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Prediction of a spin-polarized two-dimensional electron gas at the LaAlO₃/EuO(001) interface

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First-principles calculations predict the existence of a spin-polarized two-dimensional electron gas (2DEG) at the LaO/EuO interface in a LaAlO₃/EuO(001) heterostructure. This polar interface favors electron doping into the Eu-5*d* conduction bands resulting in a 2DEG formed at the interface. Due to the exchange splitting of the Eu-5*d* states, the 2DEG is spin polarized below the Curie temperature of EuO. The predicted mechanism for the formation of a spin-polarized 2DEG at the interface between polar and ferromagnetic insulators may provide a robust magnetism of the 2DEG which is interesting for spintronics applications.

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The field of research related to complex oxide materials has been developing vigorously in the last few decades. Oxides exhibit an abundance of macroscopic physical properties, often involving the interplay between magnetism, ferroelectricity, and conductivity.¹ Even a more rich spectrum of physical phenomena occurs if two or more oxides are combined with atomic-scale precision in a heterostructure to form a new nanoscale material.^{2,3} Recent advances in thin-film deposition and characterization techniques made possible the experimental realization of such oxide heterostructures, promising advanced functionalities and device concepts. A prominent example is the recent discovery of the formation of a metallic phase at the interface between two oxide insulators, SrTiO₃ and LaAlO₃.⁴ It was found that this metallic phase is confined within a couple of nanometers near the interface⁵ and therefore can be regarded as a two-dimensional electron gas (2DEG). The 2DEG has a very high carrier density and a relatively high carrier mobility making it attractive for applications in nanoelectronics, e.g., as all-oxide field-effect transistors.^{6,7} These properties of the 2DEG at oxide interfaces have stimulated significant research activity both in experiment⁸⁻¹³ and in theory.¹⁴⁻²¹

Recently, it was found that at ultralow temperatures the 2DEG occurring at the interface between the nonmagnetic LaAlO₃ and SrTiO₃ materials may become magnetic.²² This behavior was attributed to the exchange splitting of the induced electrons in the Ti-3*d* conduction band, which is corroborated by spin-polarized first-principles calculations of LaAlO₃/SrTiO₃,^{18,19} as well as LaTiO₃/SrTiO₃ interfaces.¹⁶ Making 2DEG gas spin polarized is a very exciting prospect for spintronics applications, where the involvement of the spin degree of freedom broadens the spectrum of potential applications.²³ Very recently it was proposed that it may be possible to create a fully spin-polarized 2D electron gas by replacing one monolayer of SrO by LaO in SrMnO₃.²⁴

Here, we pursue a different route to achieve a spin-polarized two-dimensional electron gas by employing a ferromagnetic insulator as one of the constituents in the oxide heterostructure. Spin-polarized properties of the 2DEG are, in this case, expected to be inherited from the ferromagnetism of the oxide and, consequently, this approach may lead to a more robust magnetism in the 2DEG, which is

beneficial for applications. To illustrate the idea we consider EuO as a representative ferromagnetic insulator in conjunction with LaAlO₃ to form a spin-polarized 2DEG at the LaAlO₃/EuO(001) interface.

EuO has a rocksalt crystal structure and is a ferromagnetic insulator (semiconductor) with a bulk Curie temperature (T_C) of 69 K. A divalent Eu ion in EuO has a half-filled 4*f* shell leading to the ⁸S_{7/2} ground state and the magnetic moment of 7μ_B per Eu ion. The Heisenberg exchange coupling between the localized 4*f* electrons causes the ferromagnetic ordering in EuO below T_C . The half-filled 4*f* band is separated from the 5*d*-6*s* conduction bands by a band gap of 1.12 eV at room temperature.²⁵ In the ferromagnetic state of EuO, the direct exchange interaction between the localized 4*f* moments and the delocalized 5*d* conduction-band states leads to the spin splitting of the latter. The spin splitting of the 5*d* states as large as 0.6 eV produces the full spin polarization near the bottom of the conduction band.²⁶ This feature of the EuO conduction band allows creating a highly spin-polarized electron gas by appropriate *n*-doping of the material. In particular, recently it was found that the transport spin polarization of conduction electrons in EuO doped with La exceeds 90%.²⁷

