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A broadband and polarization-independent metasurface is analyzed and designed for both monostatic and bistatic radar cross section (RCS) reduction in this paper. Metasurfaces are composed of two types of electromagnetic band-gap (EBG) lattice, which is a subarray with “0” or “π” phase responses, arranged in periodic and aperiodic fashions. A new mechanism is proposed for manipulating electromagnetic (EM) scattering and realizing the best reduction of monostatic and bistatic RCS by redirecting EM energy to more directions through controlling the wavefront of EM wave reflected from the metasurface. Scattering characteristics of two kinds of metasurfaces, periodic arrangement and optimized phase layout, are studied in detail. Optimizing phase layout through particle swarm optimization (PSO) together with far field pattern prediction can produce a lot of scattering lobes, leading to a great reduction of bistatic RCS. For the designed metasurface based on optimal phase layout, a bandwidth of more than 80% is achieved at the normal incidence for the −9.5 dB RCS reduction for both monostatic and bistatic. Bistatic RCS reduction at frequency points with exactly 180° phase difference reaches 17.6 dB. Both TE and TM polarizations for oblique incidence are considered. The measured results are in good agreement with the corresponding simulations.

1. Introduction

The reflection phase control for EM wave impinging on a surface is important in both fundamental EM researches and device technologies. Many phenomena and applications are related to the reflection phase control, such as EM wave absorber [1], high impedance surface and artificial magnetic conductor (AMC) [2, 3], reflectarray antennas [4], and beam forming and beam scanning antennas.

Reshaping the scattering pattern can effectively suppress backward reflection. Scattering energy is redirected to other directions. In 2007, Paquay et al. [5] proposed a planar structure for RCS reduction, based on a combination of artificial magnetic conductors (AMC) and perfect electric conductors (PEC) in a chessboard-like configuration. The backscattering field can be effectively canceled by redirecting it along other angles. However, the narrow in-phase reflection bandwidth of the AMC restricts the RCS reduction frequency range. In [6], a planar monolayer chessboard structure is presented for broadband radar cross section reduction using AMC technology. More than 55% frequency bandwidth with a monostatic RCS reduction larger than 10 dB is obtained and four scattered beams are produced. In another similar work [7], two-dimensional phase gradient metasurfaces of a chessboard-like configuration have been proposed for wideband RCS reduction. The −5 dB RCS reduction bandwidth of 66% is achieved for TE and TM polarizations at the normal incidence. For those rectangular checkerboard surfaces with periodic arrangement, four scattering beams are produced and bistatic RCS reduction is about 8.1 dB. In 2015, Chen et al. [8] proposed a hexagonal checkerboard surface of periodic fashion, with the −10 dB monostatic RCS reduction bandwidth of about 61%, which can create six bistatic RCS lobes, leading to further bistatic RCS reduction (8.9 dB). In [9] Yang...
et al. proposed a metasurface with nonuniform phase distribution for suppression of specular reflections, with the −7 dB RCS reduction bandwidth of about 42% for \( y \)-polarized and \( x \)-polarized incident waves. In addition, metamaterial absorber has been proposed as another way to achieve RCS reduction [10,11]. In 2013, Singh et al. proposed a broadband millimeter wave metamaterial absorber based on embedding of dual resonators [12], with an absorber at 77 GHz with a bandwidth of 8 GHz and peak absorption of greater than 98%.

Metasurfaces, a new class of metamaterials that consist of only a monolayer of planar metallic structures, have shown great promise for achieving full control of the reflected wavefront with low fabrication cost [13]. Metasurfaces are capable of generating abrupt interfacial phase changes and provide a unique way to control the local wave front at the subwavelength scale [14].

In this paper, an ultrathin nonabsorptive metasurface is designed for broadband RCS reduction. A new mechanism is proposed for manipulating EM scattering and realizing monostatic and bistatic RCSs reduction through optimizing reflection phase. The main contributions of this paper are as follows.

(1) Two unit cells, square ring patch and small circular patch, are optimized to achieve 86.8% bandwidth for about 180° (±37°) reflection phase difference.

(2) Theoretical analysis of far field RCS pattern for the metasurface is derived based on phase-only planar array synthesis.

(3) Optimizing the reflection phase layout of the metasurface using array pattern synthesis (APS) and particle swarm optimization (PSO) can create a lot of bistatic RCS lobes (≫4 or 6), leading to a great reduction of bistatic RCS.

