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## Measurement of $d\sigma/dy$ for high mass Drell-Yan $e^+e^-$ pairs from $p\bar{p}$ collisions at $\sqrt{s}=1.8$ TeV

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Collider Detector at Fermilab Collaboration

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# Measurement of $d\sigma/dy$ for high mass Drell-Yan $e^+e^-$ pairs from $p\bar{p}$ collisions at $\sqrt{s}=1.8$ TeV

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We report on the first measurement of the rapidity distribution  $d\sigma/dy$  over nearly the entire kinematic region of rapidity for  $e^+e^-$  pairs in both the Z-boson region of  $66 < M_{ee} < 116$  GeV/ $c^2$  and in the high mass region of  $M_{ee} > 116$  GeV/ $c^2$ . The data sample consists of 108 pb<sup>-1</sup> of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.8$  TeV taken by the Collider Detector at Fermilab during 1992–1995. The total cross section in the Z-boson region is measured to be  $252 \pm 11$  pb. The measured total cross section and  $d\sigma/dy$  are compared with quantum chromodynamics calculations in leading and higher orders.

Most measurements at high energy proton-antiproton colliders are performed in the central rapidity production region,  $|y| < 1$ . A model dependent extrapolation for  $|y| > 1$  is needed to extract the total cross section for hard processes such as top quark production or  $W$  and  $Z$  boson production. This extrapolation is made using Monte Carlo programs (e.g. PYTHIA [1]), which incorporate quantum chromodynamics (QCD) calculations in leading order (LO) or next to leading order (NLO). A previous measurement of the rapidity distribution,  $d\sigma/dy$ , using dimuon pairs in the  $Z$ -boson mass region of 66–116 GeV/ $c^2$ , was limited to  $|y| < 1$  [2]. In this Rapid Communication, we present the first measurement of  $d\sigma/dy$  over nearly the entire kinematic region of rapidity for  $e^+e^-$  pairs in both the  $Z$ -boson region of  $66 < M_{ee} < 116$  GeV/ $c^2$  and the high mass region of  $M_{ee} > 116$  GeV/ $c^2$ . At the Fermilab Tevatron  $p\bar{p}$  collider, the kinematic limit at the  $Z$ -boson mass is  $|y| = 3.0$ , while we measure  $|y|$  up to 2.8. The  $d\sigma/dy$  distributions are compared to the predictions of QCD in LO and NLO. This measurement is also relevant for precision  $W$ -boson mass measurements at hadron colliders, where  $W$ 's are reconstructed using  $e\nu$  and  $\mu\nu$  pairs from the Drell-Yan process.

In hadron-hadron collisions at high energies, massive  $e^+e^-$  pairs are produced via the Drell-Yan [3] process. In the standard model, quark-antiquark annihilations form an intermediate  $\gamma^*$  or  $Z$  ( $\gamma^*/Z$ ) vector boson, which then decays into an  $e^+e^-$  pair. Pairs within the 66–116 GeV/ $c^2$  mass region are predominantly from the resonant production and decay of  $Z$  bosons. In LO, the momentum fraction  $x_1$  ( $x_2$ ) of the partons in the proton (antiproton) is related to the rapidity,  $y$  [4], of the dileptons:  $x_{1,2} = (M/\sqrt{s})e^{\pm y}$ , where  $M$  is the mass of the dilepton pair and  $\sqrt{s}$  is the center of mass energy. Therefore, dilepton pairs which are produced at large rapidity originate from events in which one parton is at large  $x$  and another parton is at very small  $x$ . Since the quark distributions for  $x$  up to 0.9 are well constrained by the deep-inelastic lepton scattering experiments [5], comparisons of data and theory for  $d\sigma/dy$  and the total cross sections provide a test of the theory, e.g., missing NNLO [6,7] or power correction [8] terms.

The  $e^+e^-$  pairs are from 108 pb $^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.8$  TeV taken by the Collider Detector at Fermilab [4] (CDF) during 1992–1993 ( $18.7 \pm 0.7$  pb $^{-1}$ ) and 1994–1995 ( $89.1 \pm 3.7$  pb $^{-1}$ ). CDF is a solenoidal magnetic spectrometer surrounded by projective-tower-geometry calorimeters and outer-muon detectors. Only detector components used in this measurement are described here. Charged particle momenta and directions are measured by the spectrometer, which consists of a 1.4 T axial magnetic field, an 84-layer cylindrical drift chamber (CTC), an inner vertex tracking chamber (VTX), and a silicon vertex detector (SVX). The polar coverage of the CTC tracking is  $|\eta| < 1.2$ . The  $p\bar{p}$  collision point along the beam line ( $Z_{\text{vertex}}$ ) is determined using tracks in the VTX. The energies and directions [4] of electrons, photons, and jets are measured by three separate calorimeters covering three regions: central ( $|\eta| < 1.1$ ), end

plug ( $1.1 < |\eta| < 2.4$ ), and forward ( $2.2 < |\eta| < 4.2$ ). Each region has an electromagnetic (EM) and hadronic (HAD) calorimeter.