Here we demonstrate that the electron doping may be achieved locally at the LaAlO₃/EuO(001) interface to create a spin-polarized 2DEG. The mechanism responsible for the 2DEG formation is similar to that known for the LaAlO₃/SrTiO₃ interface.²⁸ LaAlO₃ consists of alternating (LaO)⁺ and (AlO₂)⁻ charged planes whereas EuO consists of (EuO)⁰ neutral planes. When LaAlO₃ is deposited on top of EuO, the divergence in the electrostatic potential can be avoided by transferring half an electron per two-dimensional unit cell to the LaO/EuO terminated interface. A charge transfer to the interface also occurs if the LaAlO₃ layer is nonstoichiometric and terminated with the LaO monolayers on both sides. In this case an “extra” electron is introduced into the system due to the uncompensated ionic charge on the additional (LaO)⁺ monolayer. This electronic charge is accommodated by partially occupying conduction-band states near the interface producing a 2DEG. In ferromagnetic EuO, the conduction band is formed by exchange-split Eu-5*d* states and so the 2DEG is expected to be spin polarized.

Thus, by forming the $\text{LaAlO}_3/\text{EuO}(001)$ interface, one can achieve a spin-polarized 2DEG below the Curie temperature of EuO.

To quantitatively demonstrate our prediction, we perform first-principles calculations of the electronic structure of the $\text{LaAlO}_3/\text{EuO}(001)$ interface within the framework of density-functional theory²⁹ implemented within Vienna *ab initio* simulation package (VASP).³⁰ The spin-polarized calculations include the electron-ion potential which is described within the projected augmented method.³¹ The exchange-correlation effects are treated within the Perdew-Burke-Ernzerhoff form³² of the generalized gradient approximation (GGA). Self-consistent calculations are performed using a plane-wave basis set limited by a cutoff energy of 520 eV and the $6 \times 6 \times 1$ Monkhorst-Pack k -point mesh³³ with energy converged to 10^{-5} eV/cell. Atomic relaxations are performed until the Hellmann-Feynman forces on atoms have become less than 30 meV/Å.

On-site correlations for the Eu-4*f* orbitals are included within the GGA+*U* approach.³⁴ The value of $J=0.6$ eV is calculated using the constrained occupation method³⁵ by considering the 4*f* states as an open-core shell and finding the GGA energy difference between the $4f_{\uparrow}^7 4f_{\downarrow}^0$ and $4f_{\uparrow}^6 4f_{\downarrow}^1$ configurations.³⁶ We find, however, that the value of $U=5.3$ eV obtained by this method appears to be too small and leads to the 4*f* states being too shallow with respect to the conduction band of EuO. This discrepancy is due to the underestimation by GGA of the intrinsic insulating gap between the O-2*p* and Eu-5*d* states. Therefore, we adjusted the value of U empirically and found that $U=7.5$ eV results in a very reasonable agreement with experiment. In particular, $U+6J=11.1$ eV agrees well with the occupied (unoccupied) 4*f*-state splitting observed in photoemission (inverse photoemission) for Eu metal; the optical band gap at the *X* point of 0.94 eV is consistent with the zero-temperature experimental value of approximately 0.95 eV (Ref. 37); and the lattice constant $a=5.188$ Å agrees with the experimental value of $a=5.144$ Å. The exchange splitting of the Eu 5*d* orbitals is found to be $\Delta_d=0.75$ eV. Since the La-4*f* bands lie at higher energy than that predicted by GGA, we impose $U=11$ eV on these orbitals to avoid their spurious mixing with the conduction bands of LaAlO_3 .¹⁶

