January 2008

Search for supersymmetry in di-photon final states at √s = 1.96 TeV

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Available online at www.sciencedirect.com
ScienceDirect
Physics Letters B 659 (2008) 856–863
PHYSICS LETTERS B
www.elsevier.com/locate/physletb

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doi:10.1016/j.physletb.2007.12.006

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Received 25 October 2007; received in revised form 7 December 2007; accepted 9 December 2007

Available online 15 December 2007

Editor: L. Rolandi

Abstract

We report results of a search for supersymmetry (SUSY) with gauge-mediated symmetry breaking in di-photon events collected by the D0 experiment at the Fermilab Tevatron Collider in 2002–2006. In 1.1 fb^{-1} of data, we find no significant excess beyond the background expected from the standard model and set the most stringent lower limits to date for a standard benchmark model on the lightest neutralino and chargino masses of 125 GeV and 229 GeV, respectively, at 95% confidence.

PACS: 14.80.Ly; 12.60.Jv; 13.85.Rm

Low-scale SUSY is one of the most promising solutions to the hierarchy problem associated with the intrinsic disparity between the electroweak and Planck scales. It postulates that for each known particle there exists a superpartner, thereby stabilizing the radiative corrections to the Higgs boson mass. Bosons have fermion superpartners, and vice versa. None of the superpartners have yet been observed, and superpartner masses must therefore be much larger than those of their partners, *i.e.*, SUSY is a broken symmetry. Experimental signatures of supersymmetry are determined through the manner and scale of SUSY breaking. In models with gauge-mediated supersymmetry breaking (GMSB) [1,2], it is achieved through the introduction of new chiral supermultiplets, called messengers that couple to the ultimate source of supersymmetry breaking and to the SUSY particles. At colliders, assuming R-parity conservation [3], superpartners are produced in pairs ($\tilde{\chi}_1^+\tilde{\chi}_1^-$ and $\tilde{\chi}_1^\pm\tilde{\chi}_2^0$ production dominates in most cases) and decay to the standard model particles and next-to-lightest SUSY particle (NLSP), which can be either a neutralino or a slepton. In the former case, which is considered in this note, the NLSP decays into a photon and a gravitino (the lightest superpartner in GMSB SUSY models, with mass less than ≈ 1 keV). The gravitino is stable, and escapes detection, creating an apparent imbalance in transverse momentum ($\not\!\!E_T$) in the event. GMSB SUSY final states are therefore characterized by two energetic photons and large missing transverse momentum. The differences in event

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kinematics between particular GMSB SUSY models result in slightly different experimental sensitivities [4], and to obtain a quantitative measure of limits on SUSY we consider a model referred to as "Snowmass Slope SPS 8" [5]. This model has only a single dimensioned parameter: an energy scale Λ that determines the effective scale of SUSY breaking. The minimal GMSB parameters correspond to a messenger mass $M_m = 2\Lambda$, the number of messengers $N_5 = 1$, the ratio of the vacuum expectation values of the two Higgs fields $\tan\beta = 15$, and the sign of the Higgsino mass term $\mu > 0$. The neutralino lifetime is not defined within the model. For this analysis, it is assumed to be sufficiently short to yield decays with prompt photons.

Searches for GMSB SUSY were carried out by collaborations at the CERN LEP collider [6] and at the Fermilab Tevatron collider in both Run I [7] and early in Run II [4,8]. The initial limits from CDF and D0 for Run II, based on the SPS 8 model, were combined [9] to yield $\Lambda > 84.6$ TeV corresponding to the limit on the chargino mass of 209 GeV, at 95% confidence. Complementary searches for GMSB SUSY with R-parity violation were performed by the H1 experiment at HERA [10].

This analysis is an update of that described in Ref. [4], with about a factor of three more data and improved photon identification based on: (i) an electromagnetic (EM) cluster "pointing" algorithm that predicts the origin of a photon with a resolution of about 2 cm along the beam axis, thereby eliminating the largest instrumental background associated with misreconstruction of the primary interaction vertex, and (ii) an improved track veto requirement that suppresses sources of background with electrons in the final state. We also use an improved likelihood fitter [11] to set limits on the scale parameter Λ.

