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Angular and energy dependence of cross sections for ejection of electrons from water vapor. II. 15–150-keV proton impact

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Absolute values of cross sections for electron ejection by protons of 15–150 keV energy have been measured as a function of the angle and energy of the electrons. The range of angles was 10° to 160° and the electron energy range was 1–300 eV. The doubly differential cross sections were also integrated over energy or angle to obtain singly differential and total cross sections and also average ejected electron energies. Good agreement is obtained with Senger's DDCS-MT (doubly differential cross section—mixed treatment) theoretical treatment using the Salin factor.

I. INTRODUCTION

Measurements of the angular and energy dependence of the double differential cross sections (DDCS) for ejection of electrons in ion-atom collisions have uncovered new mechanisms of ionization¹ and the availability of such measurements has stimulated theoretical work on ionization.² Unfortunately, while there are partially successful theoretical treatments at high energies (above say, 100 keV), the low-energy region remains without an accurate method of calculating electron-ejection cross sections. Thus, there is still a need for measurements in the energy range below the maximum in the total cross-section curve. In searching for systematics in the process, it is desirable to have measurements from a wide variety of targets with various degrees of complexity, different ionization potentials, different outer shell angular momenta, etc.

While low-energy measurements of this kind are available for a number of simple gases,^{3,4} there have been none for molecules more complex than oxygen. Besides its theoretical interest as a triatomic molecule, water vapor is of great practical interest to those modeling radiation damage in living tissue. Toburen and Wilson⁵ have made DDCS measurements on water vapor from 300–1500 keV. The present work extends this range downward into the region where present theory has little to say.

II. EXPERIMENTAL PROCEDURE

The apparatus will be only briefly described here since a fuller description is already available elsewhere.^{1,6} Protons accelerated from an rf ion source were magnetically analyzed and finely collimated before entering the target gas. A biased Faraday cup caught the proton beam and the collected charge was integrated. Electrons ejected from the static gas (typically at 0.5 mTorr) entered a 127° electrostatic analyzer which could be placed at any of the eight angles from 10° to 160° relative to the beam direction. The angular acceptance was 2.16° and the energy resolution 4.4%. No preacceleration of electrons was used. A 17-stage electron multiplier with its first stage held at 80 V detected individual electrons. The earth's

magnetic field was annulled to within a few mG by the use of three mutually perpendicular pairs of Helmholtz coils.

The efficiency of the detector was determined in an auxiliary measurement using the apparatus of Cacak and Jorgensen⁷ between an electron gun and the electron multiplier detector. Two different apertures, the sizes of which had been carefully measured, could be moved in front of the first dynode. A defocused beam from an electron gun was then directed to the apertures. The uniformity of the beam was checked by moving the smaller aperture across the larger one. The efficiency of the multiplier was determined by comparing the current through the larger aperture, using the first dynode as a Faraday cup, and the count rate through the smaller one when the multiplier was operated in its normal way. By this method the efficiency was found to be 0.63 ± 0.05 .

Corrections were made for neutralization of the proton beam as it traversed the gas and absorption of the ejected electrons by the target gas, as well as for electrons originating from the background gas. Target pressure was measured with a capacitance manometer. As in earlier work¹ cross sections differential in angle $\sigma(\theta)$ or energy $\sigma(W)$ were obtained by numerical integration of the DDCS $\sigma(W, \theta)$. Total electron-ejection cross sections σ_{-} were calculated by integrating over both energy and angle.

There was a 7% uncertainty in the measurement of target gas density and an 8% uncertainty in the detector efficiency. The uncertainties in other quantities were small so the combined uncertainty in the DDCS was 11% except at the upper end of the electron energy range where low count rates resulted in larger statistical errors. Measurements of the cross sections below about 10 eV may also have additional systematic errors due to deflection of electrons by stray electric and magnetic fields and by the possible presence of spurious low-energy secondary electrons from surfaces. Since the former effect tends to reduce the measured cross sections and the latter to increase them, the results may be either too low or too high. The magnitude of the error is difficult to determine but is probably small above 10 eV. However, since low-energy cross sections contribute strongly to $\sigma(\theta)$ and σ_{-} , there is

an additional uncertainty in the integrated values which we estimate at 15% yielding a total uncertainty in $\sigma(\theta)$ and σ_{-} of about 19%.