We consider the LaO/EuO terminated interface of a $(\text{LaAlO}_3)_{8.5}/(\text{EuO})_{15}$ superlattice (i.e., a supercell containing 8.5 unit cells of LaAlO_3 and 15 monolayers of EuO) stacked in the [001] direction, as shown in Fig. 1. We use periodic boundary conditions and impose mirror plane symmetry at the central EuO monolayer. The in-plane lattice constant of the superlattice is fixed to the calculated lattice constant of cubic LaAlO_3 , $a=3.81$ Å, which is in good agreement with the experimental value $a=3.79$ Å. Under this constraint, EuO exhibits tetragonal distortion of $c/a=0.947$. This distortion does not change significantly the electronic structure of bulk EuO. In particular, the band gap of EuO becomes $E_g=1.15$ eV and the exchange splitting of the *d* orbitals becomes $\Delta_d=0.66$ eV. The out-of-plane lattice constant of the supercell is determined by optimizing the interface separation distance between the LaAlO_3 and EuO subunits keeping their bulk lattice constants fixed. Then, we fix the supercell dimensions and relax the atomic positions of all the atoms in the $(\text{LaAlO}_3)_{8.5}/(\text{EuO})_{15}$ superlattice.

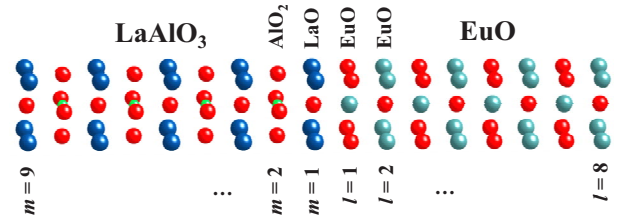


FIG. 1. (Color online) Atomic structure of the LaO/EuO interface in the $(\text{LaAlO}_3)_{8.5}/(\text{EuO})_{15}(001)$ superlattice containing 8.5 unit cells of LaAlO_3 and 15 monolayers of EuO within the supercell. Indices l and m denote atomic monolayers and are increasing with separation from the interface. $l=8$ in EuO and $m=9$ in LaAlO_3 correspond to the middle of the respective layers.

Figure 2 shows the calculated atomic displacements within the $(\text{LaAlO}_3)_{8.5}/(\text{EuO})_{15}$ supercell. It is seen that the largest structural relaxations occur in the vicinity of the LaO/EuO interfaces and involve a polar distortion in which the negatively charged O anions are displaced with respect to the positively charged cations (either La, Al, or Eu). When moving away from the interface, the magnitude of the displacements decays and the La-La and Eu-Eu distances revert to constant values close to those in the bulk.

Figure 3 shows the layer-resolved majority- and minority-spin densities of states (DOS) on the EuO [Fig. 3(a)] and LaO [Fig. 3(b)] monolayers located at different planes l and m away from the LaO/EuO interface (see labeling in Fig. 1). It is seen that there are occupied states below the Fermi energy on both the LaO and EuO monolayers near the interface which indicate the formation of the *n*-type 2DEG at the LaO/EuO interface. The DOS on AlO_2 monolayers [Fig. 3(c)] is negligible near the Fermi energy compared to that on the LaO and EuO monolayers. Far away from the interface, the DOS at the Fermi energy on both the LaO and EuO monolayers drops to zero, reflecting the insulating nature of bulk LaAlO_3 and EuO.

The central result of our calculation is the formation of spin-polarized 2DEG at the $\text{LaAlO}_3/\text{EuO}(001)$ interface.

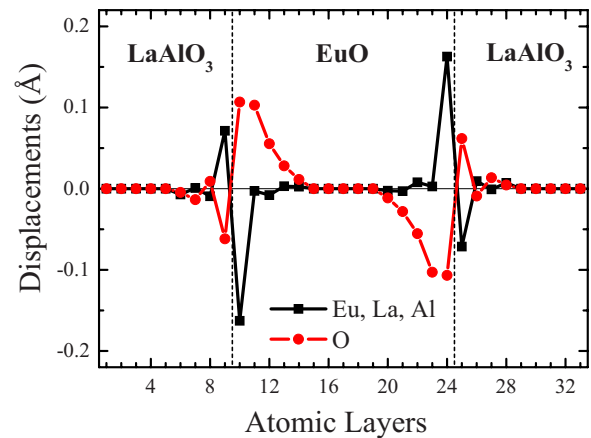


FIG. 2. (Color online) Atomic displacements in the $(\text{LaAlO}_3)_{8.5}/(\text{EuO})_{15}(001)$ superlattice with respect to the atomic “bulk” positions. The latter are determined by fixing the in-plane lattice constant and optimizing the interlayer distance, as described in the text. The vertical dashed lines indicate interfaces.