(4) The number of bistatic RCS lobes increases as the lattice number increases. Thus, as the size of metasurface increases, the reduced magnitude of monostatic and bistatic RCSs becomes larger.

(5) The designed metasurface is polarization-independent.

(6) Specular reflection and bistatic RCS for oblique incidence with TE and TM polarization can also be effectively suppressed.

2. EBG Unit Cell Design

The variation of reflection coefficient phase is one of the unique properties of EBG structure, because the phase of the reflected field changes continuously from 180° to −180° versus frequency while, ideally, the magnitude of the reflection coefficient of a lossless ground plane is unity. We consider the metasurface consists of only two EBG structures under the normal incidence of a plane wave.

For a scatterer, scattered power is proportional to the incident power and its RCS (\(|E|^2 \sim \sigma |E|^2\)). Relative to a perfectly electrical conductor (PEC) plate with the same size, the RCS reduction of the metasurface can be expressed as

\[
\text{RCS reduction} = 10 \log \frac{\sigma_{\text{metasurface}}}{\sigma_{\text{PEC}}} = 10 \log \left( \frac{|E_{\text{metasurface}}|^2}{|E_{\text{PEC}}|^2} \right)
\]

\[
= 10 \log \left( \frac{|E_0 e^{j\varphi_1} + E_0 e^{j\varphi_2}|^2}{4 |E_0 e^{j\varphi_1}|^2} \right)
\]

\[
= 10 \log \frac{1 + \cos(\varphi_1 - \varphi_2)}{2}.
\]

In order to achieve 10 dB RCS reduction, reflection phase difference between two EBG structures needs to satisfy

\[
143° \leq |\Delta\varphi| = |\varphi_1 - \varphi_2| \leq 217°.
\]

Here, Rogers RT/duroid 5880 material with thickness \( h = 62 \) mil and dielectric constant \( \varepsilon_r = 2.2 \) (loss tangent \( \tan \delta = 0.0009 \)) has been chosen as substrate for the ultrathin metasurface. Square ring patch and small circular patch were chosen as the EBG unit cell. The reflection phase of small circular patch changes approximately linearly with frequency. For square ring patch, when the frequency is fixed, the reflection phase change with side length is relatively larger. Thus the dimensions of both unit cells can be easily optimized to achieve about 180° reflection phase difference in the largest possible frequency band. Two EBG unit cells are simulated by method of moments (MoM) based Ansys Designer using periodic boundary condition (PBC). This optimization process results in the dimensions depicted in Figure 1. The phases of the two EBG structures, as well as the phase difference between them, are plotted in Figure 2. The reflection phase of EBG structure reaches 0° at the resonant frequencies, which are 14.7 GHz and 32.1 GHz for square ring patch and small circular patch, respectively. In particular, from 15.2 GHz to 38.5 GHz (86.8% bandwidth), the phase difference ranges from 143° to 217° and is exactly 180° at 16.9 GHz and 31.1 GHz. Therefore, a broadband behavior of the proposed metasurface is expected. Here, the square ring patch is considered as “0°” element; then small circular patch is “\( \pi \)” element.

3. Far Field Pattern Prediction

Consider a metasurface under normal incidence of plane wave, which consists of a two-dimensional array of \( M \times N \) EBG lattices. The lattices are uniformly spaced with \( d_x \) in \( x \).
Direction and $d_y$ in $y$ direction. The total 2D array pattern can be expressed as [15]

$$F(\theta, \phi) = EP(\theta) \cdot AF(\theta, \phi),$$

$$EP(\theta) = \cos \frac{EF}{2} \theta,$$

$$AF(\theta, \phi) = \sum_{m=1}^{M} \sum_{n=1}^{N} e^{-j\left(\frac{2\pi}{\lambda} \sin \theta \cos \phi \cdot m d_x + \sin \phi \cdot n d_y + \theta(m, n)\right)},$$

where $EP$ and $AF$ are the element pattern and array factor, respectively. $\theta(m, n) = \theta_x(m) + \theta_y(n)$ is the initial phase of the element; $\theta$ and $\phi$ are the elevation and azimuth angles of an arbitrary direction, respectively. A good expression for modeling the element pattern is a cosine function raised to a power that is called the element factor (EF).