In a previous paper [9], we presented the  $d\sigma/dP_T$  of  $e^+e^-$  pairs in the  $Z$ -boson mass region. We have extended that analysis by adding very forward end plug-end plug (PP) and end plug-forward (PF)  $e^+e^-$  pairs to its sample of central-central (CC), central-end plug (CP), and central-forward (CF)  $e^+e^-$  pairs. The inclusion of these very forward pairs increases the event sample by 20% and extends the  $Z$ -boson rapidity coverage up to  $|y|$  of 2.8. For these very forward pairs, an improved VTX track finder that significantly reduces backgrounds is used. The VTX covers the entire rapidity range and is important for the removal of background in the high  $\eta$  region which is not covered by the CTC.

The sample of  $e^+e^-$  events was collected by a three-level online trigger that required an electron in either the central or the plug calorimeter. The offline analysis selects events with two or more electron candidates. Since the electrons from the Drell-Yan process are typically isolated, both electrons are required to be isolated from any other activity in the calorimeters. Electrons in the central, end plug, and forward regions are required to be within the fiducial area of the calorimeters and have a minimum  $E_T$  of 22, 20, and 15 GeV, respectively. To improve the purity of the sample, electron identification cuts are applied [9]. For CC, CP, and CF events, the central electron (or one of them if there are two) is required to pass strict criteria. The criteria on the other electron are looser. A central electron must have a CTC track that extrapolates to the electron's shower clusters in the EM calorimeter. These clusters must have EM-like transverse shower profiles. The track momentum and the EM shower energy must be consistent with one another. The track is also used to determine the position and direction of the central electron. The fraction of energy in the HAD calorimeter towers behind the EM shower is required to be consistent with that expected for an EM shower ( $E_{\text{HAD}}/E_{\text{EM}}$ ). A plug electron must have an EM-like transverse shower profile and pass the  $E_{\text{HAD}}/E_{\text{EM}}$  requirement. A forward electron is only required to pass the  $E_{\text{HAD}}/E_{\text{EM}}$  requirement. For PP and PF events, the plug and forward electrons are additionally required to have a track in the VTX which originates from the same vertex as the other electron and points to the position of the electromagnetic cluster in the calorimeter. The ratio of found to expected hits in the VTX is required to be greater than 70% and 50% for plug and forward electron tracks, respectively. The tracking efficiency is  $(97.8 \pm 0.3)\%$  for plug electrons and  $(97.0 \pm 0.9)\%$  for forward electrons. The track finder is an improvement over the previous VTX tracker, as it reduces the background rates from 11% to 2% for PP events and from 22% to 4% for PF events.

After all cuts, the numbers of CC, CP, CF, PP, and PF events in the  $Z$ -boson mass region ( $66 < M_{ee} < 116$  GeV/ $c^2$ ) are 2894, 3811, 621, 1236, and 589, respectively. The backgrounds are low and are estimated using the data. The backgrounds in the CP, CF, PP, and PF topologies are dominated by jets and are  $28 \pm 5$ ,  $13 \pm 3$ ,  $24 \pm 2$ , and  $23 \pm 3$  events, respectively. Because of the CTC tracking re-



quirement, the jet background for the CC sample is negligible. The CC background is mainly from  $e^+e^-$  pairs from  $W^+W^-$ ,  $\tau^+\tau^-$ ,  $c\bar{c}$ ,  $b\bar{b}$ , and  $t\bar{t}$  sources. This background, estimated using  $e^\pm\mu^\mp$  pairs [10], is  $3 \pm 2$  events. The numbers of CC, CP, CF, PP, and PF events in the high mass region ( $M_{ee} > 116$  GeV/ $c^2$ ) are 61, 59, 9, 18, and 5, respectively. Within this region, the mean mass is 153 GeV/ $c^2$ .

