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Abstract

Using a newly developed first-principles approach to simulations of ionic molecular crystals, we performed static relaxation, molecular-dynamics simulation, and lattice-dynamics calculations, and measurements of the Raman spectrum, for the *Pnam* structure of K_2SO_4 . It was found that the structure does not have the zone-center instability present in isomorphous K_2SeO_4 found in an earlier study. This difference between the two systems is attributed to the different charge distributions in the molecular ions. The calculated Raman-active zone-center frequencies for the *Pnam* structure of K_2SO_4 were found to be in general agreement with the experimental Raman frequencies.

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Using a newly developed first-principles approach to simulations of ionic molecular crystals, we performed static relaxation, molecular-dynamics simulation, and lattice-dynamics calculations, and measurements of the Raman spectrum, for the *Pnam* structure of K_2SO_4 . It was found that the structure does not have the zone-center instability present in isomorphous K_2SeO_4 found in an earlier study. This difference between the two systems is attributed to the different charge distributions in the molecular ions. The calculated Raman-active zone-center frequencies for the *Pnam* structure of K_2SO_4 were found to be in general agreement with the experimental Raman frequencies.

First-principles simulations using *ab initio* Gordon-Kim¹ pair potentials have proven very successful in the studies of phase transitions, lattice dynamics, and other important properties of complex ionic solids.²⁻⁸ Based on the rigid-free-ion electron charge densities and the electron-gas model, these pair potentials give accurate descriptions of the potential-energy surfaces for ionic crystals. Recently this approach has been extended to treat molecular ionic solids, i.e., solids that contain molecular ions, e.g., SO_4^{2-} , SeO_4^{2-} , NO_2^- , etc., where the dominant bonding within the molecular ions is not ionic, but covalent.^{9,10} By starting from *ab initio* quantum-chemistry calculations for the whole molecular ion, and then using the resultant electron charge density to calculate pair potentials, the approach accounts properly for the effects of intramolecular covalency on the intermolecular interactions. For the intramolecular interactions, which are beyond the Gordon-Kim model and the pairwise interaction scheme, a harmonic expansion is employed, with the expansion coefficients also determined by *ab initio* quantum-chemistry calculations. When applied to study the phase transitions in K_2SeO_4 , an important member of the family of A_2BX_4 compounds, the method successfully reproduced, via supercell molecular-dynamics simulations, the transition from the room-temperature paraelectric phase to the lower-temperature ferroelectric superstructure. Also, first-principles lattice-dynamics calculations for the static room-temperature structure clearly revealed the structural instabilities in the system which were in excellent agreement with the experimentally observed soft-phonon behavior of the incommensurate phase transition.

In the present work we apply this method to the study of K_2SO_4 , another member in the A_2BX_4 family. At room temperature, K_2SO_4 is isomorphous with K_2SeO_4 and has the same orthorhombic to hexagonal phase transition at high temperatures.^{11,12} However, it does not have any known incommensurate phase transition at lower temperatures. The only transition below room temperature was found to be at 56 K, and showed no obvious ferroelectric properties.¹³ Thus this low temperature phase transition in K_2SO_4 is definitely different from

that for K_2SeO_4 . It is thus important to see if our *ab initio* theoretical studies can account for the similarities and differences between the two chemically similar systems and reveal the origin of these differences at the level of interionic interactions.

One of the most interesting findings in the previous study¹⁰ of K_2SeO_4 is that the room-temperature *Pnam* (D_{2h}^{16}) structure has not only the instability that drives the incommensurate and ferroelectric cell-tripling transition seen experimentally, but also a zone-center instability. This is a direct reflection of the double-well structure of the potential-energy surface responsible for the observed phase transitions. Therefore, in the present study, we put our emphasis on the lattice dynamics of K_2SO_4 at the zone center. Also we compare our first-principles normal-mode frequencies with the measured Raman frequencies; such a comparison will provide a valuable test of the theoretical description of the system.

Raman scattering measurements on K_2SO_4 have been reported by a number of workers.¹⁴⁻¹⁶ At room temperature, the spectral resolution is rather poor and there is disagreement among these results. In order to achieve higher resolution for better comparison with the theoretical frequencies, we performed Raman scattering measurements at liquid-nitrogen temperature. The experimental details will be given elsewhere.¹⁷ The measured Raman frequencies are listed in Table II. These frequencies agree well with a previous Raman measurement at the same temperature.¹⁶