Array factor can be separated as

$$AF(\theta, \phi) = AF_U(\theta, \phi) \cdot AF_V(\theta, \phi)$$

$$= \sum_{m=1}^{M} e^{-j\frac{2\pi}{\lambda} m d_U U/\lambda + \theta_U(m)} \cdot \sum_{n=1}^{N} e^{-j\frac{2\pi}{\lambda} n d_V V/\lambda + \theta_V(n)},$$

where $U = \sin \theta \cos \varphi$ and $V = \sin \theta \sin \varphi$. $AF_U$ and $AF_V$ are periodic functions with the periods $P_U = \lambda/d_x$ and $P_V = \lambda/d_y$, respectively.

The equation for the directive gain is

$$D(\theta, \phi) = \frac{4\pi F^2(\theta, \phi)}{\int_{0}^{2\pi} \int_{0}^{\pi} F^2(\theta, \phi) \sin \theta \, d\theta \, d\phi},$$

Suppose the impinging plane wave has a power density $w_i$; then the intercepted EM energy by metasurface is $A w_i$. $A = MN d_x d_y$ is the aperture area [16]. The scattered power density is

$$w_s(\theta, \phi) = \frac{A w_i D(\theta, \phi)}{4\pi R^2},$$

where $R$ is the far field distance. Therefore, the far field RCS pattern of the metasurface can be calculated by

$$\sigma(\theta, \phi) = A \cdot D(\theta, \phi).$$

In this paper, the maximum RCS was defined as follows:

$$\text{maximum RCS} = \max_{\theta, \phi} [\sigma(\theta, \phi)].$$

Maximum RCS is utilized to evaluate the suppression degree of bistatic RCS.

4. Periodic Arrangement

Nonabsorptive RCS reduction of specular surface can be achieved by cancellation of the reflected field. This can be
accomplished by covering the object by the metasurface which consisted of “0” and “π” lattices in a periodic fashion [5–8]. A general rectangular metasurface with the dimension of 436.8 mm × 312 mm has been fabricated with LPKF ProtoLaser. It is composed by 14 × 10 lattices. Each lattice is a subarray composed by 12 × 12 “0” or “π” elements. “0” and “π” lattices are arranged periodically as shown in Figure 3(a). In this paper, full-wave simulation of the full structure was carried by CST Microwave Studio.

The monostatic RCS reduction was presented at Figure 4. The metasurface and the equal-sized PEC ground were measured at x-polarized wave normal incidence. The −10 dB measured bandwidth is 89% (from 14.8 to 38.5 GHz) for the metasurface with periodic arrangement. The tendency of simulated result is similar to the measurements. The maximum RCS reduction has been obtained at 17 GHz with a value around −35 dB. Note that the simulated bandwidth has been calculated by reflection coefficient. The achieved bandwidth is in good agreement with the predicted bandwidth based on the 180° phase difference between two unit cells as shown in Figure 2.

The 3D scattering patterns of the periodic arrangement for normal incidence at 17 GHz are presented in Figure 5. For the simulation result, the normalization value is 20.77 dB. The EM wave scatters along four diverging directions. Hence, in the incident direction the scattered field was greatly suppressed, and monostatic RCS was effectively reduced. The analytical predictions are in good agreement with the full-wave simulations. Theoretical directions of four scattered beams can be found in [6, 8]. The achieved bandwidth is 89%, which is in good agreement with the predicted bandwidth based on the 180° phase difference between two unit cells as shown in Figure 2.

5. Optimized Phase Layouts

From the foregoing, periodic arrangement produces four or six scattered beams and can greatly reduce the monostatic RCS [5–8]. However, bistatic RCS reduction is not very obvious and is −8.1 and −8.9 dB for the rectangular and hexagonal surfaces, respectively. Optimizing the lattice layouts aims to redirect EM scattering to more directions and further reduce the bistatic RCS based on the energy conservation.

We consider a rectangular metasurface with the same size as the metasurface of periodic phase arrangement that contains 14 × 10 equal-sized lattices in which each lattice is occupied by a subarray of “0” or “π” elements. The distribution of “0” and “π” lattices can be arbitrary. For simplicity, the planar array was considered to be the combination of two separable linear arrays; the optimized phase sequences can be designed independently along both the horizontal and vertical directions. The main aim of the study is to find a set of phase sequences in order to produce a scattering pattern with lowest possible value of maximum lobe level (MLL) and deep null placed in monostatic direction.