The acceptance for Drell-Yan  $e^+e^-$  pairs is obtained using the Monte Carlo event generator, PYTHIA [1], as well as CDF detector simulation programs. PYTHIA generates the LO QCD interaction,  $q + \bar{q} \rightarrow \gamma^*/Z$ , simulates initial state QCD radiation via its parton shower algorithms, and generates the decay,  $\gamma^*/Z \rightarrow e^+e^-$ . The LO  $e^+e^-$  pair mass spectrum is brought into better agreement with the NLO spectrum by using a multiplicative  $K$ -factor for the one-loop virtual gluon correction to the LO process [11],  $K(M^2) = 1 + \frac{4}{3}(1 + \frac{4}{3}\pi^2)\alpha_s(M^2)/2\pi$ , where  $\alpha_s$  is the two-loop QCD coupling. For  $M > 50$  GeV/ $c^2$ ,  $1.25 < K < 1.36$ . The CTEQ3L [12] parton distribution functions (PDFs) are used in the acceptance calculations. Final state QED radiation [13] from the  $\gamma^*/Z \rightarrow e^+e^-$  vertex is added by the PHOTOS [14] Monte Carlo program. Generated events are processed by CDF detector simulation programs and are reconstructed as data. The calorimetry energy scales and resolutions used in the detector simulations are extracted from the data, as are the cut efficiencies and corresponding errors. Simulated events are accepted if, after the reconstruction, they pass the  $e^+e^-$  pair mass, the detector fiducial, and kinematic ( $E_T$ ) cuts.

In the  $d\sigma/dy$  measurement, the various samples are combined and binned in  $y$ . The  $d\sigma/dy$  is calculated with

$$\frac{d\sigma}{dy} = \frac{\Delta N}{C \Delta y \sum_r \mathcal{L}_r \epsilon A_r}.$$

The  $\Delta N$  is the background-subtracted event count in a  $y$  bin,  $C$  is a bin centering correction,  $\Delta y$  is the bin width, the sum  $r$  is over the 1992–1993 and 1994–1995 runs,  $\mathcal{L}_r$  is the integrated luminosity, and  $\epsilon A_r$  is the run's combined event selection efficiency and acceptance. The backgrounds subtracted from the event count are predicted using the data and background samples. The factor  $C$  corrects the average value of the cross section in the bin to its bin center value. Acceptances are calculated separately for CC, CP, CF, PP, and PF pairs. They are combined with the corresponding event selection efficiencies to give  $\epsilon A_r$ . Figure 1 shows the  $\epsilon A_r$  for events in the Z-boson mass region as a function of  $y$  for the CC sample, the CC+CP+CF sample, and the CC+CP+CF+PP+PF sample, respectively. The PP+PF events extend the acceptance in  $|y|$  from 1.2 to 2.8, and significantly increase our statistics. The  $d\sigma/dy$  of the  $e^+e^-$  events in the Z-boson mass region is shown in Table I.

The systematic errors considered are from variations in the background estimates using different methods, the background in the efficiency sample, the uncertainty in energy resolution of the calorimeter, the choice of PDFs, and the distribution of  $\gamma^*/Z$   $P_T$  used in the Monte Carlo event generator. The systematic error on the acceptance from the calo-

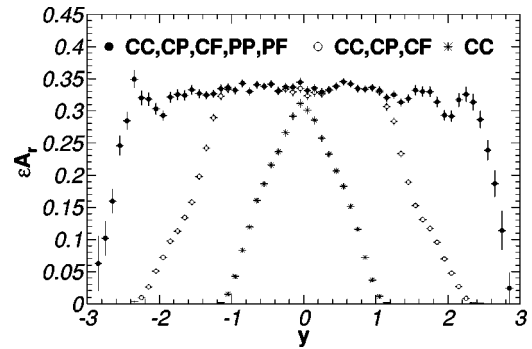


FIG. 1. The efficiency times acceptance of  $e^+e^-$  pairs in the Z-boson mass region for the CC sample (asterisks), the CC+CP+CF sample (open circles), and the CC+CP+CF+PP+PF sample (solid circles).

rimeter resolution is 0.2% in the low  $y$  region and increases to 0.7% at  $|y| > 2.5$ . This error is related to the  $y$  resolution, which has an rms of  $\sim 0.025$ . The systematic error from variations in the  $\gamma^*/Z$   $P_T$  is 0.5% at  $|y|=0$  and 2.0% at  $|y|=2.8$ . The systematic error from the choice of PDFs is under 0.5% for  $|y| < 1.9$  and increases to 2% at  $|y|=2.8$ . The errors are uncorrelated. For the total cross section measurement, the combined systematic error is 0.6% (excluding the luminosity uncertainty). The  $p\bar{p}$  collision luminosity is derived with the CDF's beam-beam cross section,  $\sigma_{\text{BBC}} = 51.15 \pm 1.60$  mb [15,16]. The luminosity error of 3.9% contains the  $\sigma_{\text{BBC}}$  error and uncertainties specific to running conditions.