III. EXPERIMENTAL RESULTS

DDCS are shown in Figs. 1–3. In Fig. 1, showing the electron-energy distributions at various angles, the 10° curve shows the beginning of the broad binary-encounter peak which is expected to be centered at 327 eV. A more prominent peak at about 80 eV can be attributed to electron transfer to the continuum.^{8,9} Interestingly, a small shoulder at the same energy also shows up in the 160° curve. This peak is similar to the electron-loss peak seen when using projectiles carrying electrons. It is possible that some protons in the beam in this experiment were neutralized in the target gas before reaching the collision center, thus leading to the possibility that this mechanism could explain the 160° peak. However, when we attempted to verify this by reducing the target-gas pressure, we found no change in the shape of the curve. Since the pressure in the 2-m-long beam line was only $1\text{--}2 \times 10^{-6}$ Torr, the neutralization of the beam there should not be greater than 1–2% at low energies and would be far too small at 150 keV to produce a noticeable effect. Another possible explanation is that it is due to a second-order electron transfer to the continuum.

A comparison of the angular distributions of electrons given in Figs. 2 and 3 shows that at 100 keV the electrons are more strongly peaked in the forward direction than at lower impact energies. Also one notes that at the lower impact energy there is a rise in the cross section in the backward direction while there is little if any at 100 keV.

Senger^{10,11} has made calculations of the singly and doubly differential cross sections for electron ejection from

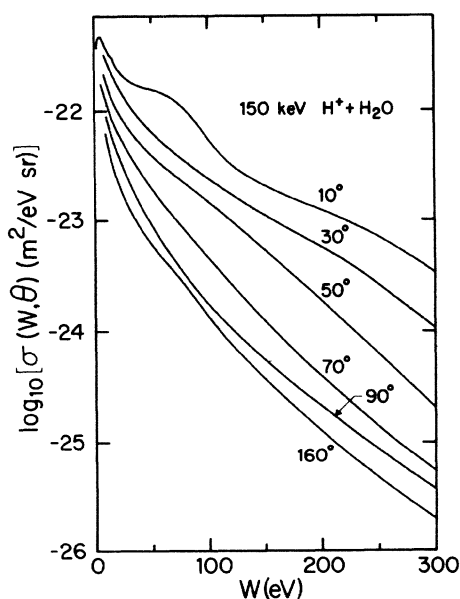


FIG. 1. Doubly differential cross sections for ejection of electrons by 150-keV protons from water vapor as a function of ejection energy for various angles of ejection.

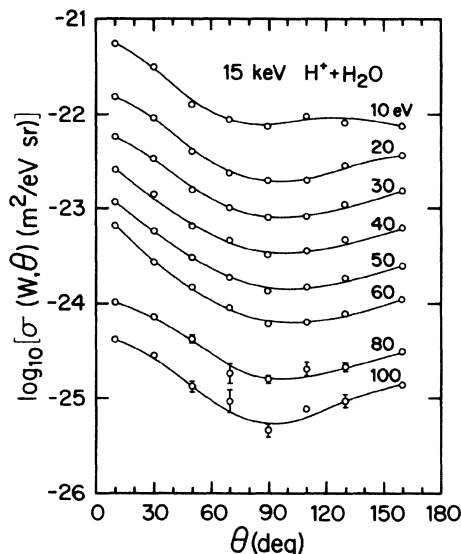


FIG. 2. Doubly differential cross sections for ejection of electrons by 15-keV protons from water vapor as a function of angle of ejection for various ejection energies. Error bars on some data points show the error due to statistical fluctuations in the count.