To get best performance of the metasurface, a particle swarm optimization (PSO) algorithm is utilized together with far field pattern prediction [17] to achieve the optimal phase arrangement. $d_x$ and $d_y$ are both set to 2λ. Accordingly, the periods of array factor are $P_{1x} = P_{1y} = 0.5$. The RCS pattern of one lattice is shown in Figure 6. As electrical size of the lattice composed of 12 × 12 “0” or “π” elements is relatively larger, in order to exactly model the scattering pattern of one lattice the EF is set to be 15. Fast Fourier transformation (FFT) is used to
Figure 5: Scattering patterns of the metasurface with periodic arrangement at 17 GHz. (a) The normalized analytical results calculated by (3). (b) The full-wave simulation results.

Figure 6: RCS pattern of one lattice composed by 12 × 12 "0" or "π" elements.

accelerate the calculation of array factor. The score function is defined as

\[
\text{score} = \text{bistatic RCS} + k \cdot \text{(monostatic RCS)},
\]

where \( k \) is weighting factor that can balance the optimization results for both monostatic and bistatic. Through numerical test, we can find that the best optimization results can be obtained when \( k = 0.25 \).

The optimized scattering patterns within the main period for lattice number of 14 and 10 are shown in Figure 8. Two corresponding phase sequences are "0π00π0π0πππ0" and "00π0π0π0π", respectively. Figure 7 depicts the convergence curves of the PSO optimization. The flat scattering pattern means that scattered fields are suppressed in all directions, so that bistatic RCS can be effectively reduced. There is deep null at the direction \( \theta = 0 \); thus the monostatic RCS can be greatly reduced. The main beam is suppressed and the EM energy is dispersed to the other direction. Our proposed physical mechanism follows the law of conservation of energy. As shown in Figure 8, the number of maximum points of the optimized results along the horizontal and vertical directions is 10 and 8, respectively. Thus, in theory, the number of scattering beams is 80 within the main period. The scattering patterns are both symmetrical along the horizontal and vertical directions.

For a uniformly excited \( N \)-element array, the array factor pattern in one period has \( N - 1 \) lobes including main beam, according to array theory. It means that a flat scattering pattern with more scattering beams can be synthesized by optimizing the phase arrangement as lattice number \( (N) \) increases. Thus, as the size of the metasurface increases, the magnitude of bistatic RCS reduction becomes larger. It is worth mentioning that periodic phase arrangement produces a fixed number of scattering beams. Four and six fixed scattering beams are created for square and hexagonal AMC/EBG surfaces [5–8], respectively. So the magnitude of maximum RCS reduction is also fixed and does not vary with the size changes.

To illustrate the above physical mechanism, we manufacture a metasurface based on two optimized phase sequences.
as shown in Figure 9. The simulated and measured results of monostatic and bistatic RCSs reductions under $x$-polarized and $y$-polarized normal incidences are presented in Figures 10 and 11, respectively. The same results of monostatic and bistatic RCSs reductions for $x$-polarization and $y$-polarization exhibit the polarization-independent feature of the metasurface. Due to the symmetric geometry of EBG unit cell, its phase responses are the same for both $x$-polarized and $y$-polarized waves; therefore metasurface is polarization-independent. As shown, more than 80% frequency bandwidth with both monostatic and bistatic RCSs reductions larger than 9.5 dB is obtained by optimized phase layout. For monostatic RCS reduction shown in Figure 10, the measurements agree well with the simulated predictions. As shown in Figure 11, the bistatic RCS for periodic arrangement reduces at the level of about $-8.1$ dB, which is analogous to previous results reported in [7]. However, bistatic RCS can be further greatly reduced in broadband by optimized phase layout. Particularly at the frequency with exact $180^\circ$ phase difference, bistatic RCS reduction reaches $17.6$ dB.

The 3D RCS patterns of the metasurface for $x$-polarization at $17$ GHz are shown in Figures 12 and 13. Nearly the same results, which are not included here to minimize duplication, were obtained for $y$-polarization. As shown, scattering energy...
focuses on the main period and its vicinity, because the scattering energy away from the main period is weakened by the lattice pattern. The scattered fields are suppressed in low levels in all directions, because the optimized phase layout has been designed to redirect the intercepted EM energies to more directions. A lot of bistatic RCS lobes are produced in space. For frequency point at 17 GHz, the lattice space $d_x = d_y = 1.768 \lambda_0$. Accordingly, the periods of array factors $AF_U$ and $AF_V$ are $P_U = P_V = 0.566$. As shown in Figure 13, the scattering pattern is symmetric about $xoz$ plane and $yoz$ plane. The spatial distribution of scattering beams of the main period is consistent with far field pattern predictions shown in Figure 8. 80 scattering beams are created within the main period, so in the whole space more than 80 scattering beams are produced, leading to a great reduction of bistatic RCS.