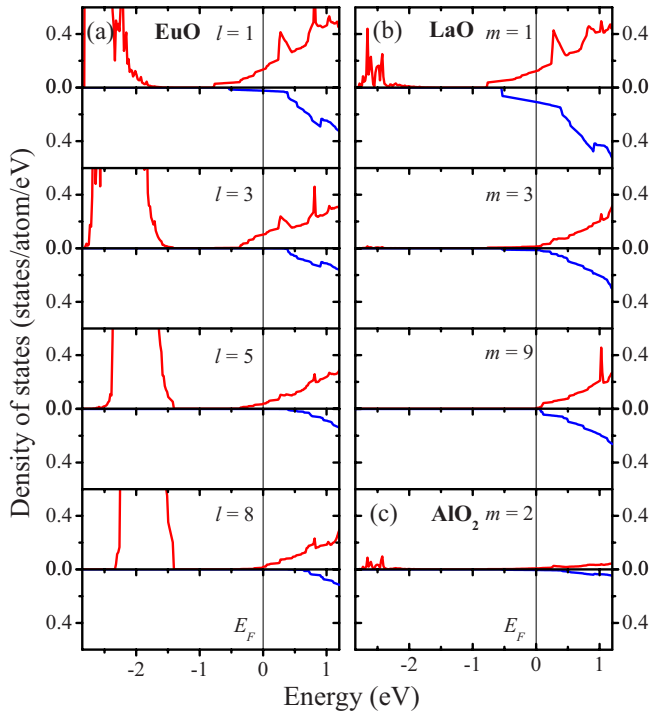


FIG. 3. (Color online) Layer- and spin-resolved DOS on (a) EuO, (b) LaO, and (c) AlO_2 monolayers located at different planes l and m away from the LaO/EuO interface (as labeled in Fig. 1). Top (bottom) panels show the majority (minority) spin. In panel (a), the majority-spin states at energies below -1 eV are the occupied Eu-4*f* states. The vertical lines denote the Fermi energy (E_F).

This fact is evident from Fig. 3(a) indicating a significant difference in the occupation of the EuO majority- and minority-spin conduction bands near the interface. This is the direct consequence of the exchange splitting of the conduction 5*d* states in EuO and the occupation of these states due to the electron doping of the interface. As seen from Fig. 3(a), starting from the third EuO monolayer ($l=3$) away from the LaO/EuO interface only majority-spin states are occupied in the conduction band of EuO indicating a tendency to half metallicity known for the *n*-doped bulk EuO.²⁵ Figure 3(b) reveals the spin splitting of the conduction bands of LaAlO_3 which is the result of the exchange interaction between the Eu-4*f* and La-5*d* states across the interface.

For the $\text{LaAlO}_3/\text{EuO}$ interface, we find that the top of the O-2*p* valence bands lies at about -3.6 eV in LaAlO_3 and at about -3.0 eV in EuO with respect to the Fermi energy. Due to a much larger energy gap between the O-*p* valence bands and the conduction bands in LaAlO_3 than in EuO (the calculated values are 3.7 and 2.7 eV, respectively, and the experimental values are 5.6 and 4.4 eV,³⁸ respectively), the conduction-band minimum lies lower in EuO than in LaAlO_3 resulting in the charge accumulating mainly within the EuO layer. This is evident from Fig. 4 which shows the distribution of the spin-dependent charge across the unit cell. These charges are calculated by integrating the spin- and layer-resolved DOS from the conduction-band minimum up to the Fermi energy.³⁹ The large spin polarization of the 2DEG is

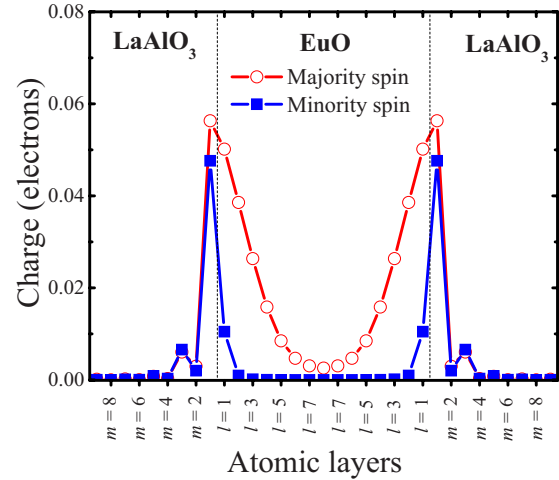


FIG. 4. (Color online) Spin-dependent charge distribution across the $(\text{LaAlO}_3)_{8.5}/(\text{EuO})_{15}(001)$ supercell. The notation for the atomic layers is the same as in Fig. 1. The dashed lines indicate interfaces.