Figures 2(a) and 2(b) compare the measured  $d\sigma/dy$  distributions to theoretical predictions in the Z-boson and high mass regions, respectively. The predictions are LO calculations with CTEQ5L [17] PDFs and NLO [18] calculations with 1999 Martin-Roberts-Stirling-Thorne (MRST99) [19] and CTEQ5M-1 [17] PDFs. The predictions in Fig. 2(a) have

TABLE I.  $d\sigma/dy$  distribution of  $e^+e^-$  events in the mass range  $66 < M_{ee} < 116$  GeV/ $c^2$ . The first and second errors are statistical and systematic, respectively. The 3.9% luminosity error is not included. Here,  $y$  is the bin center value.

| $y$  | $d\sigma/dy$ [pb]         | $y$  | $d\sigma/dy$ [pb]         |
|------|---------------------------|------|---------------------------|
| 0.05 | $70.78 \pm 3.27 \pm 0.37$ | 1.45 | $51.56 \pm 2.82 \pm 0.45$ |
| 0.15 | $70.19 \pm 3.26 \pm 0.37$ | 1.55 | $43.51 \pm 2.56 \pm 0.39$ |
| 0.25 | $72.03 \pm 3.30 \pm 0.41$ | 1.65 | $41.62 \pm 2.54 \pm 0.45$ |
| 0.35 | $72.18 \pm 3.30 \pm 0.45$ | 1.75 | $41.24 \pm 2.53 \pm 0.46$ |
| 0.45 | $69.90 \pm 3.20 \pm 0.50$ | 1.85 | $35.15 \pm 2.35 \pm 0.40$ |
| 0.55 | $72.86 \pm 3.25 \pm 0.59$ | 1.95 | $29.72 \pm 2.23 \pm 0.36$ |
| 0.65 | $73.68 \pm 3.27 \pm 0.64$ | 2.05 | $23.80 \pm 1.98 \pm 0.33$ |
| 0.75 | $64.64 \pm 3.09 \pm 0.58$ | 2.15 | $18.04 \pm 1.68 \pm 0.29$ |
| 0.85 | $66.59 \pm 3.13 \pm 0.63$ | 2.25 | $18.47 \pm 1.70 \pm 0.33$ |
| 0.95 | $64.59 \pm 3.10 \pm 0.67$ | 2.35 | $9.35 \pm 1.17 \pm 0.19$  |
| 1.05 | $64.04 \pm 3.09 \pm 0.70$ | 2.45 | $8.20 \pm 1.17 \pm 0.19$  |
| 1.15 | $63.11 \pm 3.10 \pm 0.67$ | 2.55 | $3.24 \pm 0.80 \pm 0.08$  |
| 1.25 | $56.26 \pm 2.92 \pm 0.55$ | 2.65 | $2.15 \pm 0.85 \pm 0.06$  |
| 1.35 | $51.57 \pm 2.82 \pm 0.45$ | 2.75 | $2.14 \pm 1.02 \pm 0.06$  |

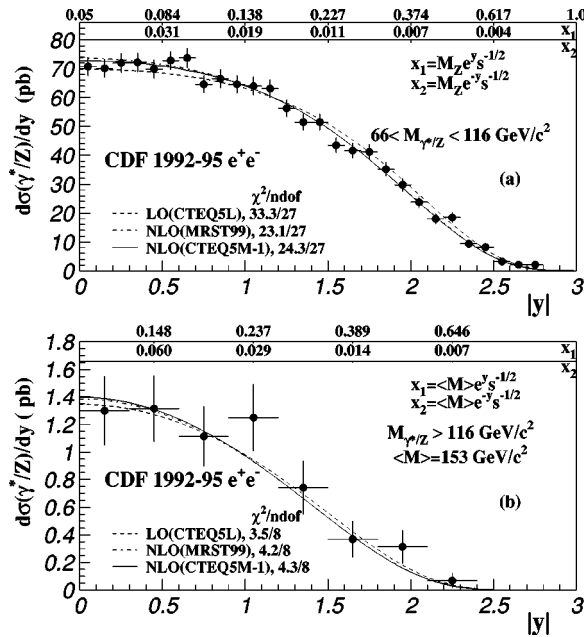


FIG. 2.  $d\sigma/dy$  distribution of  $e^+e^-$  pairs in: (a) the Z-boson mass ( $66 < M_{ee} < 116 \text{ GeV}/c^2$ ) region, and (b) the high mass ( $M_{ee} > 116 \text{ GeV}/c^2$ ) region. The error bars on the data include statistical errors only. The theoretical predictions have been normalized to the data in the Z-boson mass region. The top horizontal axes on the figures are the corresponding values of  $x_1$  and  $x_2$  as a function of  $y$ . The  $M$  used to obtain  $x_1$  and  $x_2$  in (b) is the mean mass over the bin.

