, Moscow, Russia.

ⁿAlso at Paul Scherrer Institut, Villigen, Switzerland.

^oAlso at Vinca Institute of Nuclear Sciences, Belgrade, Serbia.

^pAlso at Rutherford Appleton Laboratory, Didcot, United Kingdom.

^qAlso at INFN Sezione di Perugia, Università di Perugia, Perugia, Italy.

^rAlso at KFKI Research Institute for Particle and Nuclear Physics, Budapest, Hungary.

^sAlso at Institute for Nuclear Research, Moscow, Russia.