The data in this analysis were recorded using single EM triggers with the D0 detector [12], the main components of which are an inner tracker, liquid-argon/uranium calorimeters, and a muon spectrometer. The inner tracker consists of silicon microstrip and central scintillating-fiber trackers located in a 2 T superconducting solenoidal magnet, providing measurements up to pseudorapidities[8] of $|\eta| \approx 3.0$ and $|\eta| \approx 1.8$, respectively. The calorimeters are finely segmented and consist of a central section (CC) covering $|\eta| < 1.2$ and two endcap calorimeters extending coverage to $|\eta| \approx 4$, all housed in separate cryostats [13]. The electromagnetic section of the calorimeter has four longitudinal layers and transverse segmentation of 0.1×0.1 in η–ϕ space (where ϕ is the azimuthal angle), except in the third layer, where it is 0.05×0.05. The central preshower (CPS) system is placed between the solenoid and the calorimeter cryostat and covers $|\eta| \lesssim 1.2$. The CPS provides precise measurement of positions of EM showers. The axes of EM showers are reconstructed by fitting straight lines to shower positions measured in the four longitudinal calorimeter layers and the CPS (EM "pointing"). The data for this study were collected between 2002 and summer 2006, using inclusive single EM triggers that are almost 100% efficient to select signal data. The integrated luminosity [14] of the sample is 1100 ± 70 pb^{-1}.

[8] Pseudorapidity is defined as $-\log(\tan(\frac{\theta}{2}))$, where θ is the angle between the particle and the proton beam direction.

Photons and electrons are identified based on reconstructed EM clusters using calorimetric information and further classified into electron and photon candidates, based on tracking information. The EM clusters are selected from calorimeter clusters using the simple cone method (of radius $\mathcal{R} = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} = 0.4$) by requiring that (i) at least 90% of the energy is deposited in the EM section of the calorimeter, (ii) the calorimeter isolation variable $I = [E_{\text{tot}}(0.4) - E_{\text{EM}}(0.2)]/E_{\text{EM}}(0.2)$ is less than 0.07, where $E_{\text{tot}}(0.4)$ is the total shower energy in a cone of radius $\mathcal{R} = 0.4$, and $E_{\text{EM}}(0.2)$ is the EM energy in a cone of radius $\mathcal{R} = 0.2$, (iii) the transverse, energy-weighted, width of the EM cluster in the third EM calorimeter layer is smaller than 0.04 rad, and (iv) the scalar sum of the transverse momenta (p_T) of all tracks originating from the primary vertex in an annulus of $0.05 < \mathcal{R} < 0.4$ around the cluster is less than 2 GeV. The isolation criteria are tuned so that photons that convert in the tracker material are not rejected. The EM cluster is further defined as an electron candidate if it is spatially matched to activity in the tracker, and as a photon candidate otherwise. The tracker activity can be either a reconstructed track or a density of hits in the silicon microstrip and central fiber trackers consistent with a track, i.e., an electron. The latter requirement allows for increasing electron track-matching efficiency, ϵ_{trk}, measured in $Z \to ee$ data, from $(93.0 \pm 0.1)\%$ to $(98.6 \pm 0.1)\%$ by identifying electrons with lost tracks due to hard bremsstrahlung and/or inefficiency of the inner trackers. This reduces electron backgrounds to photons by a factor of five, while keeping the efficiency of anti-track activity requirement high. We measure that $(91 \pm 3)\%$ of photon candidates in $Z \to ee\gamma$ data satisfy the anti-track activity requirement.

Jets are reconstructed using the iterative, midpoint cone algorithm [15] with a cone size of $\mathcal{R} = 0.5$. The missing transverse energy is determined from the energy deposited in the calorimeter for $|\eta| < 4$ and is corrected for the EM and jet energy scales.

We select $\gamma\gamma$ candidates by requiring events to have two photon candidates, each with transverse energy $E_T > 25$ GeV identified in the CC with $|\eta| < 1.1$. We require that at least one of the photon candidates be matched to a CPS cluster, and that the primary vertex be consistent with that of the photon candidate (obtained from the EM pointing). The accuracy of the determination of the photon vertex is measured using photons from final state radiation in $Z \to ee\gamma$ data sample and found to be 2.3 ± 0.3 cm. The requirement of consistency between the photon and primary vertices ensures correct calculation of the transverse energies and tracking isolation requirements. The accuracy of primary vertex association is studied in GMSB SUSY Monte Carlo simulated events, where the primary vertex is identified correctly in $(98.5 \pm 0.1)\%$ of the events while the photon vertex matches the primary vertex in $(95.8 \pm 0.1)\%$.

To reduce potential bias in the measurement of $\not\!\!E_T$ from mismeasurement of jet transverse momentum, we also require that the jet with the highest E_T (if jets are present in the event) be separated from the $\not\!\!E_T$ in azimuth by no more than 2.5 radians. This selection yields 2341 events (the $\gamma\gamma$ sample).