Our theoretical study started from *ab initio* quantum-chemistry calculations for the whole SO_4^{2-} ion. These calculations did not include a background crystal field. This is justified by the finding that, while the oxygen sites are the only sites with non-negligible crystal fields in the *Pnam* structure of K_2SO_4 , over 90% of the contribution to these fields is from the other ions in the same SO_4^{2-} group, rather than from the rest of the lattice. Although the background crystal field may affect, to some degree, the shapes of the tails of the electron charge distributions on the outskirts of the SO_4^{2-} ion, it is relatively unimportant in determining the electron charge distribution within the SO_4^{2-} ion which is most crucial for both intra-

and intermolecular interactions. A full description of the procedure for obtaining interionic potentials from these quantum-chemistry calculations has been given in Ref. 10. Briefly, we first performed a structural optimization for a SO_4^{2-} ion; the optimized structure is a perfect tetrahedron with the S—O bond length of 2.81 a.u., which is fairly close to the experimental value of 2.72 a.u. at room temperature. Then the second derivatives of the energy with respect to distortions of the SO_4^{2-} ion are calculated to construct a harmonic expansion for the description of the intramolecular interactions. The resultant electron charge density for the SO_4^{2-} ion is decomposed into approximate charge densities for the individual ions.¹⁰ Using these charge densities and the free-ion charge density for the K^+ ion,¹⁸ we then computed the short-range pair potentials between these ions according to the Gordon-Kim¹ electron-gas model. The charges on the ions were +1 for K, +1.6372 for S, and -0.9093 for O, where the fractional charges were obtained from the computed charge density of SO_4^{2-} by a Mulliken population analysis.¹⁹

With the intra- and intermolecular potentials determined, we first performed static relaxation for the room-temperature *Pnam* structure of K_2SO_4 . Our relaxation started from the experimental structure and was subject to the constraints of *Pnam* symmetry, i.e., only the structural parameters listed in Table I were allowed to vary. The energy minimization was for an infinite lattice, obtained by applying periodic boundary conditions, and followed a Newton-Raphson algorithm. The standard technique of the Ewald sum was used for the calculation of the lattice energy and forces, etc.

The relaxation quickly reaches the structure with zero forces on the basis ions and zero stresses. The parameters for the relaxed structure are given in Table I, along with the experimental values. As can be seen, all the lattice constants in the theoretical structure are shorter than the experimental values,¹¹ 5% for *a*, 4% for *b*, and 3%

for *c*, which is a rather general feature for large unit cell simulations using Gordon-Kim potentials. Although these deviations are certainly not negligible in magnitude, the fact that they are percentage-wise comparable in all three directions indicates that the theoretical structure is almost uniformly contracted and thus still preserves the basic physical character of the system.

The theoretical basis parameters are very close to the experimental values. The largest difference is 0.02 for x/a of *O*(2,1), which gives a displacement of about 0.15 Å, well below the thermal fluctuations in the positions of these ions at room temperature.

In the previous study¹⁰ of K_2SeO_4 a similar *Pnam*-constrained static relaxation was performed and equally good agreement with the experimental structure was found, which seemed to indicate the stability of the *Pnam* structure. However, when a molecular-dynamics relaxation was performed for a single *Pnam* unit cell (without the symmetry constraints), the system transformed into a monoclinic structure, which directly revealed the zone-center instability in the room-temperature phase. We therefore performed a molecular-dynamics relaxation for a *Pnam* unit cell of K_2SO_4 . Our relaxation follows a constant (zero) -pressure algorithm²⁰ with the time step size of 0.005 ps. Starting from the experimental room-temperature structure, we gradually reduced the kinetic energy of the sample and thus reached $T=0$ K. We found that the atomic positions and the potential energy of the relaxed structure are almost identical with those for the theoretical *Pnam* structure given in Table I. Consequently, there is no structural transformation and therefore no zone-center instabilities, as will also be apparent from subsequent lattice-dynamical calculations.

Next we computed the normal-mode frequencies at the zone center for the theoretical *Pnam* structure of K_2SO_4 . The resultant normal modes are classified according to symmetry and the Raman-active frequencies are listed in Table II, along with our measured Raman frequencies at liquid-nitrogen temperature and the percentage differences $\delta = [(\omega_t - \omega_e)/\omega_e] \times 100\%$, with the ω_t and ω_e being the theoretical and experimental frequencies, respectively. No imaginary frequencies were found, and therefore there is no zone-center instability in the system, consistent with the molecular-dynamics result. This constitutes a major difference between K_2SO_4 and K_2SeO_4 .

The agreement between the theoretical and experimental frequencies in Table II is generally acceptable, considering the parameter-free nature of the potentials. In terms of the absolute values, the lattice modes seem to agree better than the internal modes that originate from the internal vibrations of the SO_4^{2-} molecular ions. But the percentage differences are comparable, and for almost all modes are within $\pm 20\%$, with only a few exceptions. It should be pointed out that since the present calculation was for a statically fully relaxed structure, it did not include the effects of temperature and therefore direct comparison with experiment can be misleading. We found that if we increase the lattice constants in the theoretical structure by about 0.1% to their values at 80 K (obtained by a molecular-dynamics simulation) to include temperature effects, the fit between theory and experiment clearly

TABLE I. Experimental (Ref. 11) and theoretical structural parameters for the *Pnam* structure of K_2SO_4 . (The lattice constants are given in angstroms.)