proton impact on water molecules. He applied the plane-wave Born approximation to the various molecular orbitals, modifying the form factors and making corrections for binding energies of inner shells. In addition, he has applied the Salin factor¹² to account for the mechanism of electron transfer to the continuum. His results are com-

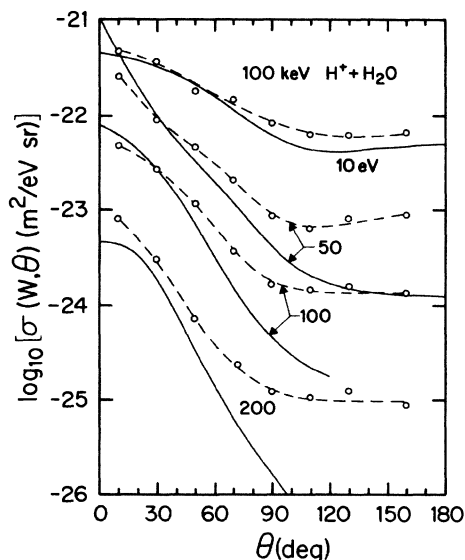


FIG. 3. Doubly differential cross sections for ejection of electrons by 100-keV protons from water vapor as a function of angle. Points are experimental values with the dashed lines drawn to guide the eye. Solid lines are theoretical calculations by Senger *et al.* (Refs. 10 and 11).

TABLE I. Values of $\sigma(W)$ in m^2/eV , σ_- in m^2 , and W_{av} in eV for $\text{H}^+ + \text{H}_2\text{O}$ collisions.

W (eV)	Projectile energy (keV)						
	15	20	30	50	70	100	150
2	1.5(-21) ^a	2.1(-21)	2.2(-21)	2.2(-21)	2.6(-21)	1.6(-21)	1.6(-21)
3	2.1(-21)	2.4(-21)	2.5(-21)	2.3(-21)	2.5(-21)	2.2(-21)	1.8(-21)
5	2.5(-21)	2.8(-21)	2.7(-21)	2.5(-21)	2.5(-21)	2.3(-21)	2.2(-21)
8	2.0(-21)	2.2(-21)	2.4(-21)	2.2(-21)	2.1(-21)	1.9(-21)	1.7(-21)
13	1.0(-21)	1.2(-21)	1.7(-21)	1.6(-21)	1.5(-21)	1.4(-21)	1.2(-21)
20	4.5(-22)	6.2(-22)	8.8(-22)	1.2(-21)	1.1(-21)	9.4(-22)	7.7(-22)
30	1.8(-22)	2.8(-22)	4.4(-22)	6.7(-22)	7.4(-22)	6.4(-22)	5.0(-22)
50	3.2(-23)	6.2(-23)	1.3(-22)	2.5(-22)	3.2(-22)	3.6(-22)	2.8(-22)
80	3.8(-24)	8.2(-24)	2.5(-23)	7.4(-23)	1.1(-22)	1.4(-22)	1.4(-22)
130	6.7(-25)	8.1(-25)	2.0(-24)	9.5(-24)	2.3(-23)	4.3(-23)	4.8(-23)
200	1.9(-25)	2.0(-25)	3.1(-25)	7.9(-25)	2.6(-24)	8.7(-24)	1.6(-23)
300	5.6(-26)	7.2(-26)	9.3(-26)	1.1(-25)	2.2(-25)	7.1(-25)	3.0(-24)
σ_-	3.2(-20)	4.0(-20)	5.1(-20)	6.0(-20)	6.4(-20)	6.2(-20)	5.5(-20)
W_{av}	12.33	13.6	16.7	22.4	26.0	32.2	37.0

^aThe designation 1.4(-22) means 1.4×10^{-22} .

pared with the present experimental data in Fig. 3. While the Salin factor improves the agreement at the forward angles, it worsens the agreement at angles greater than 90° .

SDCS are shown in Figs. 4–6 and in Table I. While there are no earlier experiments in our energy range with which we can compare directly, plots may be made of the SDCS versus proton energy to compare with the corresponding data of Toburen and Wilson⁵ as shown in Fig. 4. While the trend of their data at the lowest energies indicates an extrapolation to somewhat lower values than our data, the general agreement is satisfactory.