seen at the EuO monolayers near the $\text{LaAlO}_3/\text{EuO}$ interface. The estimated value of the spin polarization of the free charge density is about 50%.

We note that our calculation is performed for the nonstoichiometric LaAlO_3 layer which is assumed to be LaO terminated on both sides. Similar to the previous theoretical studies,^{14–19} in this geometry an “extra” electron is introduced in the system due to the uncompensated ionic charge on the additional $(\text{LaO})^+$ monolayer. For a stoichiometric LaAlO_3 layer deposited on top of EuO, we expect the formation of 2DEG due to electronic reconstruction resulting from the charge transfer to the interface that eliminates the increasing electrostatic potential in LaAlO_3 .²⁹ Due to dissimilar band alignments at the $\text{LaAlO}_3/\text{EuO}$ interface compared to the $\text{LaAlO}_3/\text{SrTiO}_3$ interface, we anticipate a different critical thickness of LaAlO_3 required to form a 2DEG.⁶

In summary, based on first-principles calculations, we have predicted the possibility to create a spin-polarized 2DEG at the LaO/EuO interface in the $\text{LaAlO}_3/\text{EuO}(001)$ heterostructure. We demonstrated that this polar interface favors electron doping into the Eu-5*d* conduction bands, rendering a 2DEG formed at the interface. Due to the exchange splitting of the Eu-5*d* states, the 2DEG becomes spin polarized below the Curie temperature of EuO. The predicted mechanism for the formation of a spin-polarized 2DEG at the interface between polar and ferromagnetic insulators may lead to a robust magnetism of a 2DEG which is interesting for spintronics applications.