been normalized by a factor  $F$ , the ratio of the measured total cross section to the prediction ( $F=1.51, 1.14$ , and  $1.13$  for the CTEQ5L, MRST99 and CTEQ5M-1 PDFs, respectively). The predictions in Fig. 2(b) are normalized with the factor  $F$  from the Z-boson mass region. We compare the data to the theory using statistical errors only. As the  $\chi^2$  values listed in Fig. 2(a) indicate, the LO calculation for the Z-boson mass region does not fit the shape as well as the NLO calculations. The LO calculation with CTEQ3L PDFs, the basis of the acceptance calculation, fits the Z-boson mass region data well: it has a  $\chi^2/\text{ndof}=21.2/27$  ( $F=1.55$ ).

Model independent measurements of the total production cross sections for  $e^+e^-$  pairs are extracted by integrating the measured values of  $d\sigma/dy$ . Because there are no data for  $|y| > 2.8$ , we use a NLO calculation with the CTEQ5M-1 PDFs to estimate  $\gamma^*/Z$  production in that region. The cross section in the region of  $|y| > 2.8$  is about 0.02% of the predicted total cross section. The extracted cross section for the Z-boson mass region is  $252.1 \pm 3.9$  (stat)  $\pm 1.6$  (syst)  $\pm 9.8$  (lum) pb. The corresponding  $\sigma(p\bar{p} \rightarrow Z) \cdot Br(Z$

$\rightarrow ee$ ) is  $253 \pm 4$  (stat+syst)  $\pm 10$  (lum) pb. This measurement is in agreement with our previous measurements in the dielectron [9] [ $248 \pm 5$  (stat+syst)  $\pm 10$  (lum) pb, using only the CC+CP+CF  $e^+e^-$  sample] and dimuon [2] [ $237 \pm 9$  (stat+syst)  $\pm 9$  (lum) pb, using a CC sample] channels. These previous measurements use a QCD model to correct for the missing events at high rapidity. The combined  $e^+e^-$  and  $\mu^+\mu^-$  cross section is  $250 \pm 4$  (stat+syst)  $\pm 10$  (lum) pb. Since the  $p\bar{p}$  inelastic cross section used by the CDF in luminosity calculations differs from DØ's by +5.9% [16], measured cross sections must be renormalized before comparisons to other experiments. The DØ measurement [20] of  $221 \pm 11$  pb, when renormalized to the CDF luminosity measurement, is  $234 \pm 12$  pb.

The total cross section measurements can also be compared to QCD calculations. Fixed order QCD calculations have uncertainties from PDF measurements and corrections from higher orders of perturbation theory, i.e., the  $K$ -factor. The NLO-to-LO total cross section correction is significant:  $K \sim 1.4$ . The NNLO cross section [6], calculated with recent estimates of NNLO PDFs [7,21], is 230 pb [21], or  $\sim 4\%$  larger than the NLO cross section. Using the MRST99 NLO PDFs, the NNLO calculation gives a cross section of  $227 \pm 9$  pb, where the 4% error is mostly from uncertainties [19] in the NLO PDFs. Given these uncertainties, the theoretical expectation is consistent with the Z cross section measurements.

The measurement of the Drell-Yan total cross section in the high mass region is  $4.0 \pm 0.4$  (stat+syst)  $\pm 0.2$  (lum) pb. The corresponding prediction of the total cross section from the NNLO QCD theory using MRST99 PDFs is 3.3 pb.

In summary, the rapidity distributions of  $e^+e^-$  pairs in the Z-boson and high mass regions have been measured for the first time over nearly the entire kinematic region. This measurement uses an improved track finder in the high rapidity region to reduce the background and the uncertainties associated with it. In addition, unlike the previous measurements of the total cross section, this measurement is model independent.

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- component of a particle's momentum,  $P_T = P \sin \theta$ , and  $E_T = E \sin \theta$ , where  $E$  is the energy measured in the calorimeter.
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