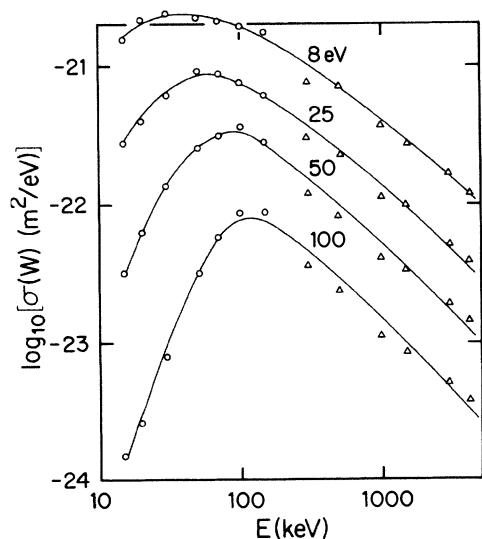


FIG. 4. Singly differential cross sections for ejection of electrons by protons in water vapor as a function of proton energy. Circles, present data; triangles, data of Toburen and Wilson (Ref. 5).

Senger's calculations of SDCS are shown in Fig. 5 where they are compared to the experimental values. Also plotted are calculations made using the binary encounter model¹³ integrated over a Fock distribution of orbital energies.¹⁴ Partial cross sections for each molecular subshell were added to obtain the cross sections shown. Senger's calculations are in excellent agreement with experiment in this case while the binary-encounter approximation (BEA) overestimates the cross sections in the low-energy region.

Figure 6 shows a comparison of the SDCS versus electron energy for various impact energies. While in every case we took the precaution of taking a background count

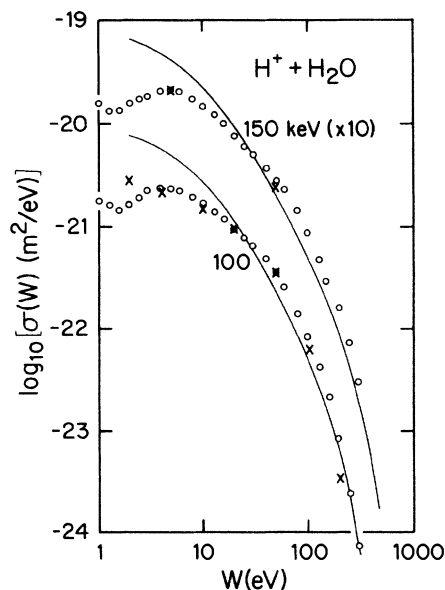


FIG. 5. Singly differential cross sections for ejection of electrons by 100- and 150-keV protons in water vapor as a function of electron energy. Circles, present data; \times 's, calculations by Senger *et al.* (Refs. 10 and 11); solid lines, BEA calculations.

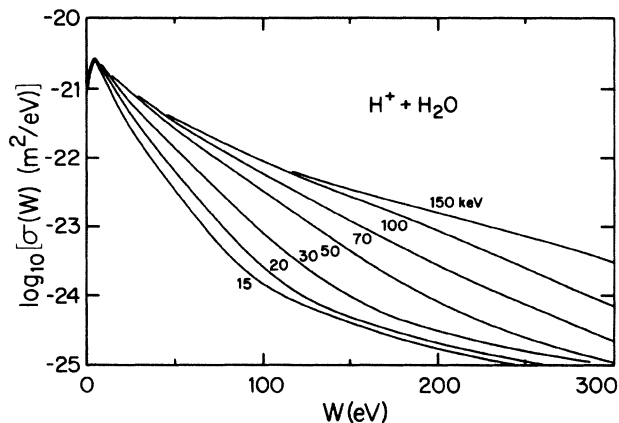


FIG. 6. Singly differential cross sections vs ejected electron energy for various proton energies.

without the target gas, there may have been a target pressure-dependent background which could not be eliminated in this way. This could, e.g., be due to uv photons generated in the collision region¹⁵ which managed to be reflected into the electron multiplier detector. This may account for the leveling off of the cross sections for low impact energy and high ejection energy. In earlier work⁴ using a different apparatus and different target gases, the low-energy cross sections decreased approximately exponentially with electron energy.