Parameters	Experiment	Theory
<i>a</i>	7.497	7.097
<i>b</i>	10.083	9.667
<i>c</i>	5.735	5.581
<i>x/a</i> of <i>S</i> (1)	0.2315	0.2330
<i>y/b</i> of <i>S</i> (1)	0.4208	0.4195
<i>x/a</i> of <i>K</i> (1,1)	0.1755	0.1683
<i>y/b</i> of <i>K</i> (1,1)	0.0884	0.0882
<i>x/a</i> of <i>K</i> (2,1)	0.9906	0.9897
<i>y/b</i> of <i>K</i> (2,1)	0.7052	0.7008
<i>x/a</i> of <i>O</i> (2,1)	0.2899	0.3100
<i>y/b</i> of <i>O</i> (2,1)	0.5576	0.5613
<i>x/a</i> of <i>O</i> (3,1)	0.0454	0.0258
<i>y/b</i> of <i>O</i> (3,1)	0.4157	0.4229
<i>x/a</i> of <i>O</i> (1,1)	0.2996	0.3010
<i>y/b</i> of <i>O</i> (1,1)	0.3530	0.3468
<i>z/c</i> of <i>O</i> (1,1)	0.0426	0.0348

TABLE II. Experimental (four of the B_{1g} modes were not observed due to their weak intensities) and calculated Raman frequencies (in cm^{-1}) for $Pnam$ K_2SO_4 .

A_g			B_{1g}			B_{2g}			B_{3g}		
Expt.	Calc.	Err. (%)	Expt.	Calc.	Err. (%)	Expt.	Calc.	Err. (%)	Expt.	Calc.	Err. (%)
52	58	12	92	109	18	78	69	-11	74	34	-54
97	119	23	96	140	46	107	105	-2	94	73	-22
105	129	23	110	147	34	141	121	-14	105	120	14
113	141	25	117	169	44	147	161	10	144	154	7
134	170	27		196		165	199	21	148	180	22
162	187	15		208							
172	220	28		225							
450	491	9	455	507	11	457	486	6	458	500	9
617	681	10	620	687	11	623	684	10	622	691	11
630	693	10	635	707	11						
989	1049	6		1049							
1099	1208	10	1114	1219	9	1115	1214	9	1112	1210	9
1152	1220	6	1171	1231	5						

improves.

Furthermore, it would be unwise to infer from these discrepancies, that the potential-energy surface has only comparable accuracy, since there is good reason to believe that such differences are to be expected. This stems from the fact that oxygen ions are very light. As a consequence, when moving at low frequencies, in a shallow well, they travel relatively very large distances and the motion is highly anharmonic. Consequently the curvature at the bottom of the well, which corresponds to our theoretical frequencies, may give only a semiquantitative agreement with the observed Raman frequency. As an example, the worst discrepancy, between 74 (observed) and 34 cm^{-1} (theory) for the lowest B_{3g} mode, is almost certainly explicable in these terms. At the other extreme, as has been noted, the very-high-frequency internal modes display (in absolute terms) even worse discrepancies. However, this is well known in quantum chemistry, and is again due to anharmonicity owing to the light oxygens having large zero-point amplitudes which (a) displace their mean positions from the potential well minima, and (b) soften the effective force constant to reflect the average curvature of the potential well at the ground-state level. In addition, these factors will also indirectly influence the low-frequency motions.

The zone-center instability found in K_2SeO_4 was due to the rotation of the SeO_4^{2-} ions as rigid bodies from their orientations in the room-temperature structure.¹⁰ Thus we compared the intermolecular potentials in K_2SO_4 with

those in K_2SeO_4 for an explanation for the absence of such zone-center instability in K_2SO_4 . It was found that the short-range potentials in the two compounds are essentially the same in their effective ranges and that the only differences are the effective ionicities for the atoms within the molecular ions. In K_2SeO_4 , they are 1.1504 for Se and -0.7876 for O,¹⁰ while in K_2SO_4 they are 1.6370 for S and -0.9093 for O. A close inspection of $Pnam$ structures of these compounds shows that of all the intermolecular ion pairs, the K-O interactions are the most important, since they are the shortest and are directly responsible for the rotation of the molecular ions. Therefore the difference between the lattice-dynamics properties in K_2SeO_4 and in K_2SO_4 can be mainly attributed to the difference in the ionicities of the oxygen atoms, which originates from the different chemistry between SeO_4^{2-} and SO_4^{2-} ions.

In summary, we have performed static relaxation, molecular-dynamics, and lattice-dynamics calculations for the room-temperature structure of the K_2SO_4 crystal. Unlike the isomorphous K_2SeO_4 , no zone-center instability has been found. This implies that any soft-phonon type of phase transition below room temperature has to be associated with a non-zone-center instability and the resultant structure should be a superlattice.

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