All instrumental backgrounds arise from standard model processes, with either genuine $\not\!\!E_T$ ($W\gamma$, $W+$jet, and $t\bar{t}$ production) or without inherent $\not\!\!E_T$ (direct photon, multi-jet, and $Z \to ee$ production). All these backgrounds are measured using data.

The former source always has an electron in the final state which is misidentified as a photon. The contribution of this background to the $\not\!\!E_T$ distribution in data can be estimated using an $e\gamma$ sample (selected by requiring an electron and a photon candidate and using the same kinematical requirements as for the $\gamma\gamma$ sample) scaled by the probability of an electron–photon misidentification which is measured using $Z \to ee$ data. First, the $\not\!\!E_T$ distribution in the $e\gamma$ sample must be corrected for the contribution from events with no real $\not\!\!E_T$. The contribution from Drell–Yan events is taken into account by obtaining the $\not\!\!E_T$ distribution for the ee sample (selected by requiring two electron candidates and applying the same kinematical requirements as for the $\gamma\gamma$ sample) which is dominated by Drell–Yan events. The Drell–Yan $\not\!\!E_T$ distribution is further normalized to the number of Z boson events in the $e\gamma$ sample (the latter is determined by fitting the $e\gamma$ invariant mass spectrum to the Z boson mass peak).

The contribution from the multi-jet processes is estimated from a data sample (referred to as the *QCD* sample) selected by requiring two EM clusters that (a) satisfy all the kinematic selection used to select $\gamma\gamma$ sample and (b) satisfy all the photon identification criteria but fail the shower shape requirement. The $\not\!\!E_T$ distribution in the *QCD* sample is normalized to the number of the events in the $e\gamma$ sample with $\not\!\!E_T < 12$ GeV after subtraction of the Drell–Yan contribution as determined above. The expected number of $W\gamma$, $W+$jet, and $t\bar{t}$ events with $\not\!\!E_T < 12$ GeV is negligible.

After the Drell–Yan and multi-jet contributions to the $e\gamma$ sample are subtracted, the resulting $\not\!\!E_T$ distribution is scaled by $(1-\epsilon_{\text{trk}})/\epsilon_{\text{trk}}$, where ϵ_{trk} is the efficiency of the track-matching requirement to obtain the estimate of $\not\!\!E_T$ distribution for the background with genuine $\not\!\!E_T$.

The background from events with no inherent $\not\!\!E_T$ is divided into multi-jet events with two real isolated photons and events where one or both photons are misidentified jets. Since the $\not\!\!E_T$ resolution for both sources is dominated by the photon energy resolution, the $\not\!\!E_T$ distributions for the two sources are very similar. However, misidentified jets have a different energy response compared with that of real photons which leads to a slight difference in the shapes of the $\not\!\!E_T$ distributions. For the real di-photon events, the $\not\!\!E_T$ is assumed to have the same shape as that of the Drell–Yan events. For misidentified jets, the shape of the $\not\!\!E_T$ distribution is taken from the *QCD* sample. Relative normalization of the two sources is obtained using a fit to the $\not\!\!E_T$ distribution in the $\gamma\gamma$ sample. We check that the fit is not sensitive to possible signal contribution, and cross-check with a method that estimates the $\gamma\gamma$ sample purity using the measured shower shape in the CPS. The relative fraction of di-photons is $(60 \pm 20)\%$ and this uncertainty is propagated as a systematic uncertainty for the limit setting. Absolute normalization of the $\not\!\!E_T$ distributions from both sources is determined so that the number of events with $\not\!\!E_T < 12$ GeV matches that in the $\gamma\gamma$ sample.

The largest physics backgrounds are from $Z\gamma\gamma \to \nu\nu\gamma\gamma$ and $W\gamma\gamma \to \ell\gamma\gamma\nu$ processes. Contributions from these backgrounds are estimated as 0.15 ± 0.06 and 0.10 ± 0.04 events, respectively, using CompHep [16] Monte Carlo simulation, cross-checked with MadGraph [17]. The contribution of these backgrounds to the $\not\!\!E_T$ distribution is taken from Monte Carlo simulation, with number of events normalized to the integrated luminosity of the data sample.