Total electron-ejection cross sections are shown in Fig. 7 compared to cross sections measured more directly which were reported earlier.¹⁶ While the present results are generally somewhat higher, the agreement is within the combined uncertainties of the two measurements.

IV. SCALING CROSS SECTIONS

In an attempt to find relationships between electron-ejection cross sections for protons on different targets, we have tried two methods of scaling. The first is based on the Born approximation and was given for SDCS by Rudd, Sautter, and Bailey.¹⁷ It may be written

$$\sigma(W, E_p, B, N) = N(R/B)^3 \sigma(WR/B, E_p R/B, R, 1), \quad (1)$$

where N is the number of target electrons with binding energy B , R is the Rydberg of energy, W is the electron-ejection energy, and E_p the proton-beam energy. This allows scaling from cross sections for atomic hydrogen to

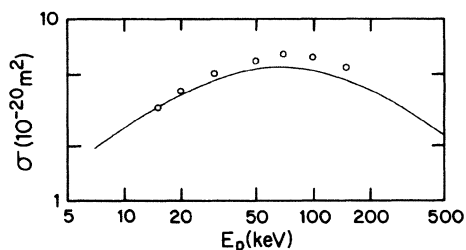


FIG. 7. Total electron-ejection cross sections vs proton energy. Circles, present data; solid line, data of Rudd *et al.* (Ref. 16).

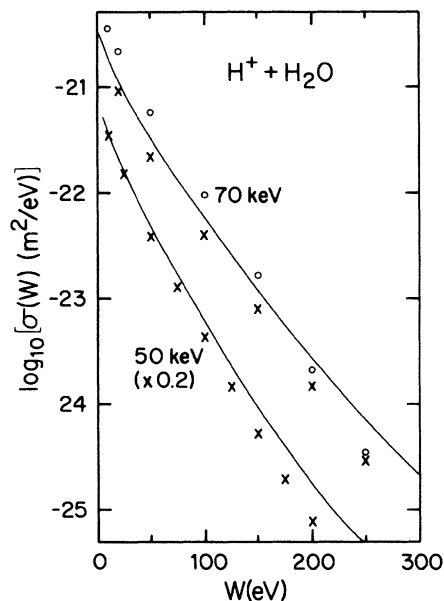


FIG. 8. Singly differential cross sections for 50- and 70-keV protons on water vapor. Solid lines, present experimental data; circles, results of scaling data on H_2 (Ref. 4) by method described in text; \times 's, results of Bragg scaling using data on H_2 (Ref. 4) and O_2 (Ref. 19).

any other target provided the values of N and B for the various subshells of the target are known.

We have taken the data of Rudd⁴ for H_2 and using 15.42 eV as the ionization potential have scaled that to atomic hydrogen using Eq. (1) and then using the values of N and B for water vapor in the preceding paper¹⁸ scaled the hydrogen data to the various shells of water vapor. Approximations had to be made because data were not available at exactly the proton energies needed in the scaling; however, the error introduced should have only a small effect on the results. The results of this scaling are shown in Fig. 8.

Another method of scaling that has been applied to stopping powers is the Bragg rule for additivity. Applied to water it yields

$$\sigma(H_2O) = \sigma(H_2) + \frac{1}{2}\sigma(O_2). \quad (2)$$

Calculations were made using the data of Rudd⁴ for H_2 and the data of Crooks and Rudd¹⁹ for O_2 . These results are also shown in Fig. 8.

For this case, Bragg scaling yields reasonably good values at low ejection energies but gives values which are too low at higher energies. The other method is slightly better at high energies, but much worse for small values of W .

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