Under normal incidence, the RCS is dramatically reduced along the principal planes $(xoz, yoz)$, as shown in Figures 14 and 15.

The RCS pattern in $\varphi = 46.5^\circ$ and $133.5^\circ$ plane, where the maximum lobe appears at 17 GHz, is shown in Figure 16. The maximum RCS of the metasurface is reduced by 17.6 dB, 28.9 dB, and 11.3 dB in $\varphi = 0^\circ$, $\varphi = 46.5^\circ$, and $\varphi = 133.5^\circ$ plane, respectively.
Table 1: Comparisons between the results of this work and earlier checkerboard surfaces.

<table>
<thead>
<tr>
<th></th>
<th>[6]</th>
<th>[7]</th>
<th>[8]</th>
<th>This work</th>
</tr>
</thead>
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<tr>
<td>NSL</td>
<td>4</td>
<td>4</td>
<td>6</td>
<td>&gt;80</td>
</tr>
<tr>
<td>MRR (dB)</td>
<td>−8.1</td>
<td>−8.1</td>
<td>−8.9</td>
<td>−17.6</td>
</tr>
<tr>
<td>BMRR</td>
<td>40.2% (&lt;−10 dB)</td>
<td>66.0% (&lt;−6 dB)</td>
<td>63.0% (&lt;−10 dB)</td>
<td>80% (&lt;−9.5 dB)</td>
</tr>
<tr>
<td>BBRR</td>
<td>—</td>
<td>66.0% (&lt;−6 dB)</td>
<td>—</td>
<td>80% (&lt;−9.5 dB)</td>
</tr>
</tbody>
</table>

NSL: the number of scattering lobes.
MRR: maximum RCS reduction.
BMRR: the bandwidth of monostatic RCS reduction.
BBRR: the bandwidth of bistatic RCS reduction.

which is much lower than that (8.9 dB) of the hexagonal EBG surface of periodic phase arrangement [8].

Scattering characteristics including the number of scattering beams, the magnitude of maximum RCS reduction, and the bandwidth of both monostatic and bistatic RCS reductions, are compiled together in Table 1 under normal incidence. Comparing our results to recent research findings suggests that the proposed design approach has overwhelming advantages, especially in wideband suppression of bistatic RCS.

For the oblique incidence with different polarization, the RCS reductions in specular direction are given in Figure 17. Incident angle and polarization of plane wave result in the change of reflection phase difference between two unit cells. For TE polarization, as the incident angle increases, the operating frequency band shifts to the higher frequency. For TM polarization, as the incident angle increases, the high frequency band shifts to the higher frequency and the magnitude of specular RCS reduction becomes smaller at center frequency. The bistatic RCS patterns of the metasurface for oblique incidence with TE polarization at 17 GHz are shown in Figure 18. The optimized phase arrangement has been designed to redirect the intercepted EM energies to more directions. The specular reflections will no longer dominate within the whole scattered waves.

6. Conclusion

A broadband ultrathin nonabsorptive metasurface is analyzed and designed for broadband monostatic and bistatic RCSs reductions in this paper. A new mechanism is proposed for manipulating EM scattering and realizing the best reduction of monostatic and bistatic RCSs through controlling reflection phase of EM wave impinging on metasurface. The bandwidth of more than 80% is achieved at the normal incidence for the −9.5 dB RCS reduction for both monostatic and bistatic. The experiment results are in good agreement with the corresponding simulations. Experimental and simulated results validate the effectiveness of our approach.

Periodic Arrangement. Because the scattered fields are redirected toward four and six directions for the rectangular and hexagonal [5–8] checkerboard surfaces, respectively, the maxima of the bistatic RCS are −8.1 and −8.9 dB compared to those of the equal-sized PEC ground planes.

Optimized Phase Layout. The intercepted EM energies were redirected to more directions, so that the scattered fields are controlled in low levels in all directions. A lot of bistatic RCS lobes (≫4 or 6) are produced in space, leading to a great bistatic RCS reduction. Furthermore, the number of RCS lobes increases as metasurface size (lattice number) increases; thus the reduced magnitude of monostatic and bistatic RCSs becomes greater. For the metasurface designed based on optimized phase layout in this paper, more than 80 scattering beams are created in space, and the bistatic RCS reduction reaches 17.6 dB.
Competing Interests

The authors declare that there are no competing interests regarding the publication of this paper.

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