The $\not\!\!E_T$ distribution for the $\gamma\gamma$ sample, with contributions from physics background ($W/Z+\gamma\gamma$), and instrumental background with genuine $\not\!\!E_T$ (processes with mis-identified electrons) and no inherent $\not\!\!E_T$ ($\gamma\gamma$ and multi-jet) is given in Fig. 1. We also illustrate the $\not\!\!E_T$ distribution expected from GMSB SUSY for two values of Λ. The number of observed events, as well as expected background and signal from GMSB SUSY for $\not\!\!E_T > 30$ GeV and > 60 GeV are given in Table 1.

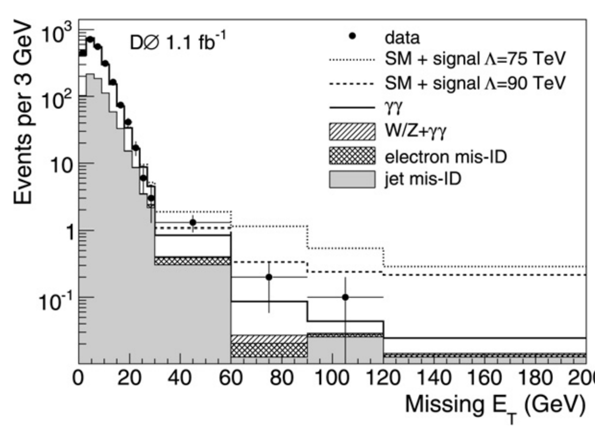

Fig. 1. The $\not\!\!E_T$ distribution in $\gamma\gamma$ data with $W/Z+\gamma\gamma$ background (hatched histogram), instrumental background with no genuine $\not\!\!E_T$: $\gamma\gamma$ (solid black line) and multi-jet (filled histogram), and background from processes with genuine $\not\!\!E_T$ and a misidentified electron (cross-hatched histogram). The expected $\not\!\!E_T$ distributions if GMSB SUSY events were present are shown as dotted and dashed lines.

Table 1
Numbers of background events from $W\gamma$, $W+$jet, and $t\bar{t}$ (Genuine $\not\!\!E_T$), no inherent $\not\!\!E_T$ (No $\not\!\!E_T$), $Z\gamma\gamma \to \nu\nu\gamma\gamma$ and $W\gamma\gamma \to \ell\gamma\gamma\nu$ (Physics) processes; the total number of expected background events; numbers of expected GMSB SUSY signal events for two values of Λ; and the observed numbers of events for $\not\!\!E_T > 30$ GeV and 60 GeV. Errors are statistical and systematic combined

	Background events				Expected signal events		Observed events
	Genuine $\not\!\!E_T$	No $\not\!\!E_T$	Physics	Total	$\Lambda = 75$ TeV	$\Lambda = 90$ TeV	
$\not\!\!E_T > 30$ GeV	0.97 ± 0.12	9.62 ± 1.12	0.19 ± 0.07	10.8 ± 1.1	28.3 ± 4.2	8.7 ± 1.3	16
$\not\!\!E_T > 60$ GeV	0.11 ± 0.04	1.44 ± 0.43	0.08 ± 0.04	1.6 ± 0.4	18.1 ± 2.7	6.4 ± 1.0	3

Table 2
Points on the GMSB Snowmass Slope model: neutralino and chargino masses, cross sections predicted by PYTHIA, k-factors, and reconstruction efficiencies with total uncertainty

Λ, TeV	$m_{\tilde{\chi}_1^0}$, GeV	$m_{\tilde{\chi}_1^+}$, GeV	σ^{LO}, fb	k-factor	Efficiency
70	93.7	168.2	215	1.21	0.17 ± 0.03
75	101.0	182.3	148	1.20	0.18 ± 0.03
80	108.5	198.1	97.5	1.19	0.18 ± 0.03
85	115.8	212.0	65.4	1.18	0.19 ± 0.03
90	123.0	225.8	41.8	1.17	0.19 ± 0.03
95	130.2	239.7	29.5	1.16	0.20 ± 0.03
100	137.4	253.4	20.6	1.15	0.20 ± 0.03
105	144.5	267.0	14.4	1.14	0.18 ± 0.03
110	151.7	280.7	10.3	1.13	0.19 ± 0.03

The expected GMSB signal efficiency is estimated from Monte Carlo simulation generated for several points on the Snowmass Slope (see Table 2), covering the neutralino mass range from 170 GeV to 280 GeV. Although $\tilde{\chi}_1^+\tilde{\chi}_1^-$ and $\tilde{\chi}_1^\pm\tilde{\chi}_2^0$ processes dominate, we consider all GMSB SUSY production channels. We used ISAJET 7.58 [18] to determine SUSY interaction eigenstate masses and couplings. PYTHIA 6.319 [19] is used to generate the events after determining the sparticle masses, branching fractions and leading order (LO) production cross sections using CTEQ6L1 parton distributions [20]. The generated events are processed through a full GEANT-based [21] detector simulation and the same reconstruction code as used for data. The LO signal cross sections are scaled to match the next-to-leading order (NLO) prediction using k-factor values (see Table 2), extracted from Ref. [22].