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- ¹E. Dagotto, *Science* **309**, 257 (2005).
- ²Y. Tokura and H. Y. Hwang, *Nature Mater.* **7**, 694 (2008).
- ³J. W. Reiner, F. J. Walker, and C. H. Ahn, *Science* **323**, 1018 (2009).
- ⁴A. Ohtomo and H. Y. Hwang, *Nature (London)* **427**, 423 (2004).
- ⁵M. Basletic, J.-L. Maurice, C. Carrétéro, G. Herranz, O. Copie, M. Bibes, E. Jacquet, K. Bouzehouane, S. Fusil, and A. Barthélémy, *Nature Mater.* **7**, 621 (2008).
- ⁶S. Thiel, G. Hammerl, A. Schmehl, C. W. Schneider, and J. Mannhart, *Science* **313**, 1942 (2006).
- ⁷C. Cen, S. Thiel, J. Mannhart, and J. Levy, *Science* **323**, 1026 (2009).
- ⁸N. Reyren *et al.*, *Science* **317**, 1196 (2007).
- ⁹Y. Hotta, T. Susaki, and H. Y. Hwang, *Phys. Rev. Lett.* **99**, 236805 (2007).
- ¹⁰A. Kalabukhov, R. Gunnarsson, J. Börjesson, E. Olsson, T. Claesson, and D. Winkler, *Phys. Rev. B* **75**, 121404(R) (2007).
- ¹¹W. Siemons, G. Koster, H. Yamamoto, W. A. Harrison, G. Lucovsky, T. H. Geballe, D. H. A. Blank, and M. R. Beasley, *Phys. Rev. Lett.* **98**, 196802 (2007).
- ¹²G. Herranz, M. Basletic, M. Bibes, C. Carrétéro, E. Tafrá, E. Jacquet, K. Bouzehouane, C. Deranolt, A. Hamzić, J.-M. Broto, A. Barthélémy, and A. Fert, *Phys. Rev. Lett.* **98**, 216803 (2007).
- ¹³K. Yoshimatsu, R. Yasuhara, H. Kumigashira, and M. Oshima, *Phys. Rev. Lett.* **101**, 026802 (2008).
- ¹⁴Z. S. Popovic and S. Satpathy, *Phys. Rev. Lett.* **94**, 176805 (2005); Z. S. Popovic, S. Satpathy, and R. M. Martin, *ibid.* **101**, 256801 (2008).
- ¹⁵M. S. Park, S. H. Rhim, and A. J. Freeman, *Phys. Rev. B* **74**, 205416 (2006).
- ¹⁶S. Okamoto, A. J. Millis, and N. A. Spaldin, *Phys. Rev. Lett.* **97**, 056802 (2006).
- ¹⁷J. M. Albina, M. Mrovec, B. Meyer, and C. Elsasser, *Phys. Rev. B* **76**, 165103 (2007).
- ¹⁸R. Pentcheva and W. E. Pickett, *Phys. Rev. B* **74**, 035112 (2006); *Phys. Rev. Lett.* **99**, 016802 (2007); **102**, 107602 (2009).
- ¹⁹K. Janicka, J. P. Velev, and E. Y. Tsybal, *J. Appl. Phys.* **103**, 07B508 (2008); *Phys. Rev. Lett.* **102**, 106803 (2009).
- ²⁰J. Lee and A. A. Demkov, *Phys. Rev. B* **78**, 193104 (2008).
- ²¹S. Ishibashi and K. Terakura, *J. Phys. Soc. Jpn.* **77**, 104706 (2008).
- ²²A. Brinkman, M. Huijben, M. van Zalk, J. Huijben, U. Zeitler, J. C. Maan, W. G. van der Wiel, G. Rijnders, D. H. A. Blank, and H. Hilgenkamp, *Nature Mater.* **6**, 493 (2007).
- ²³I. Žutić, J. Fabian, and S. D. Sarma, *Rev. Mod. Phys.* **76**, 323 (2004).
- ²⁴B. R. K. Nanda and S. Satpathy, *Phys. Rev. Lett.* **101**, 127201 (2008).
- ²⁵A. Mauger and C. Godart, *Phys. Rep.* **141**, 51 (1986).
- ²⁶P. G. Steeneken, L. H. Tjeng, I. Elfimov, G. A. Sawatzky, G. Ghiringhelli, N. B. Brookes, and D.-J. Huang, *Phys. Rev. Lett.* **88**, 047201 (2002).
- ²⁷A. Schmehl *et al.*, *Nature Mater.* **6**, 882 (2007).
- ²⁸N. Nakagawa, H. Y. Hwang, and D. A. Muller, *Nature Mater.* **5**, 204 (2006).
- ²⁹P. Hohenberg and W. Kohn, *Phys. Rev.* **136**, B864 (1964); W. Kohn and L. J. Sham, *ibid.* **140**, A1133 (1965).
- ³⁰G. Kresse and J. Furthmüller, *Phys. Rev. B* **54**, 11169 (1996).
- ³¹P. E. Blöchl, *Phys. Rev. B* **50**, 17953 (1994).
- ³²J. P. Perdew, K. Burke, and M. Ernzerhof, *Phys. Rev. Lett.* **78**, 1396 (1997).
- ³³H. J. Monkhorst and J. D. Pack, *Phys. Rev. B* **13**, 5188 (1976).
- ³⁴A. I. Liechtenstein, V. I. Anisimov, and J. Zaanen, *Phys. Rev. B* **52**, R5467 (1995).
- ³⁵P. H. Dederichs, S. Blügel, R. Zeller, and H. Akai, *Phys. Rev. Lett.* **53**, 2512 (1984).
- ³⁶I. V. Solov'yev, *J. Phys.: Condens. Matter* **20**, 293201 (2008).
- ³⁷G. Busch and P. Wachter, *Phys. Kondens. Mater.* **5**, 232 (1966); *Z. Angew. Phys.* **26**, 2 (1968).
- ³⁸D. E. Eastman, F. Holtzberg, and S. Methfessel, *Phys. Rev. Lett.* **23**, 226 (1969).
- ³⁹The local charges are evaluated within the Wigner-Seitz spheres, as specified by VASP.