The systematic error on the expected number of signal events comes from the uncertainties in photon identification efficiency (10%), statistics in MC samples (5%), track veto requirement (3%), and trigger efficiency (4%). These were obtained using $Z \to e^+e^-$ and $Z \to e^+e^-\gamma$ decays in data and in MC simulation. Variation of parton distribution functions and uncertainty in the total integrated luminosity result in additional 4% and 6.1% errors in signal yield respectively. The total uncertainty on the background is dominated by statistics.

As the observed number of events for all values of $\not{E}_T$ is in good agreement with the standard model prediction, we conclude that there is no evidence for GMSB SUSY in the data. We set limits on the production cross section by utilizing a likelihood fitter [11] that incorporates a log-likelihood ratio (LLR) test statistic method. This method utilizes binned $\not{E}_T$ distributions rather than a single-bin (fully-integrated) value, and therefore accounts for the shapes of the distributions, leading to greater sensitivity. The value of the confidence level for the signal CL_s is defined as $CL_s = CL_{s+b}/CL_b$, where CL_{s+b} and CL_b are the confidence levels for the signal plus background hypothesis and the background-only (null) hypothesis, respectively. These confidence levels are evaluated by integrating corresponding LLR distributions populated by simulating outcomes via Poisson statistics. Systematic uncertainties are treated as uncertainties on the expected numbers of signal and background events, not the outcomes of the limit calculations. The degrading effects of systematic uncertainties are reduced

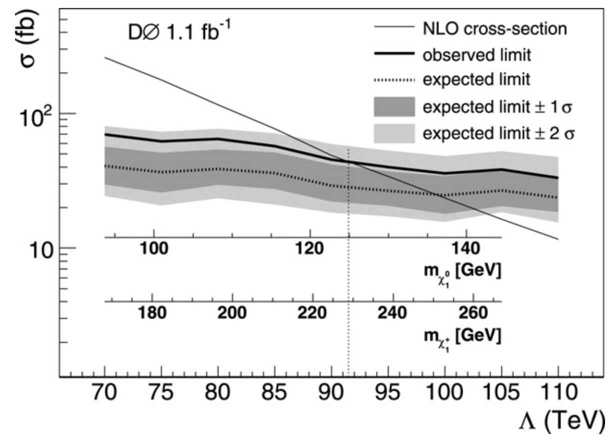

Fig. 2. Predicted cross section for the Snowmass Slope model versus Λ. The observed and expected 95% C.L. limits are shown in solid and dash-dotted lines, respectively.

by introducing a maximum likelihood fit to the missing transverse energy distribution. A separate fit is performed for both the background-only and signal-plus-background hypotheses for each data or pseudo-data distribution.

The limits are shown in Fig. 2 together with expected signal cross sections. The observed limits are statistically compatible with the expected limits. The observed upper limit on the signal cross section is below the prediction of the Snowmass Slope model for $\Lambda < 91.5$ TeV, or in terms of gaugino masses, $m_{\tilde{\chi}_1^0} < 125$ GeV and $m_{\tilde{\chi}_1^+} < 229$ GeV. These represent the most stringent limits on this particular GMSB SUSY model to date.

Acknowledgements

We thank the staffs at Fermilab and collaborating institutions, and acknowledge support from the DOE and NSF (USA); CEA and CNRS/IN2P3 (France); FASI, Rosatom and RFBR (Russia); CAPES, CNPq, FAPERJ, FAPESP and FUNDUNESP (Brazil); DAE and DST (India); Colciencias (Colombia); CONACyT (Mexico); KRF and KOSEF (South Korea); CONICET and UBACyT (Argentina); FOM (The Netherlands); Science and Technology Facilities Council (United Kingdom); MSMT and GACR (Czech Republic); CRC Program, CFI, NSERC and WestGrid Project (Canada); BMBF and DFG (Germany); SFI (Ireland); The Swedish Research Council (Sweden); CAS and CNSF (China); Alexander von Humboldt Foundation; and the Marie Curie Program.

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